

A novel analytical model of the magnetic field configuration in the Galactic Center explaining the diffuse gamma-ray emission

Mehmet Guenduez

Julia Becker Tjus

Katia Ferrière

Ralf-Jürgen Dettmar

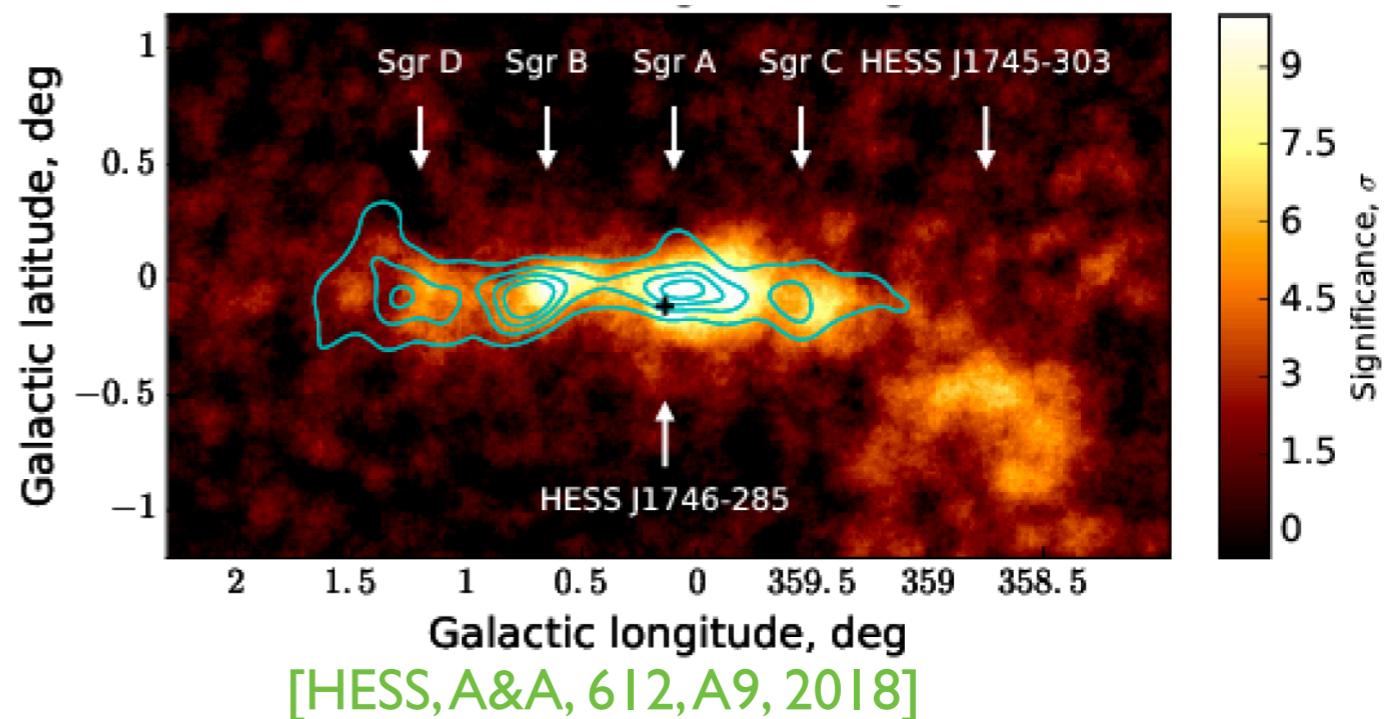
Dominik J. Bomans

Madison

July 26th 2019

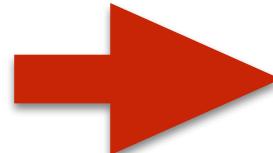
Motivation:

- In April 2016 the H.E.S.S. Collaboration reported the detection of a high energy diffuse gamma-ray flux in the Galactic Center from some GeV up to tens of TeV
- No significant allusion to a cut-off

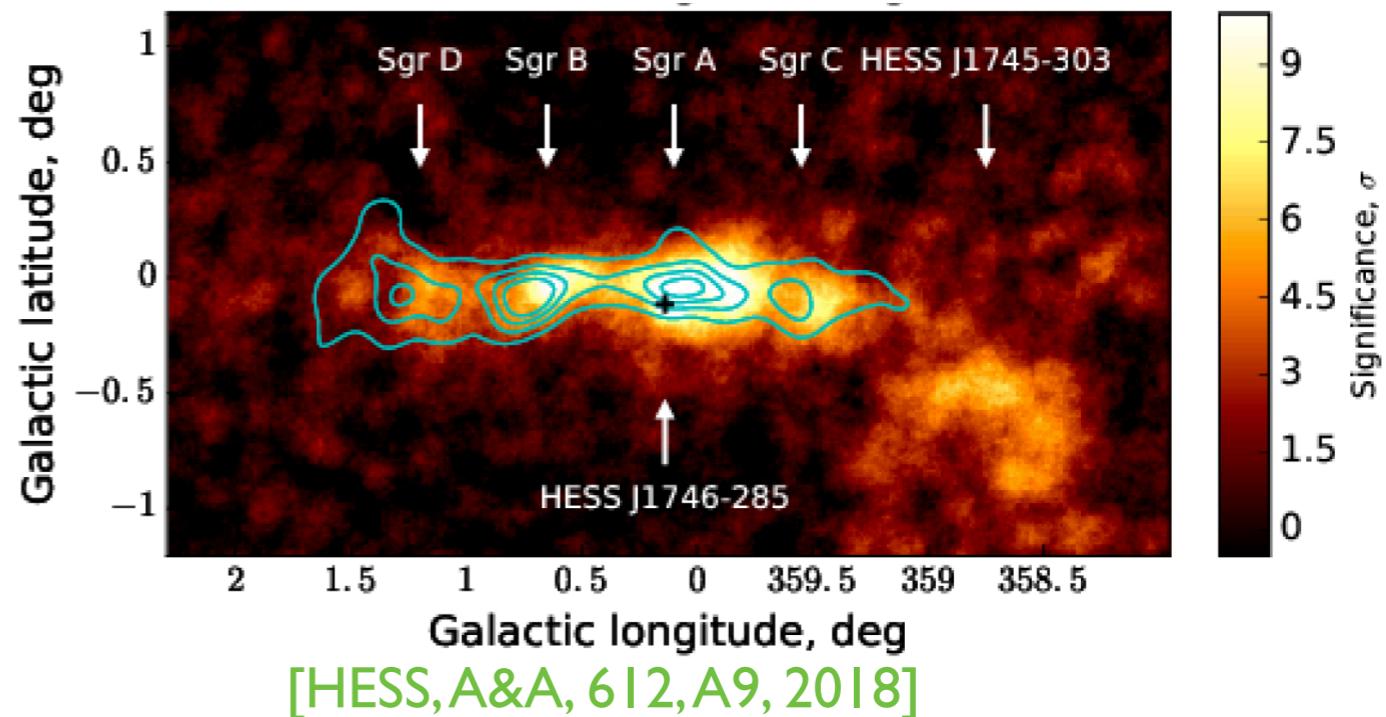


Motivation:

- In April 2016 the H.E.S.S. Collaboration reported the detection of a high energy diffuse gamma-ray flux in the Galactic Center from some GeV up to tens of TeV
- No significant allusion to a cut-off



PeVatron accelerator



[HESS,A&A,612,A9,2018]

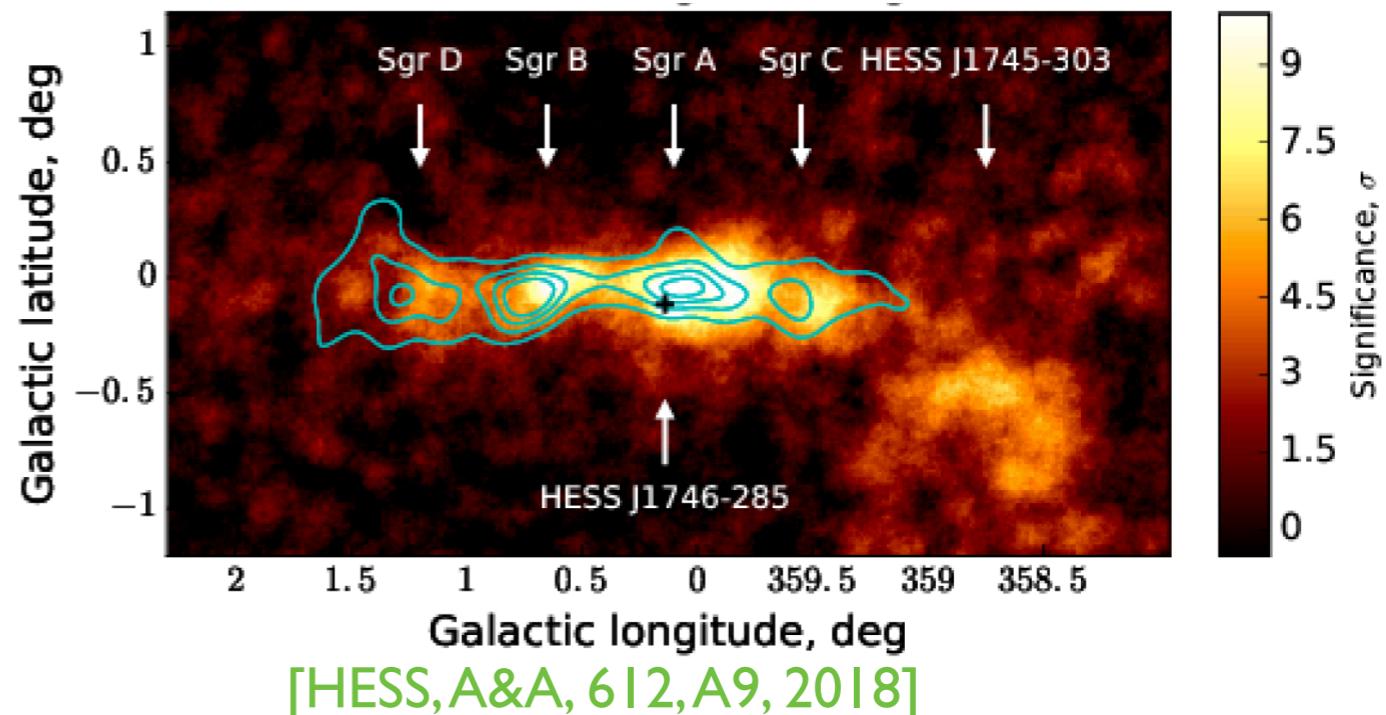
Motivation:

- In April 2016 the H.E.S.S. Collaboration reported the detection of a high energy diffuse gamma-ray flux in the Galactic Center from some GeV up to tens of TeV
- No significant allusion to a cut-off



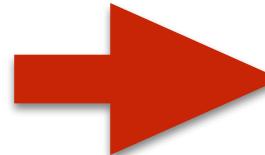
PeVatron accelerator

Semi-analytical model

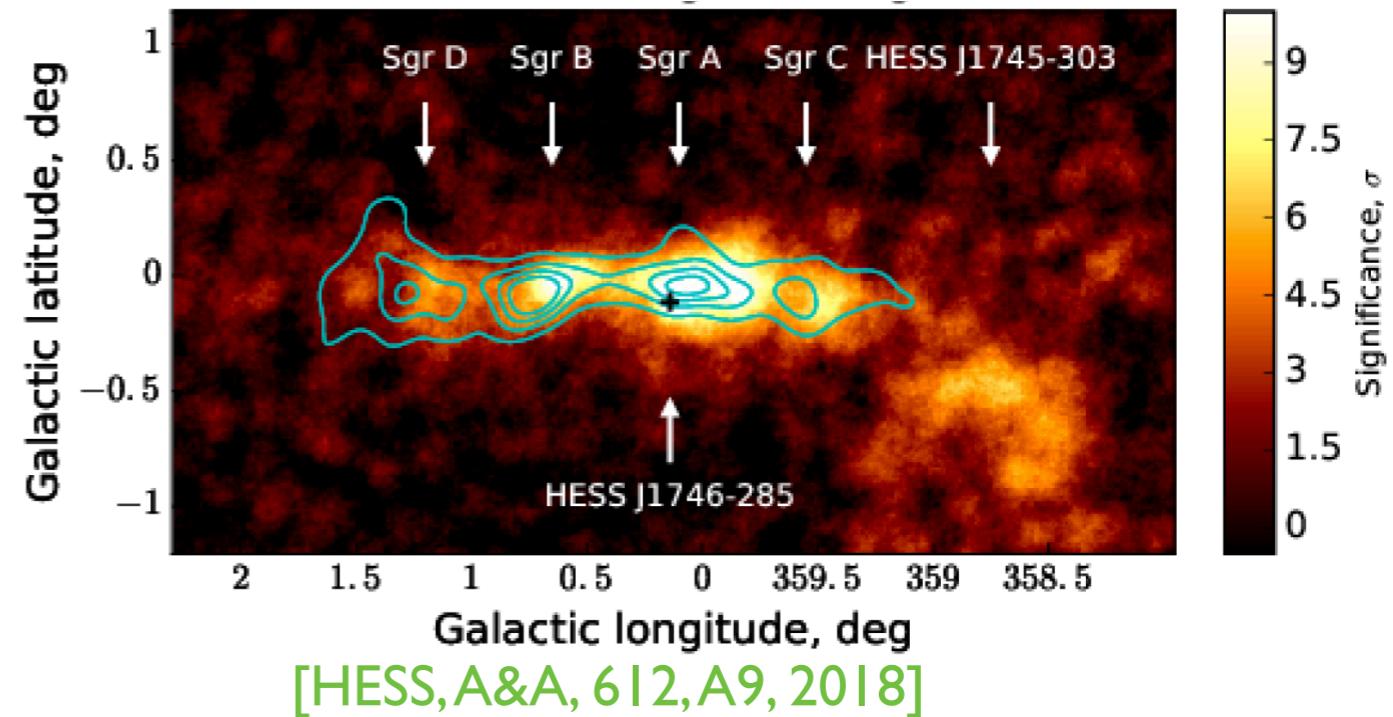


Motivation:

- In April 2016 the H.E.S.S. Collaboration reported the detection of a high energy diffuse gamma-ray flux in the Galactic Center from some GeV up to tens of TeV
- No significant allusion to a cut-off



PeVatron accelerator

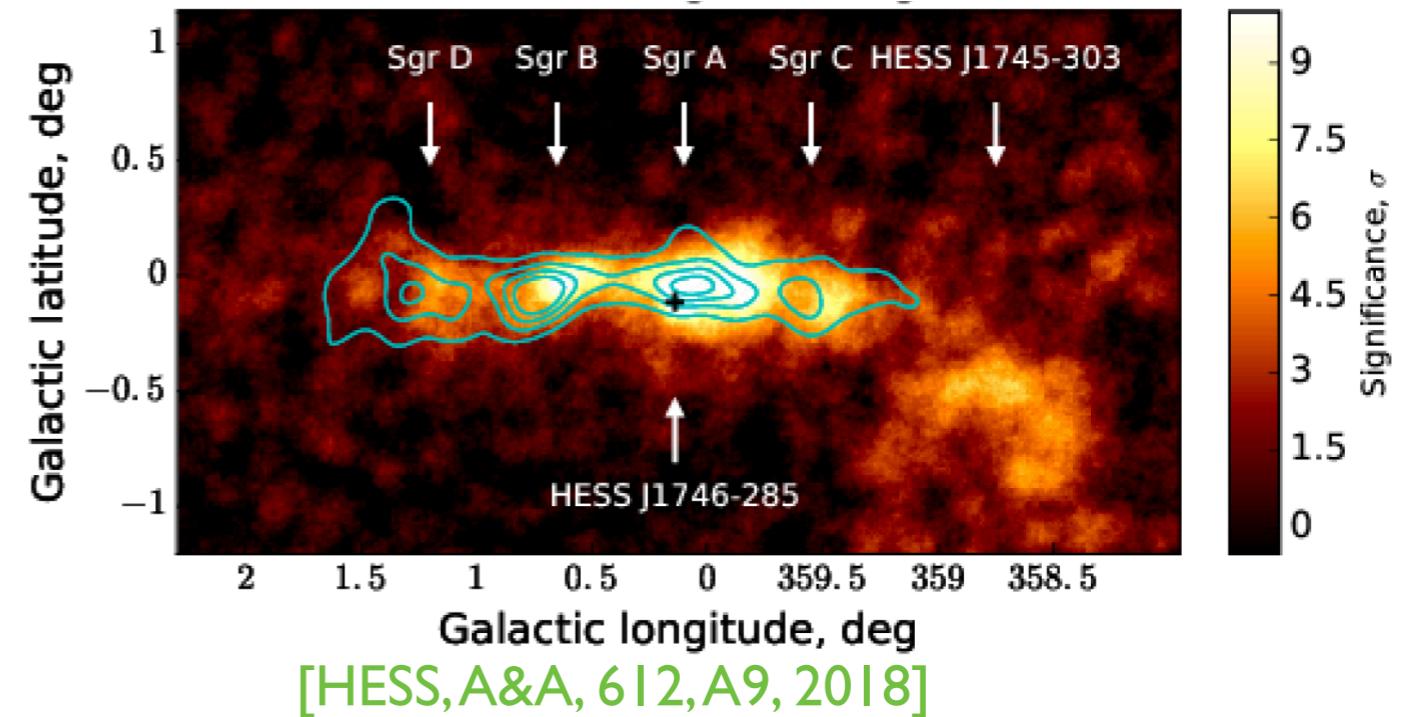


Semi-analytical model

- Diffusion with Kolmogorov spectrum
- spatially and energy dependent continuous loss and source distribution
- static and spherically symmetric
- injection at the center
- Proton-proton interaction

Motivation:

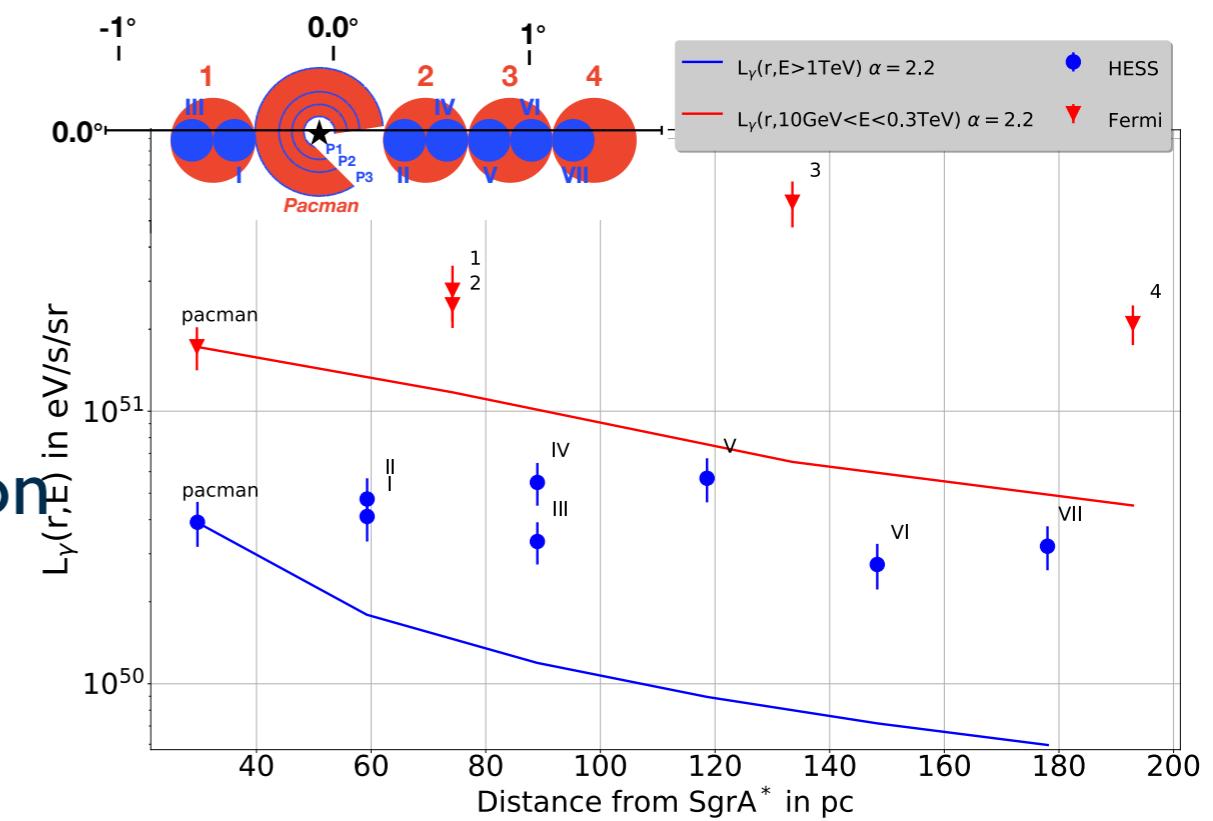
- In April 2016 the H.E.S.S. Collaboration reported the detection of a high energy diffuse gamma-ray flux in the Galactic Center from some GeV up to tens of TeV
- No significant allusion to a cut-off



PeVatron accelerator

Semi-analytical model

- Diffusion with Kolmogorov spectrum
- spatially and energy dependent continuous loss and source distribution
- static and spherically symmetric
- injection at the center
- Proton-proton interaction



Motivation:

Open questions:

Motivation:

Open questions:

- How many sources are sufficient for the observation?
- How much energy is necessary?
- What is the impact of the ambient conditions?

Problem:

- Larger discrepancies at:
 1. higher longitudes
 2. lower energies

Problem:

- Larger discrepancies at:
 1. higher longitudes
 2. lower energies
- dominant horizontal magnetic field
- additional sources?



Problem:

- Larger discrepancies at:
 1. higher longitudes
 2. lower energies
- dominant horizontal magnetic field
- additional sources?



Move away from semi-analytical approximation

Problem:

- Larger discrepancies at:
 1. higher longitudes
 2. lower energies
- dominant horizontal magnetic field
- additional sources?

Move away from semi-analytical approximation



A novel model of the magnetic field configuration
and strength in the Galactic Center

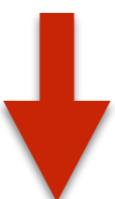
Problem:

- Larger discrepancies at:
 1. higher longitudes
 2. lower energies
- dominant horizontal magnetic field
- additional sources?

Move away from semi-analytical approximation



A novel model of the magnetic field configuration
and strength in the Galactic Center



Accurate modeling of the gas distribution

Gas distribution

An accurate 3D distribution has not been modeled yet!

Gas distribution

An accurate 3D distribution has not been modeled yet!

New model

Gas distribution

An accurate 3D distribution has not been modeled yet!

Can be split into 3 components:

New model



1. diffuse medium [Ferrière et al, A&A, 467, 611-627, 2007]
2. 10 pc around SgrA* [Ferrière, A&A, 540, A50, 2012]
3. molecular clouds

Gas distribution

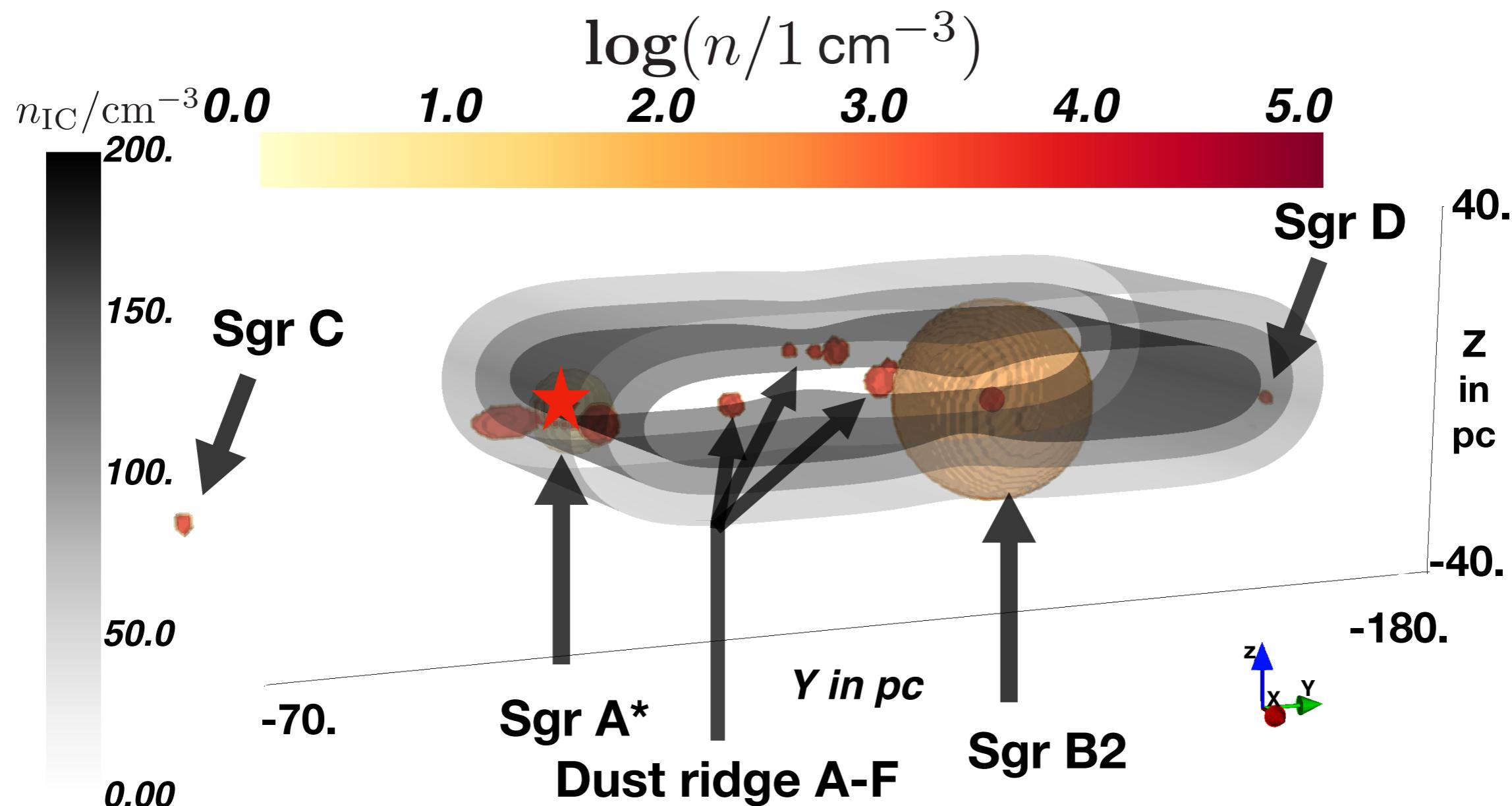
An accurate 3D distribution has not been modeled yet!

Can be split into 3 components:

New model

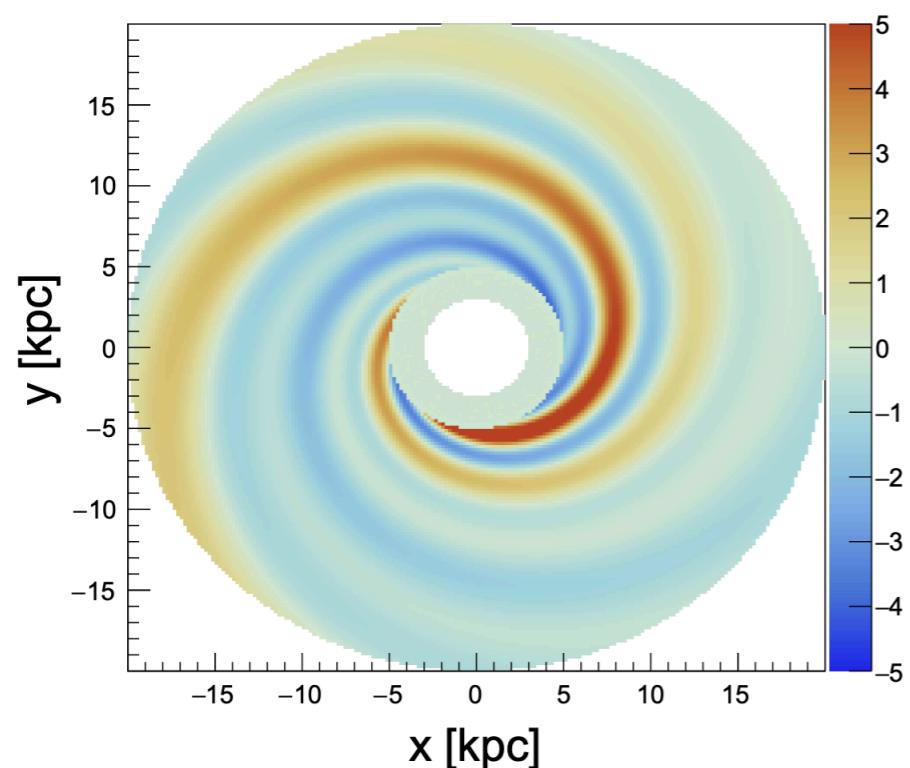


1. diffuse medium [Ferrière et al, A&A, 467, 611-627, 2007]
2. 10 pc around SgrA* [Ferrière, A&A, 540, A50, 2012]
3. molecular clouds

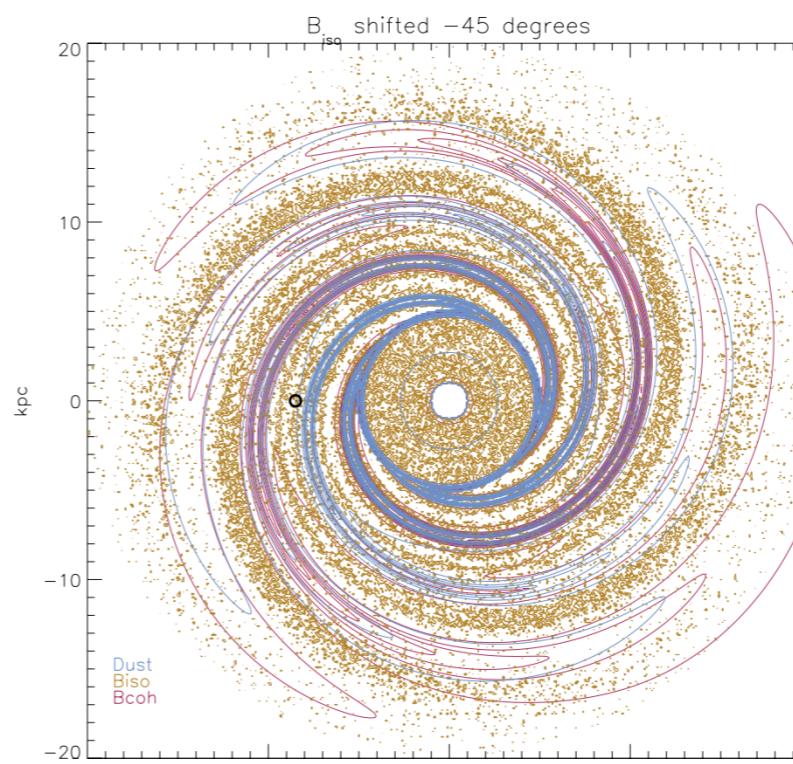


Magnetic Field in the GC

Status of research:



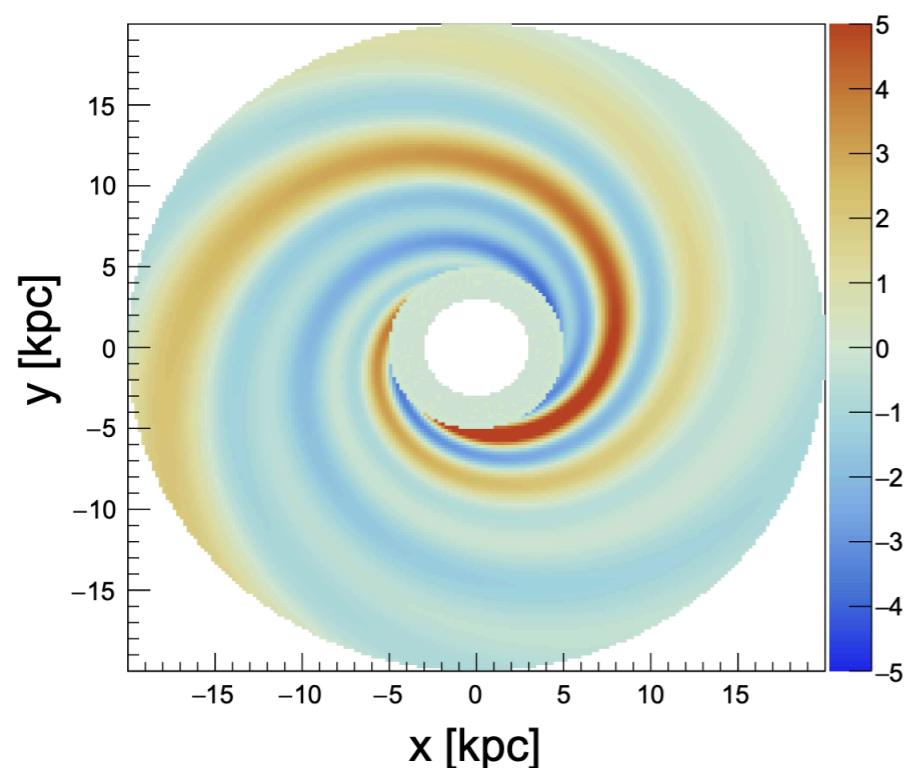
[Unger & Farrar, EPJWC, 210, 4005, 2019]



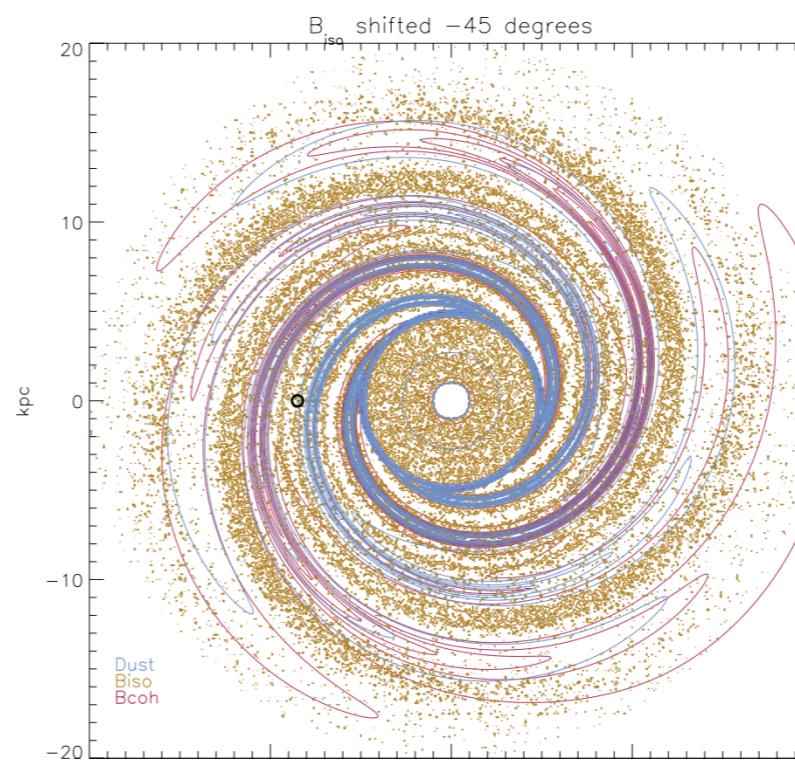
[Jaffe et al., MNRAS, 431, 683, 2013]

Magnetic Field in the GC

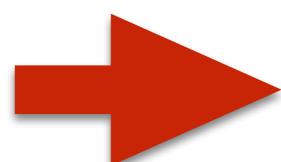
Status of research:



[Unger & Farrar, EPJWC, 210, 4005, 2019]



[Jaffe et al., MNRAS, 431, 683, 2013]



lack in the Galactic Center

Magnetic Field in the GC- GBFDI9

- Variation range between $\sim 1e-5$ G and $\sim 5e-3$ G
- **Large-scale** magnetic field in the GC initially **poloidal**
- Then, inside **molecular clouds** were sheared out **horizontally**

Magnetic Field in the GC- GBFDI9

- Variation range between $\sim 1e-5$ G and $\sim 5e-3$ G
- **Large-scale** magnetic field in the GC initially **poloidal**
- Then, inside **molecular clouds** were sheared out **horizontally**

Has not been modeled yet!

Magnetic Field in the GC- GBFDI9

- Variation range between $\sim 1e-5$ G and $\sim 5e-3$ G
- **Large-scale** magnetic field in the GC initially **poloidal**
- Then, inside **molecular clouds** were sheared out **horizontally**

Has not been modeled yet!



Can be split into 2 components:

1. Poloidal:

- Intercloud medium (ICM) → large scale field
- Non-thermal filaments (NTFs) → local field

2. Horizontal:

- Dense molecular clouds (MCs) → local field

Poloidal field

From Ferrière et al. (2014) we consider Model C (FT14-C)}

Poloidal field

From Ferrière et al. (2014) we consider Model C (FT14-C)}



analytical and divergence-free

Poloidal field

From Ferrière et al. (2014) we consider Model C (FT14-C)}



analytical and divergence-free

$$\mathbf{B}^C = \begin{pmatrix} B_r \\ B_\phi \\ B_z \end{pmatrix} = \begin{pmatrix} \frac{2a z}{(1+a z^2)^3} \\ 0 \\ \frac{1}{(1+a z^2)^2} \end{pmatrix} \cdot B_1 \cdot e^{-r/L \cdot \frac{1}{(1+a z^2)}}$$

Polooidal field

From Ferrière et al. (2014) we consider Model C (FT14-C)}



analytical and divergence-free

$$\mathbf{B}^C = \begin{pmatrix} B_r \\ B_\phi \\ B_z \end{pmatrix} = \begin{pmatrix} \frac{2a z}{(1+a z^2)^3} \\ 0 \\ \frac{1}{(1+a z^2)^2} \end{pmatrix} \cdot B_1 \cdot e^{-r/L \cdot \frac{1}{(1+a z^2)}}$$

- a opening of field lines away from the z-axis
- L exponential scale length
- B_1 normalization factor

Polooidal field

From Ferrière et al. (2014) we consider Model C (FT14-C)}



analytical and divergence-free

$$\mathbf{B}^C = \begin{pmatrix} B_r \\ B_\phi \\ B_z \end{pmatrix} = \begin{pmatrix} \frac{2a z}{(1+a z^2)^3} \\ 0 \\ \frac{1}{(1+a z^2)^2} \end{pmatrix} \cdot B_1 \cdot e^{-r/L \cdot \frac{1}{(1+a z^2)}}$$

- a → opening of field lines away from the z-axis } → geometry
 L → exponential scale length } → observation
 B_1 → normalization factor } → observation

Horizontal field

Starting with Euler's potentials:

$$\vec{B} = \vec{\nabla}\alpha \times \vec{\nabla}\beta$$

Horizontal field

Starting with Euler's potentials:

$$\vec{B} = \vec{\nabla}\alpha \times \vec{\nabla}\beta \quad \xrightarrow{\text{red arrow}} \text{analytical and divergence-free}$$

Horizontal field

Starting with Euler's potentials:

$$\vec{B} = \vec{\nabla}\alpha \times \vec{\nabla}\beta \rightarrow \text{analytical and divergence-free}$$

Consider the zonal plane $\rightarrow B_z = 0 \rightarrow \beta = 1$

Horizontal field

Starting with Euler's potentials:

$$\vec{B} = \vec{\nabla}\alpha \times \vec{\nabla}\beta \rightarrow \text{analytical and divergence-free}$$

Consider the zonal plane $\rightarrow B_z = 0 \rightarrow \beta = 1$

$$\rightarrow \vec{B} = \begin{pmatrix} B_r \\ B_\phi \\ 0 \end{pmatrix} = \begin{pmatrix} \frac{1}{r} \frac{\partial}{\partial \phi} \alpha \\ -\frac{\partial}{\partial r} \alpha \\ 0 \end{pmatrix} = \begin{pmatrix} \frac{1}{r} \frac{\partial \alpha}{\partial \psi} \frac{\partial \psi}{\partial \phi} |_r \\ -\frac{\partial \alpha}{\partial \psi} \frac{\partial \psi}{\partial r} |_\phi \\ 0 \end{pmatrix}$$

Horizontal field

Starting with Euler's potentials:

$$\vec{B} = \vec{\nabla}\alpha \times \vec{\nabla}\beta \quad \text{analytical and divergence-free}$$

Consider the zonal plane $\rightarrow B_z = 0 \rightarrow \beta = 1$

$$\rightarrow \vec{B} = \begin{pmatrix} B_r \\ B_\phi \\ 0 \end{pmatrix} = \begin{pmatrix} \frac{1}{r} \frac{\partial}{\partial \phi} \alpha \\ -\frac{\partial}{\partial r} \alpha \\ 0 \end{pmatrix} = \begin{pmatrix} \frac{1}{r} \frac{\partial \alpha}{\partial \psi} \frac{\partial \psi}{\partial \phi} |_r \\ -\frac{\partial \alpha}{\partial \psi} \frac{\partial \psi}{\partial r} |_\phi \\ 0 \end{pmatrix}$$

$$|B_r/B_\phi| = \frac{1}{r} \cdot dr/d\phi|_{\psi,\rho} \stackrel{!}{=} \eta \text{ depends on the MC characteristics}$$

Horizontal field

Starting with Euler's potentials:

$$\vec{B} = \vec{\nabla}\alpha \times \vec{\nabla}\beta \rightarrow \text{analytical and divergence-free}$$

Consider the zonal plane $\rightarrow B_z = 0 \rightarrow \beta = 1$

$$\rightarrow \vec{B} = \begin{pmatrix} B_r \\ B_\phi \\ 0 \end{pmatrix} = \begin{pmatrix} \frac{1}{r} \frac{\partial}{\partial \phi} \alpha \\ -\frac{\partial}{\partial r} \alpha \\ 0 \end{pmatrix} = \begin{pmatrix} \frac{1}{r} \frac{\partial \alpha}{\partial \psi} \frac{\partial \psi}{\partial \phi} |_r \\ -\frac{\partial \alpha}{\partial \psi} \frac{\partial \psi}{\partial r} |_\phi \\ 0 \end{pmatrix}$$

$$|B_r/B_\phi| = \frac{1}{r} \cdot dr/d\phi|_{\psi,\rho} \stackrel{!}{=} \eta \text{ depends on the MC characteristics}$$

$$\rightarrow \psi = \phi \pm \eta^{-1} \ln(r/\rho) \rightarrow \frac{\partial \alpha}{\partial \psi} = \rho \cdot B_r(\psi)$$

Horizontal field

Starting with Euler's potentials:

$$\vec{B} = \vec{\nabla}\alpha \times \vec{\nabla}\beta \rightarrow \text{analytical and divergence-free}$$

Consider the zonal plane $\rightarrow B_z = 0 \rightarrow \beta = 1$

$$\rightarrow \vec{B} = \begin{pmatrix} B_r \\ B_\phi \\ 0 \end{pmatrix} = \begin{pmatrix} \frac{1}{r} \frac{\partial}{\partial \phi} \alpha \\ -\frac{\partial}{\partial r} \alpha \\ 0 \end{pmatrix} = \begin{pmatrix} \frac{1}{r} \frac{\partial \alpha}{\partial \psi} \frac{\partial \psi}{\partial \phi} |_r \\ -\frac{\partial \alpha}{\partial \psi} \frac{\partial \psi}{\partial r} |_\phi \\ 0 \end{pmatrix}$$

$$|B_r/B_\phi| = \frac{1}{r} \cdot dr/d\phi|_{\psi,\rho} \stackrel{!}{=} \eta \text{ depends on the MC characteristics}$$

$$\rightarrow \psi = \phi \pm \eta^{-1} \ln(r/\rho) \rightarrow \partial \alpha / \partial \psi = \rho \cdot B_r(\psi)$$

we define: $B_r(\psi) = B_1 \cdot \cos(\psi) \cdot h(z)$

Horizontal field

Starting with Euler's potentials:

$$\vec{B} = \vec{\nabla}\alpha \times \vec{\nabla}\beta \rightarrow \text{analytical and divergence-free}$$

Consider the zonal plane $\rightarrow B_z = 0 \rightarrow \beta = 1$

$$\rightarrow \vec{B} = \begin{pmatrix} B_r \\ B_\phi \\ 0 \end{pmatrix} = \begin{pmatrix} \frac{1}{r} \frac{\partial}{\partial \phi} \alpha \\ -\frac{\partial}{\partial r} \alpha \\ 0 \end{pmatrix} = \begin{pmatrix} \frac{1}{r} \frac{\partial \alpha}{\partial \psi} \frac{\partial \psi}{\partial \phi} |_r \\ -\frac{\partial \alpha}{\partial \psi} \frac{\partial \psi}{\partial r} |_\phi \\ 0 \end{pmatrix}$$

$$|B_r/B_\phi| = \frac{1}{r} \cdot dr/d\phi|_{\psi,\rho} \stackrel{!}{=} \eta \text{ depends on the MC characteristics}$$

$$\rightarrow \psi = \phi \pm \eta^{-1} \ln(r/\rho) \rightarrow \partial \alpha / \partial \psi = \rho \cdot B_r(\psi)$$

we define: $B_r(\psi) = B_1 \cdot \cos(\psi) \cdot h(z)$

$$\rightarrow \text{net magnetic flux}=0 \rightarrow h(z) \text{ is arbitrary}$$

Horizontal field

Starting with Euler's potentials:

$$\vec{B} = \vec{\nabla}\alpha \times \vec{\nabla}\beta \rightarrow \text{analytical and divergence-free}$$

Consider the zonal plane $\rightarrow B_z = 0 \rightarrow \beta = 1$

$$\rightarrow \vec{B} = \begin{pmatrix} B_r \\ B_\phi \\ 0 \end{pmatrix} = \begin{pmatrix} \frac{1}{r} \frac{\partial}{\partial \phi} \alpha \\ -\frac{\partial}{\partial r} \alpha \\ 0 \end{pmatrix} = \begin{pmatrix} \frac{1}{r} \frac{\partial \alpha}{\partial \psi} \frac{\partial \psi}{\partial \phi} |_r \\ -\frac{\partial \alpha}{\partial \psi} \frac{\partial \psi}{\partial r} |_\phi \\ 0 \end{pmatrix}$$

$$|B_r/B_\phi| = \frac{1}{r} \cdot dr/d\phi|_{\psi,\rho} \stackrel{!}{=} \eta \text{ depends on the MC characteristics}$$

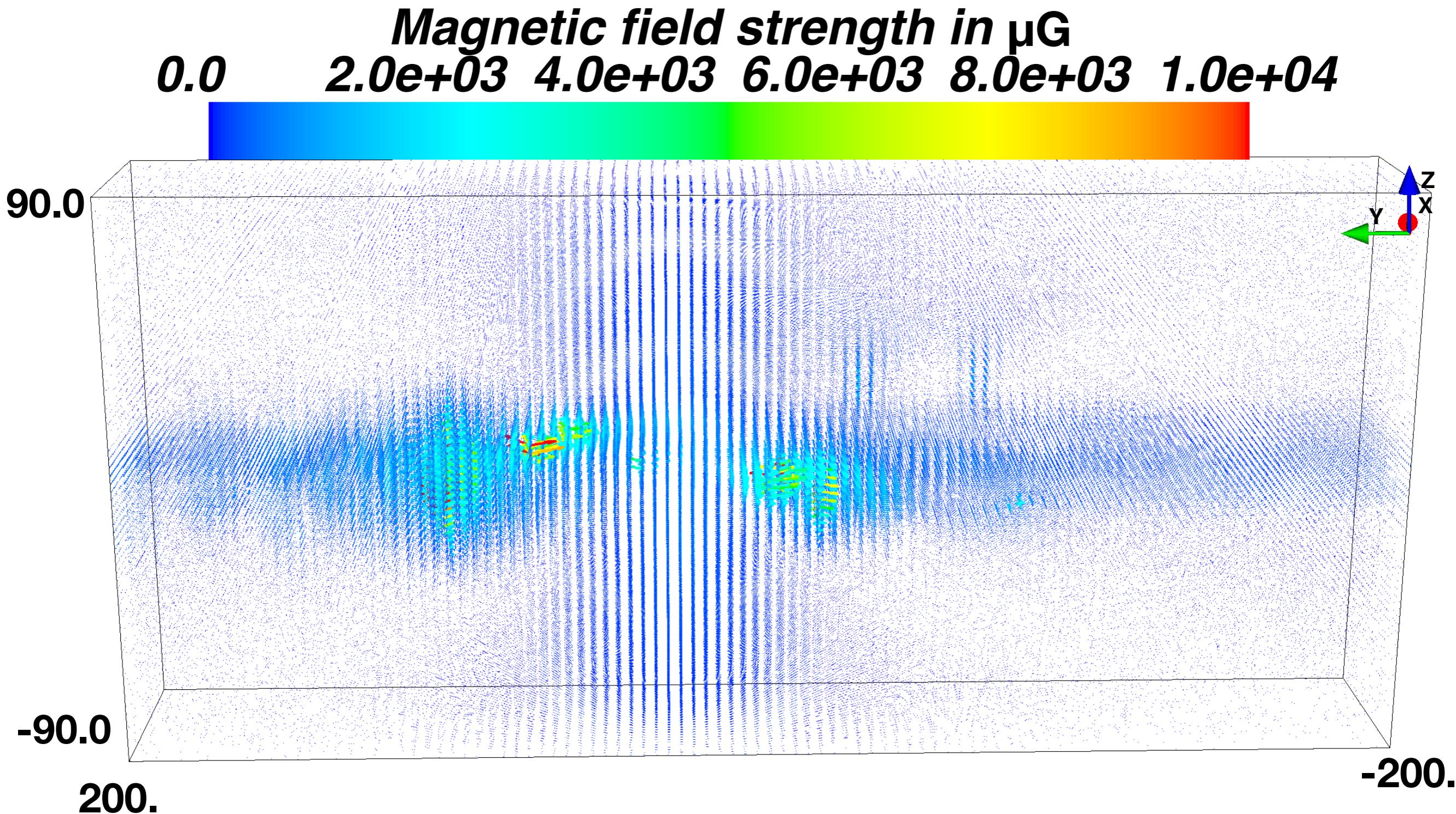
$$\rightarrow \psi = \phi \pm \eta^{-1} \ln(r/\rho) \rightarrow \partial \alpha / \partial \psi = \rho \cdot B_r(\psi)$$

we define: $B_r(\psi) = B_1 \cdot \cos(\psi) \cdot h(z)$

\rightarrow net magnetic flux=0 \rightarrow $h(z)$ is arbitrary

$$\mathbf{B}_{\text{tot}} = \mathbf{B}_{\text{IC}}^C + \sum_{i=1}^8 \mathbf{B}_{\text{NTF},i}^C + \sum_{i=1}^{12} \mathbf{B}_{\text{MC},i}$$

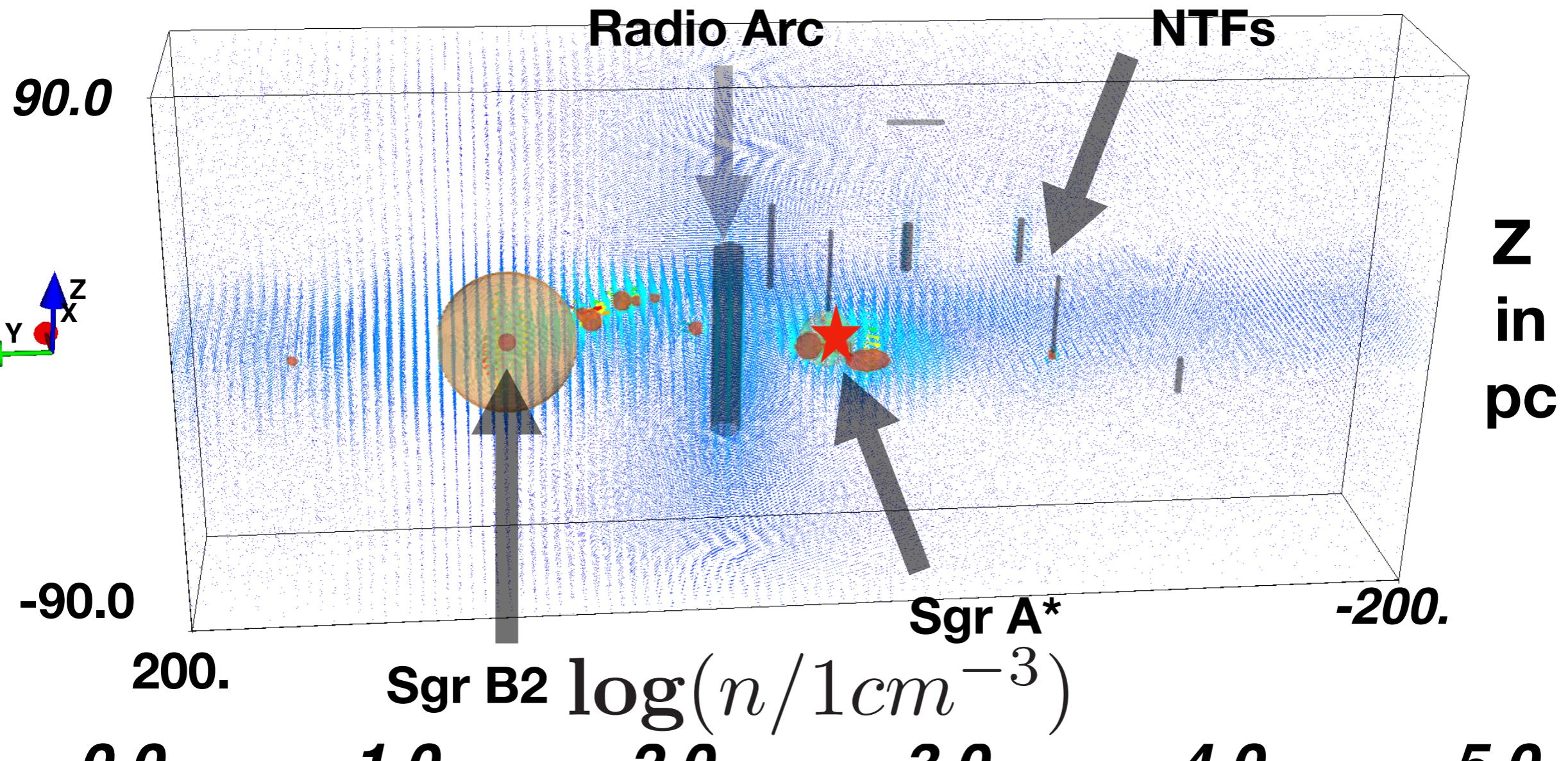
Magnetic Field in the GC - GBFDI9



Magnetic Field in the GC - GBFDI9

Magnetic field strength in μG

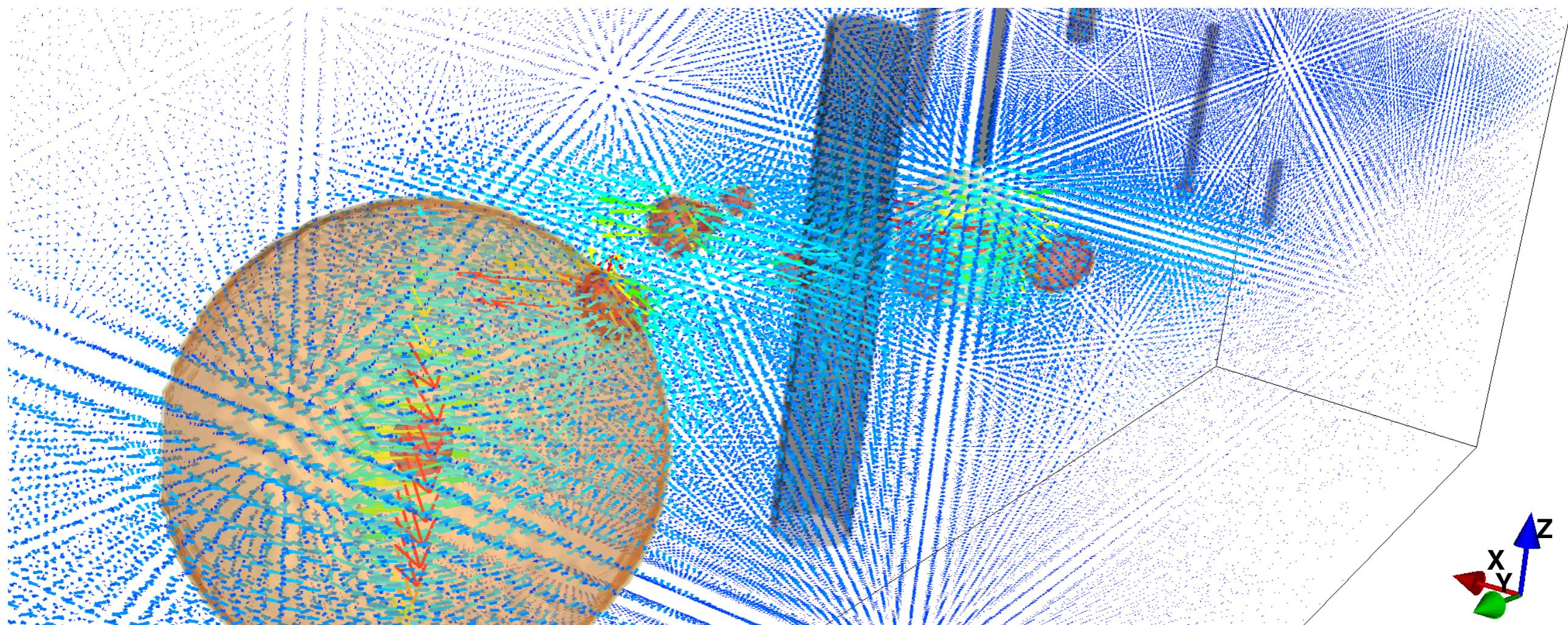
0.0 2.0e+03 4.0e+03 6.0e+03 8.0e+03 1.0e+04



Magnetic Field in the GC - GBFDI9

Magnetic field strength in μG

0.0 2.0e+03 4.0e+03 6.0e+03 8.0e+03 1.0e+04

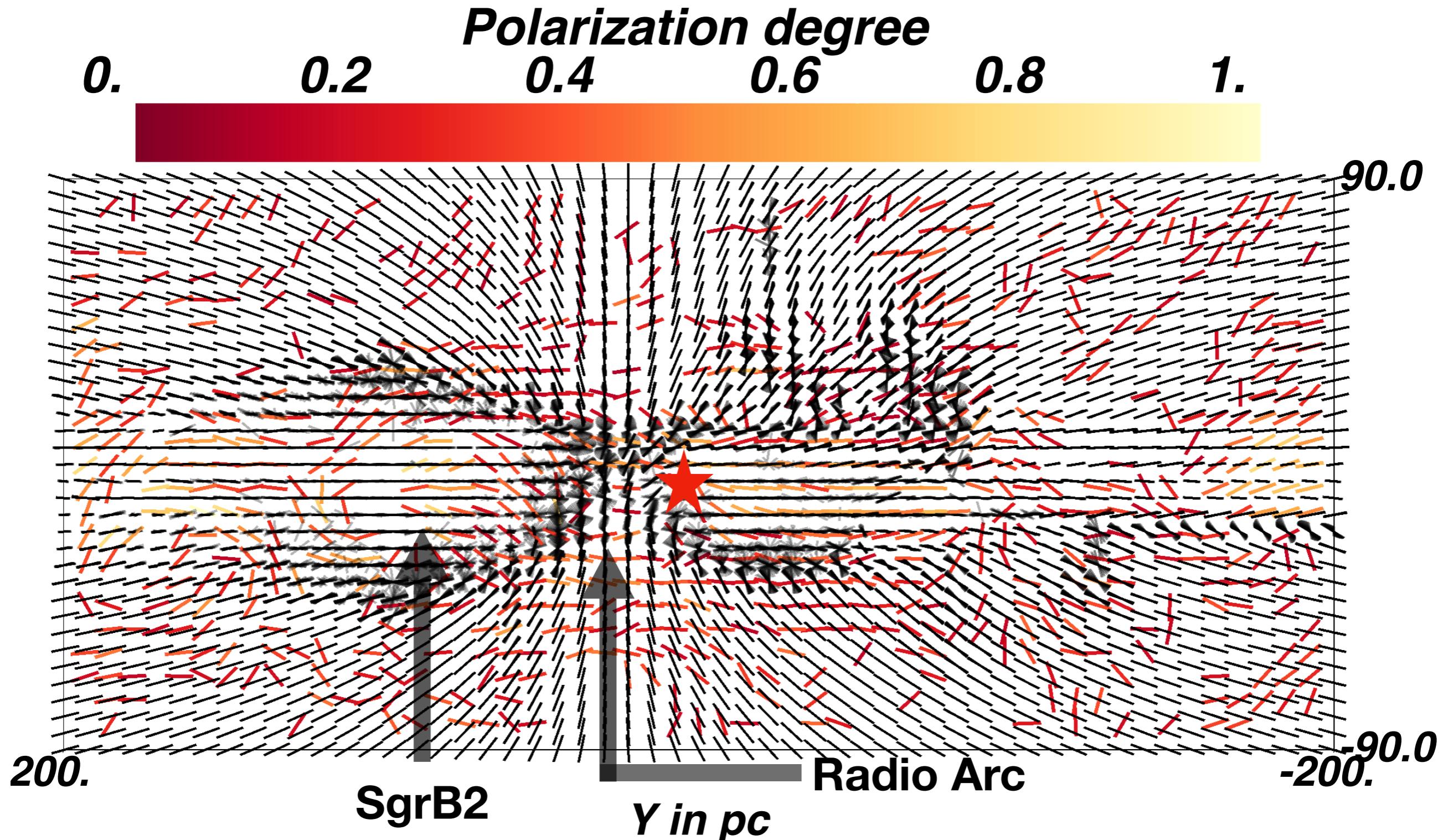


$\log(n/1\text{cm}^{-3})$

0.0 1.0 2.0 3.0 4.0 5.0

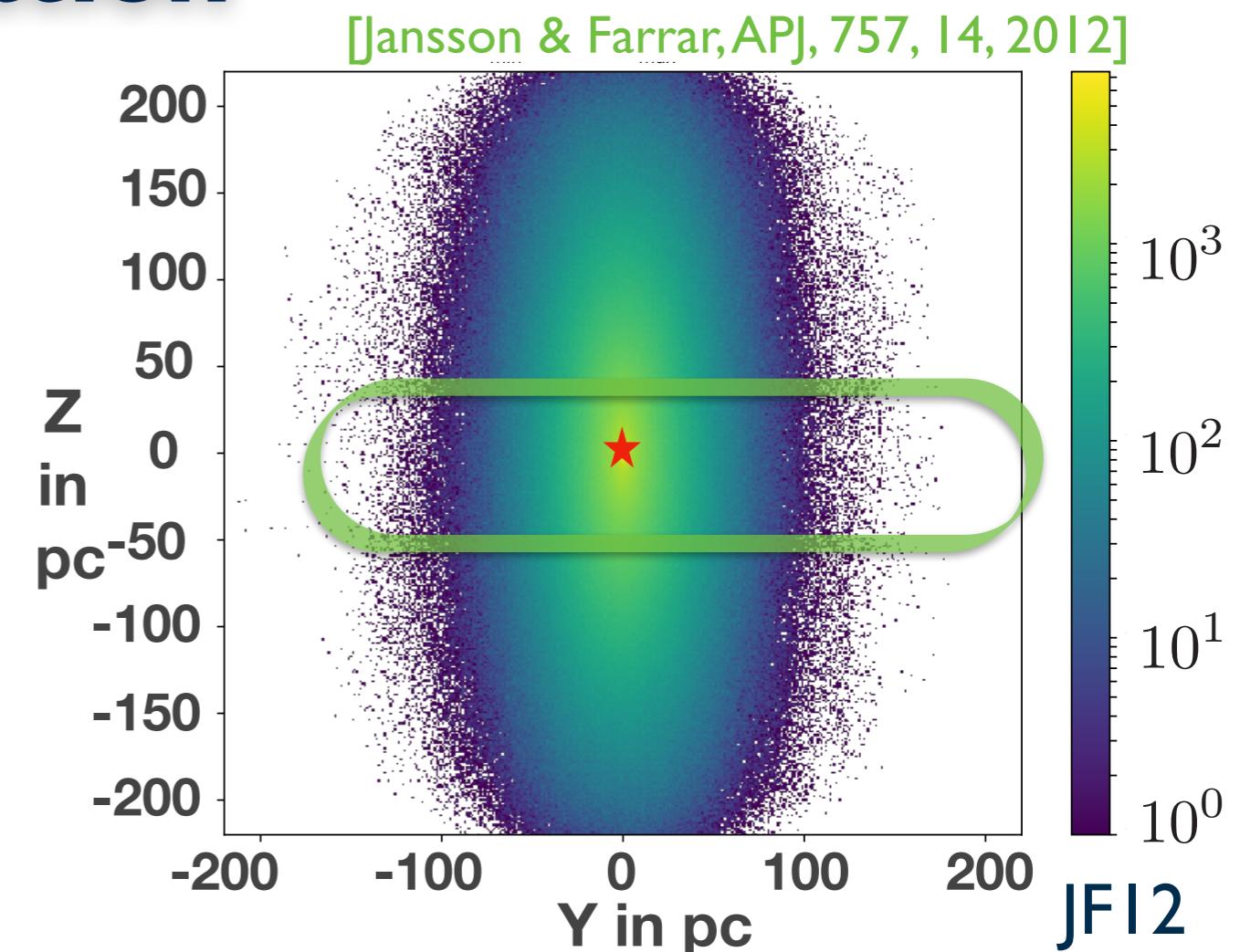
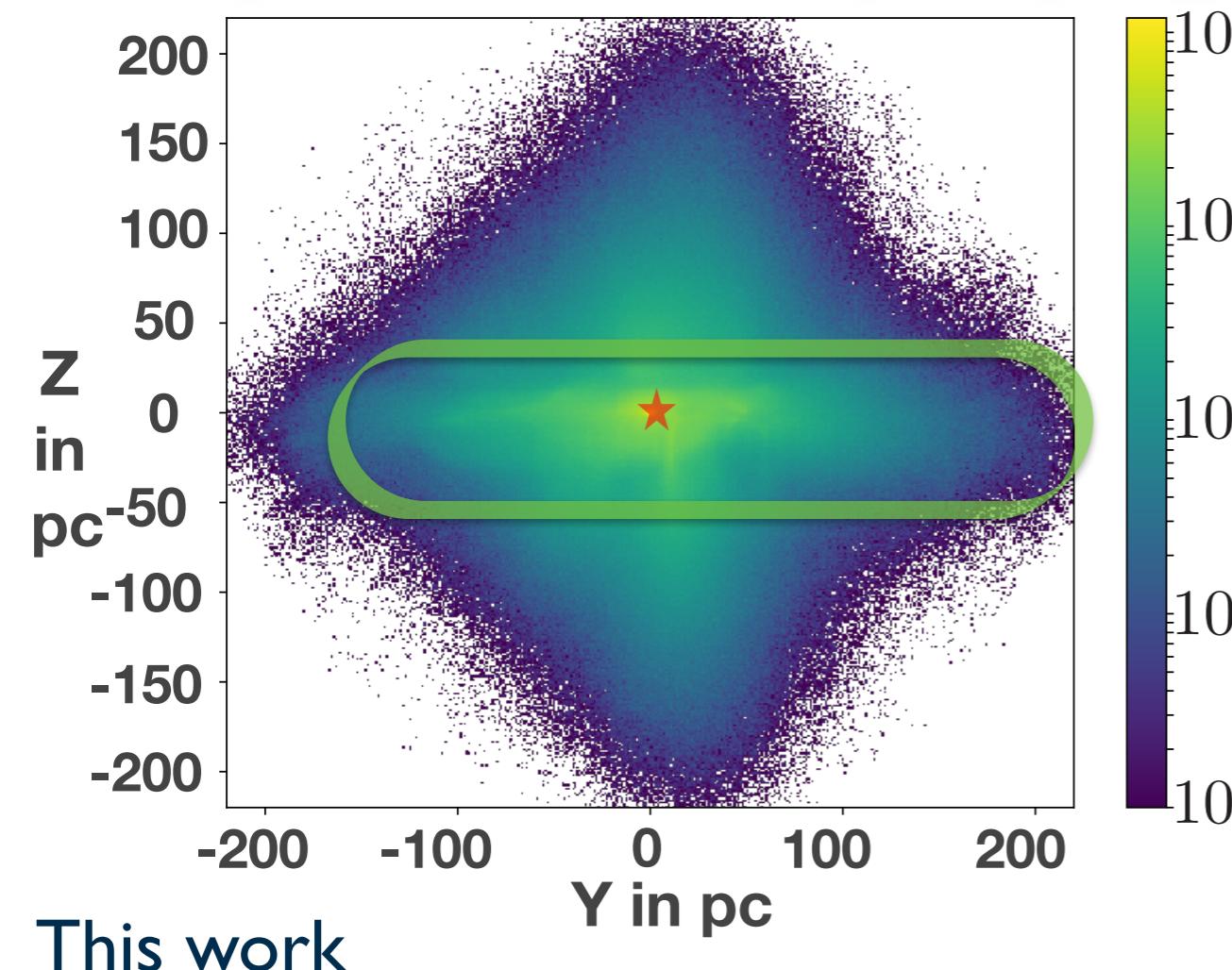


Magnetic Field in the GC - GBFDI9



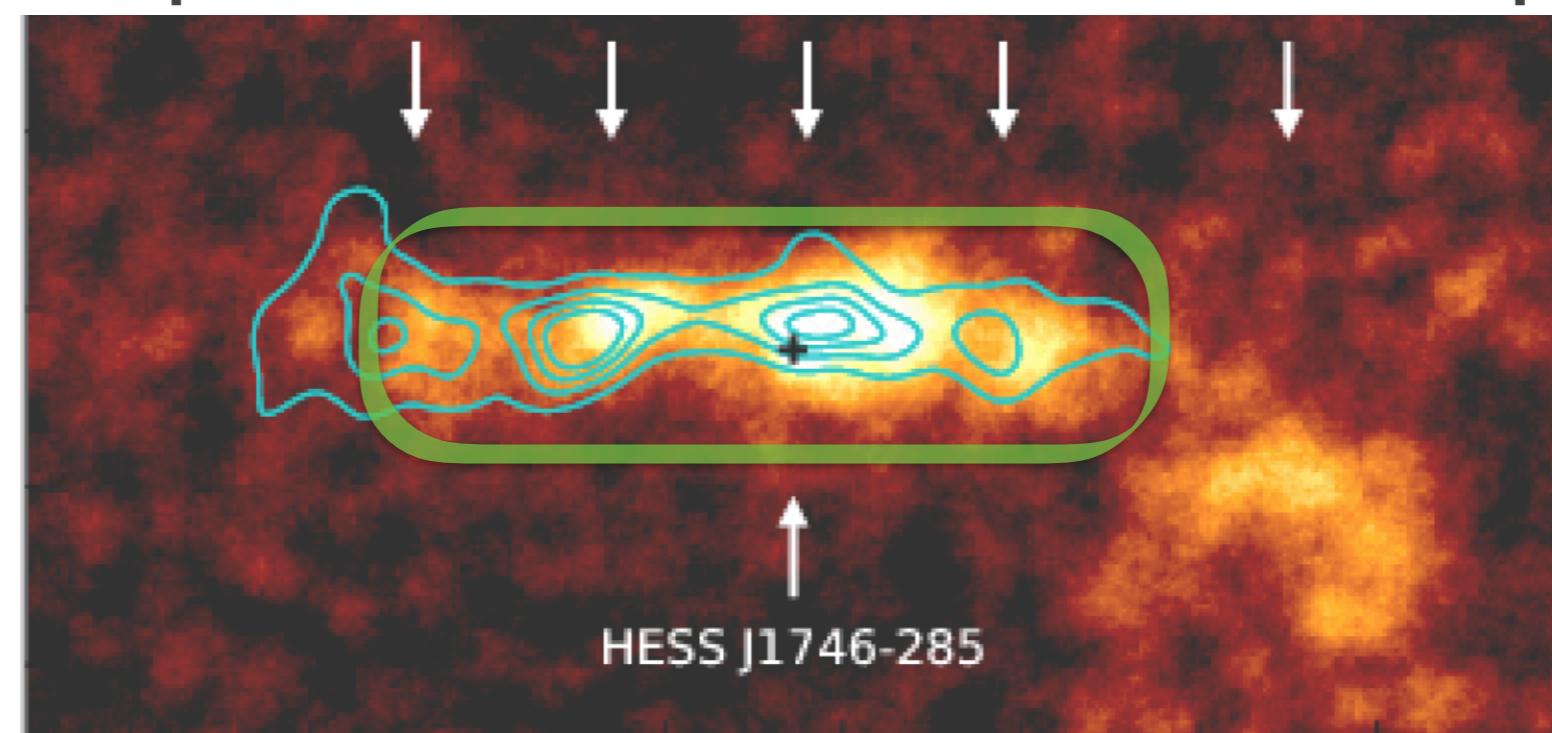
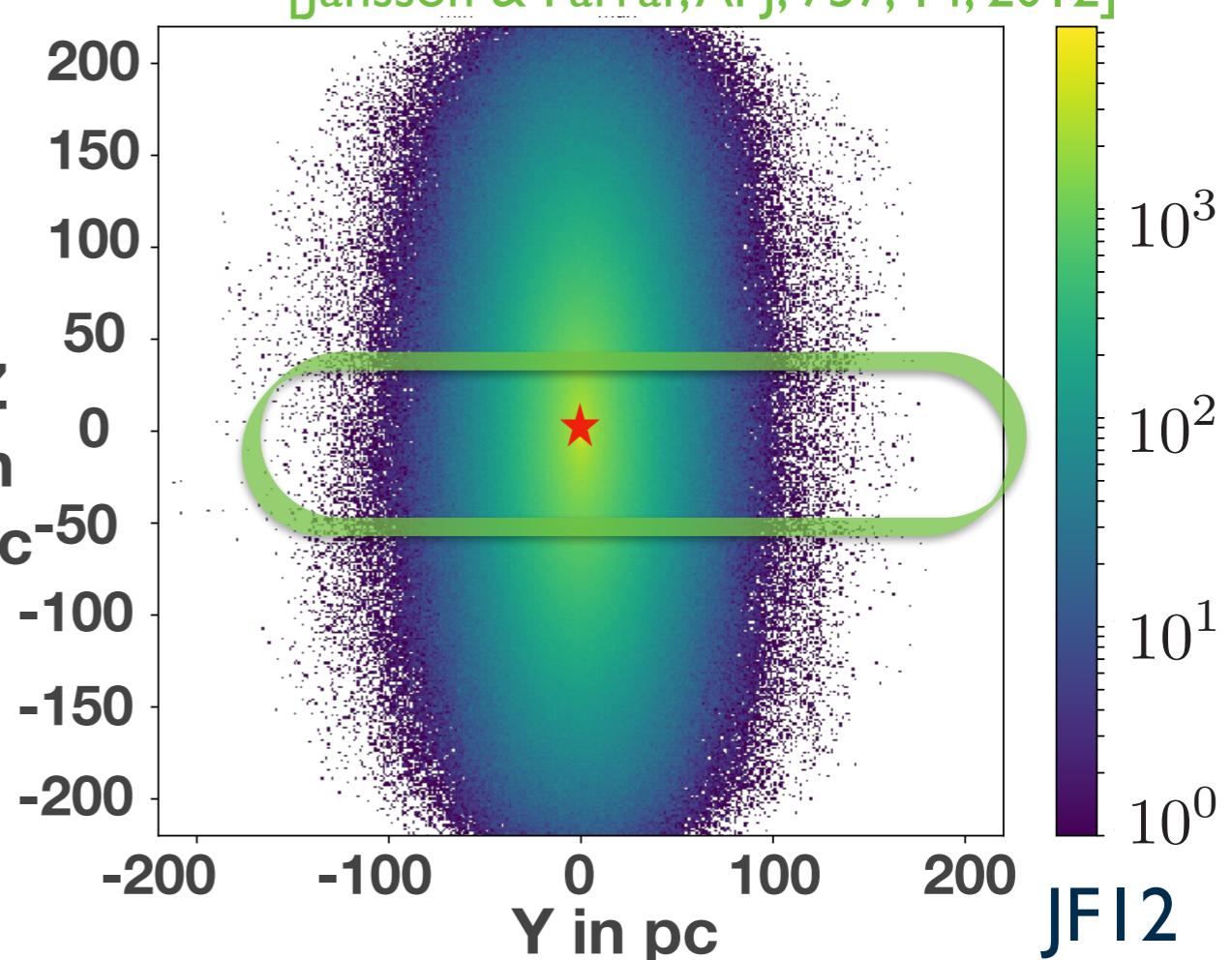
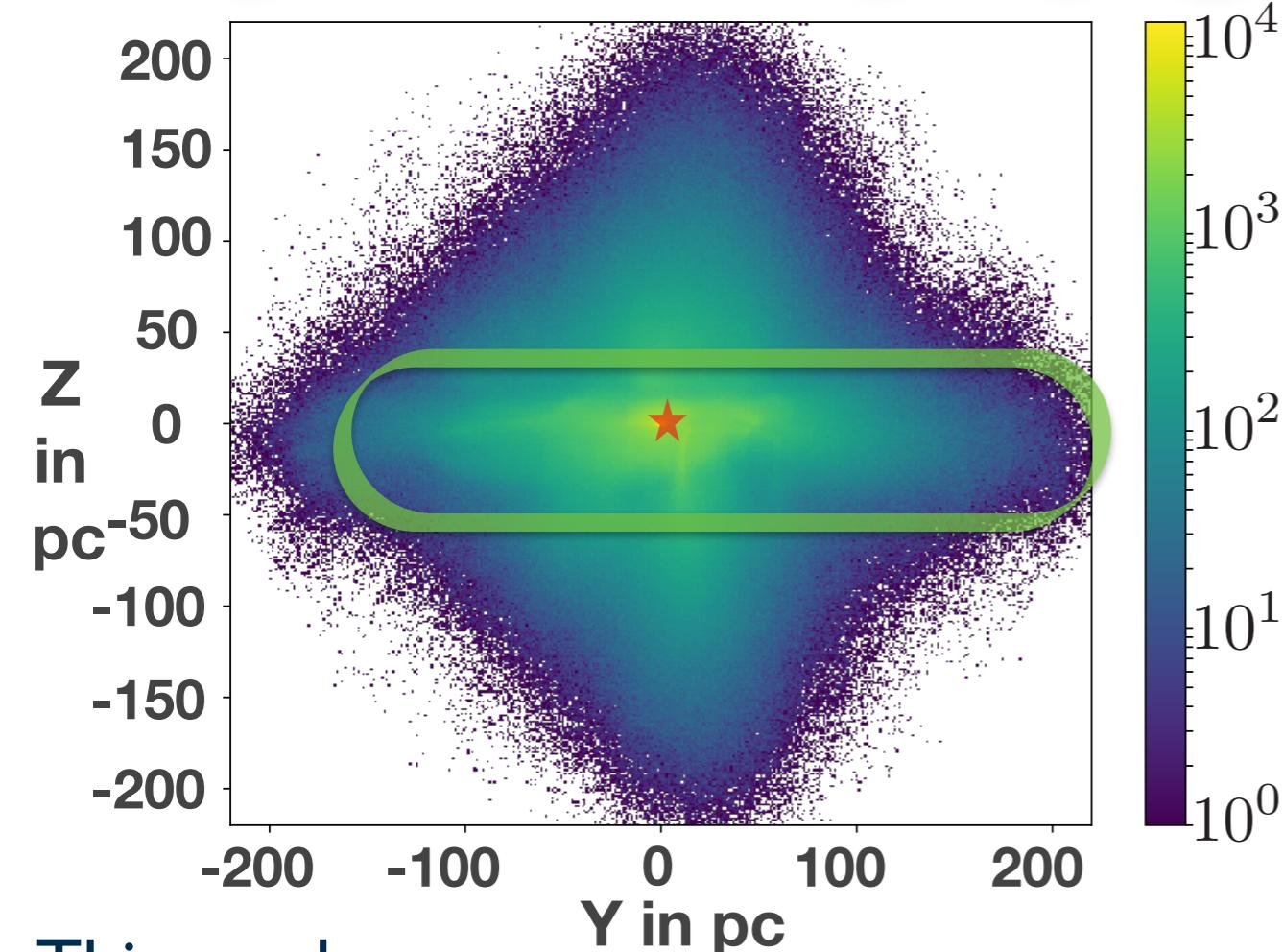
Colored data from [Nishiyama et al., APJ Letters, 712, L23, 2010]

Impact on CR propagation



Impact on CR propagation

[Jansson & Farrar, APJ, 757, 14, 2012]



Summary and conclusion

- A centrally located source is not easily sufficient for the observed gamma-ray emission
- *GBDF19*:
 - could explain the diffuse gamma-ray emission caused by a single centralized source
 - is strong in non-thermal filaments and molecular clouds
 - is compatible with polarization data and corresponds to the observations
 - has a significant impact on the longitudinal profiles of CR propagation
 - can be combined with any other Galactic B-field models
- The 3D gas distribution and *GBDF19* make a more accurate calculation accessible for the GC investigators
- For detailed modeling: M.Guenduez et al. 2019 Submitted to A&A, arXiv:1906.05211

Backup slides

- Proton as primary cosmic-rays injected from SgrA* at the origin
 - spherically symmetric description
 - diffusion dominated scenario is favored
 - spatial and energy dependent source
 - energy dependent continuous loss
- For the first time:

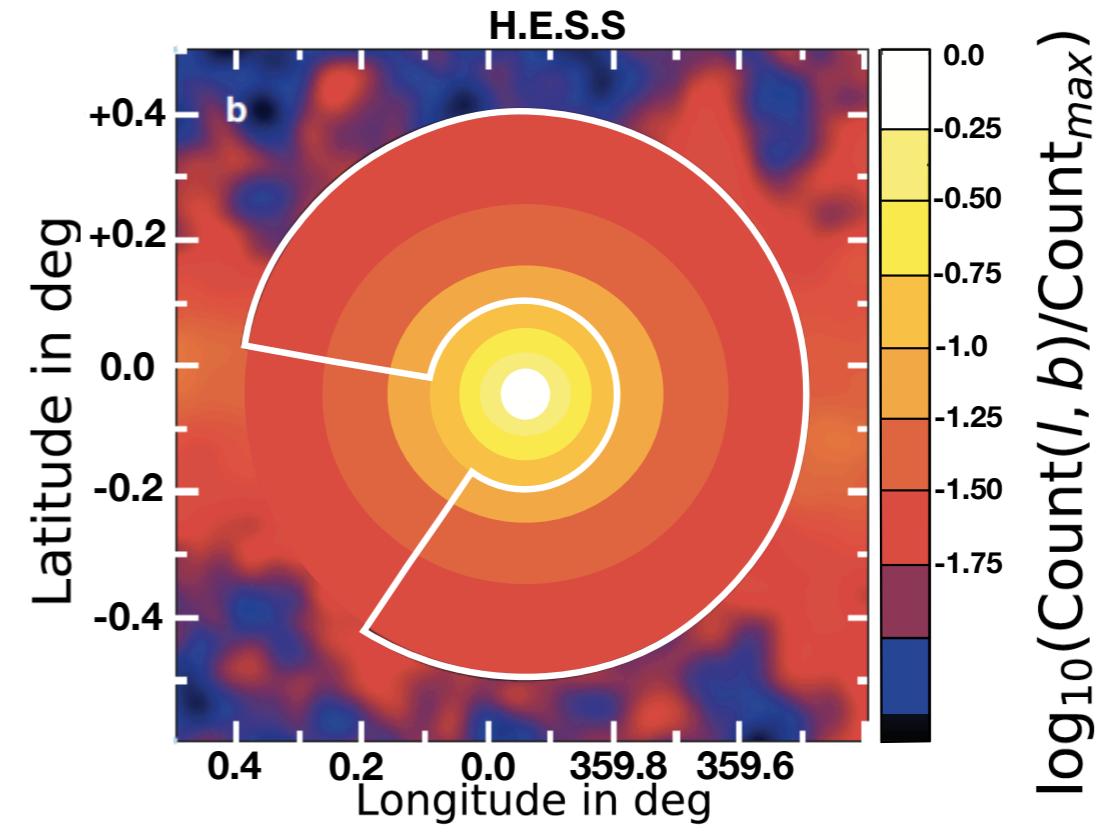
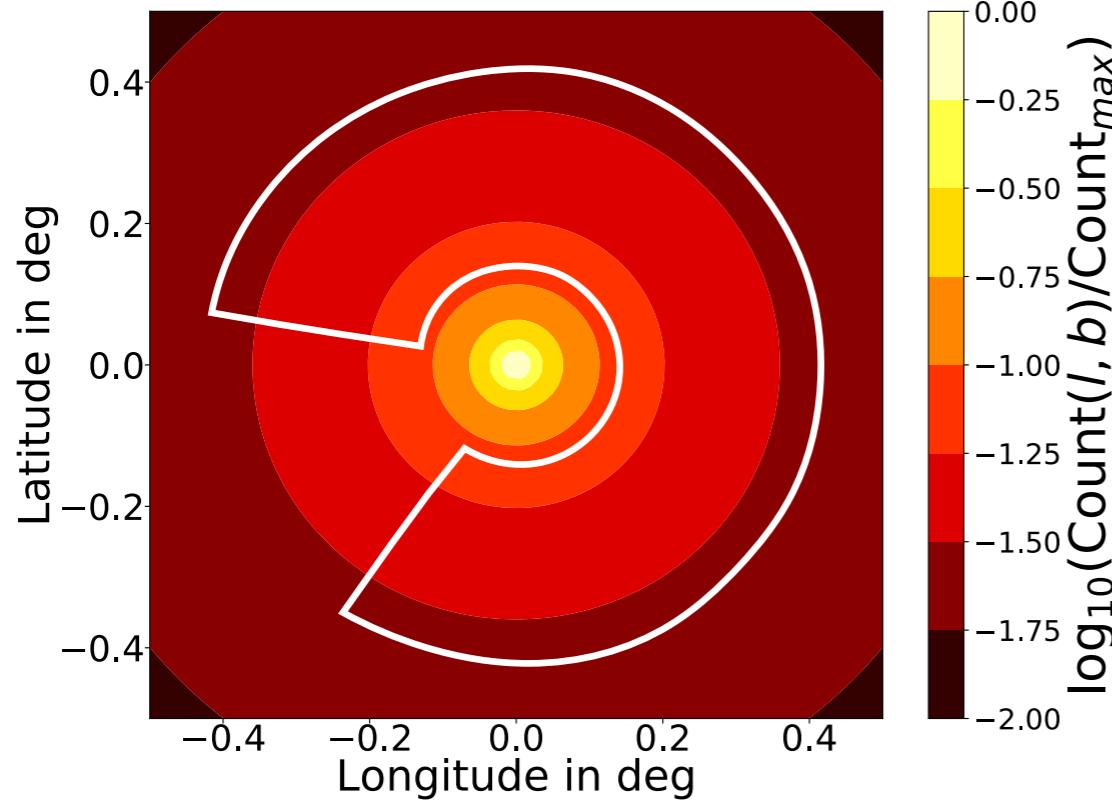
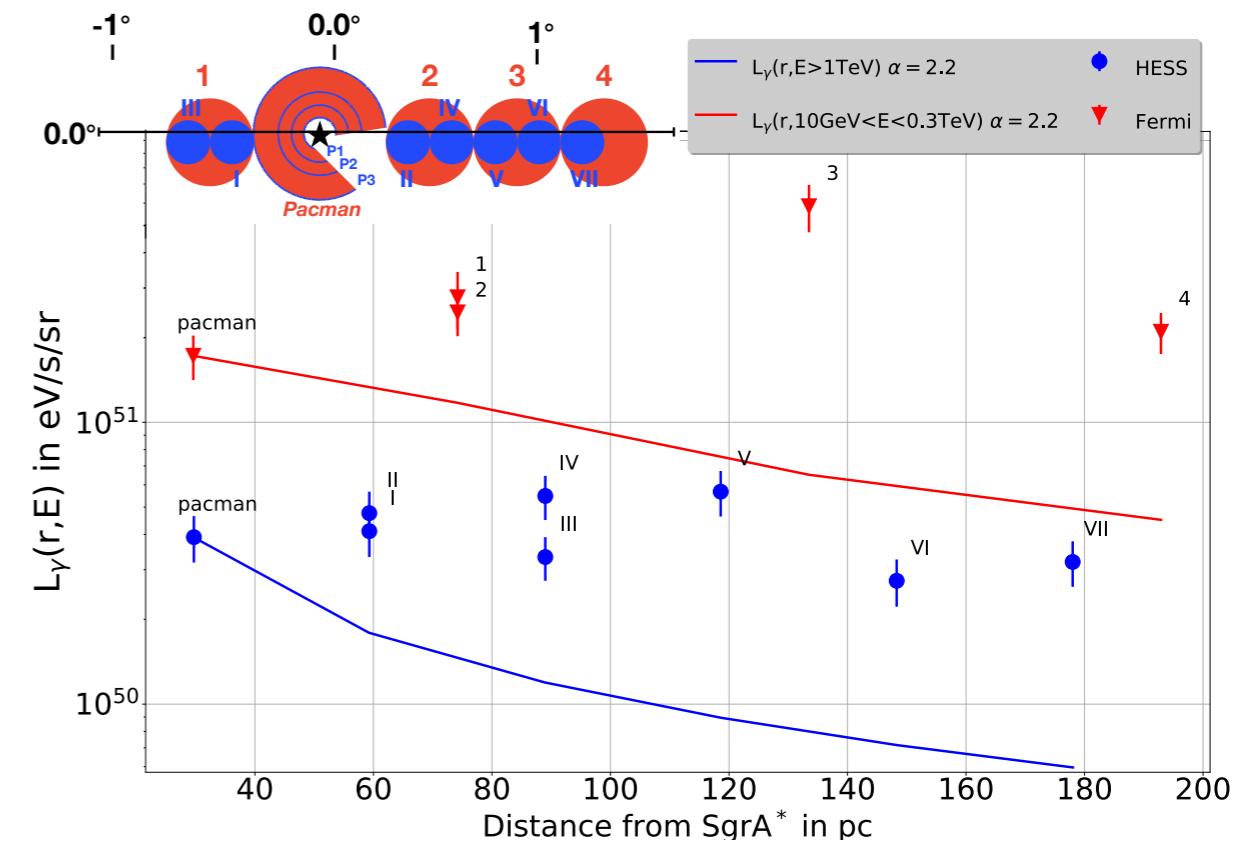
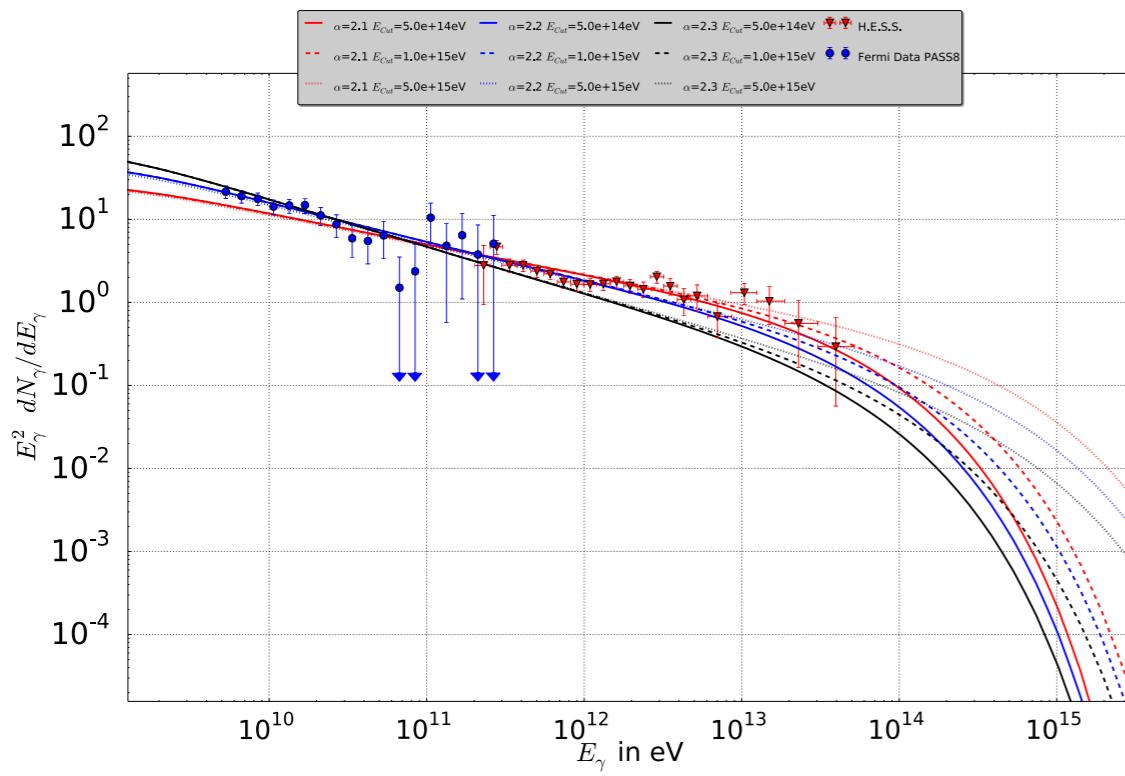
- we include a **spatially** dependent term in the hadronic pion production interaction term

$$-D_0 \gamma^\nu \frac{2}{r} \frac{\partial}{\partial r} n(r, \gamma) - D_0 \gamma^\nu \frac{\partial^2}{\partial r^2} n(r, \gamma) - \frac{\partial}{\partial \gamma} b(r) \gamma^{1+\mu} n(r, \gamma) = Q(r, \gamma)$$

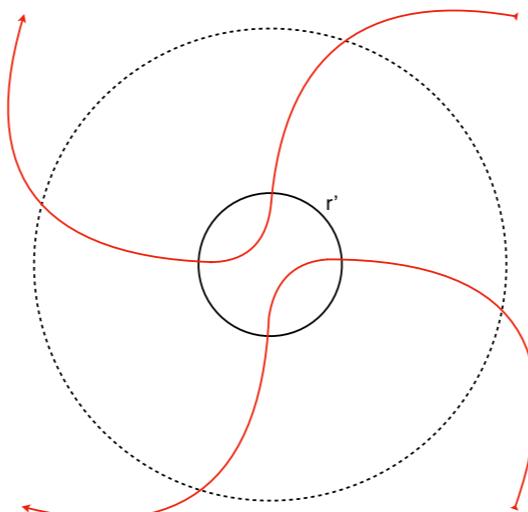
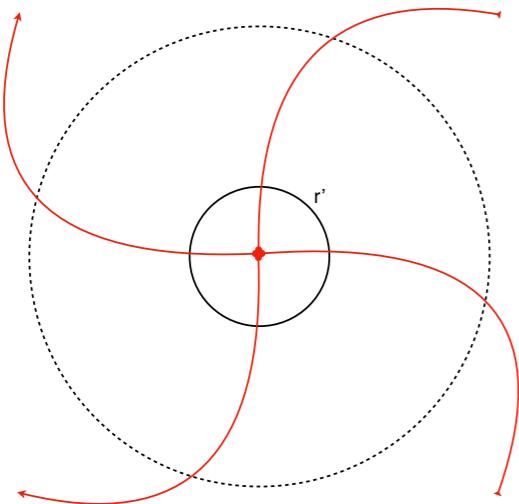
Solution:

$$n(r, \gamma) = \gamma^{-1-\mu} \int \int \frac{r_0^{\frac{3}{2}} H[r - r_0] \cdot H[\gamma_0 - \gamma]}{\sqrt{4\pi \cdot D_0 \cdot b_0 \cdot r_c^2}} \cdot \frac{\gamma_0^{\mu+2\nu}}{\sqrt{\frac{\gamma^\nu - \gamma_0^\nu}{\nu}}} \cdot \exp \left(-\frac{D_0}{4\nu \cdot b_0 \cdot r_c^2} (\gamma^\nu - \gamma_0^\nu) \right) \cdot \\ \left[\exp \left(-\frac{\log(\sqrt{r_0})^2 \nu \cdot b_0 \cdot r_c^2}{D_0 \cdot (\gamma^\nu - \gamma_0^\nu)} \right) - \exp \left(-\frac{\log(\frac{r}{\sqrt{r_0}})^2 \nu \cdot b_0 \cdot r_c^2}{D_0 \cdot (\gamma^\nu - \gamma_0^\nu)} \right) \right] \cdot Q(r_0, \gamma_0) dr_0 d\gamma_0$$

Backup slides



Backup slides



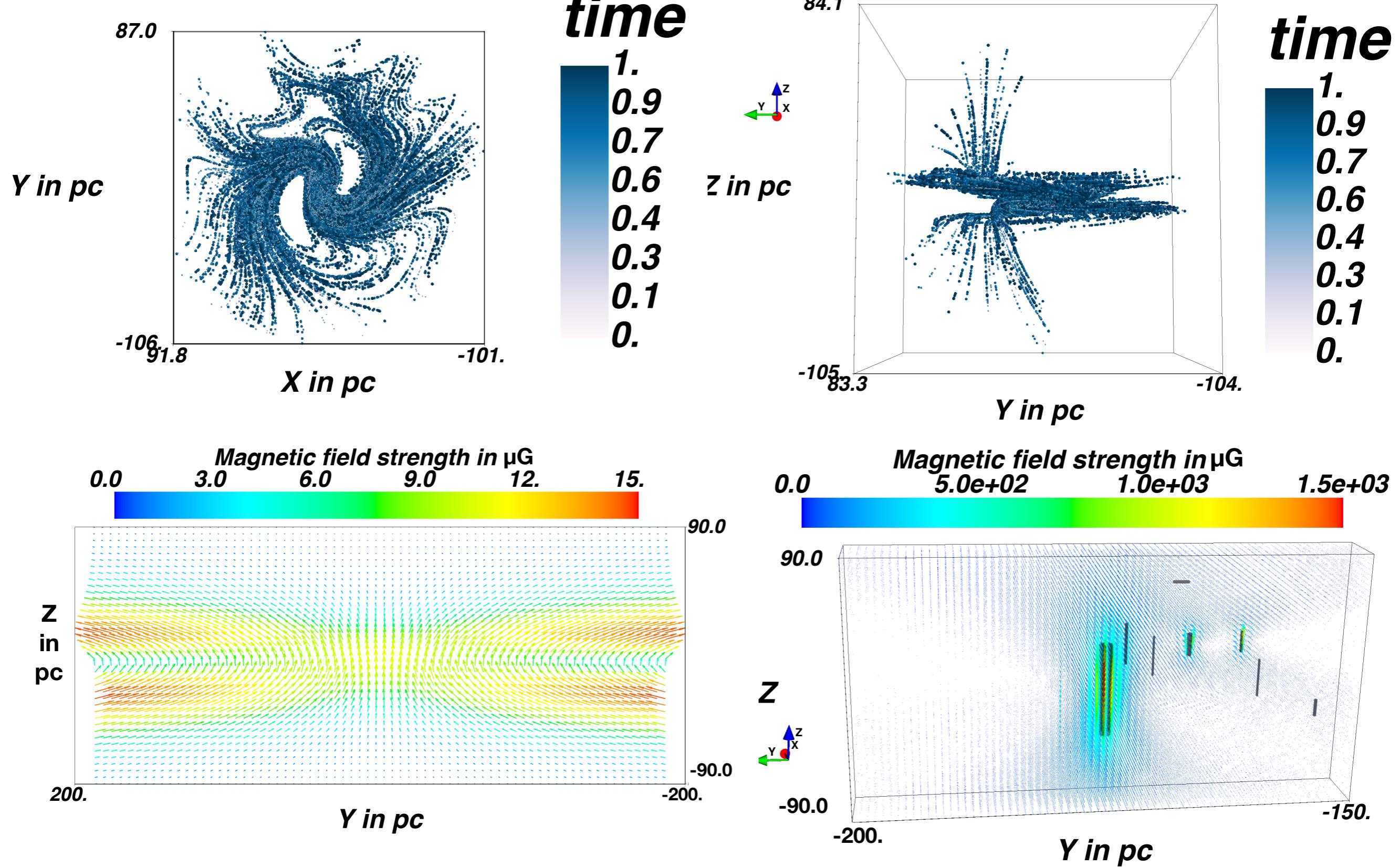
$$\vec{B}_{\text{MC}}^{r>r'} = \beta \cdot \begin{pmatrix} \frac{R}{r} \\ \mp \eta^{-1} \frac{R}{r+b} \\ 0 \end{pmatrix}$$

$$\beta = B_1 \cdot \cos(\pm v(r) + m \cdot \phi) \cdot \exp\left(-\frac{z^2}{H_c^2}\right)$$
$$v(r) = \frac{1}{\eta} \cdot \ln\left(\frac{r+b}{R+b}\right)$$

$$\vec{B}_{\text{MC}}^{r<r'} = \beta \cdot \frac{R}{r'} \left(\frac{3r}{r'} - \frac{2r^2}{r'^2} \right) \begin{pmatrix} 1 \\ \mp \frac{r}{\eta(r+b)} \left(1 + \frac{6(r-r')}{2r-3r'} \left(\frac{\sin(\pm v(r)+m\phi)-\sin(\pm v(r))}{\cos(\pm v(r)+m\phi)} \right) \right) \\ 0 \end{pmatrix}$$

$$B_{\text{tot}} = B_{\text{IC}}^C + \sum_{i=1}^8 B_{\text{NTF},i}^C + \sum_{i=1}^{12} B_{\text{MC},i}$$

Backup slides



Backup slides

