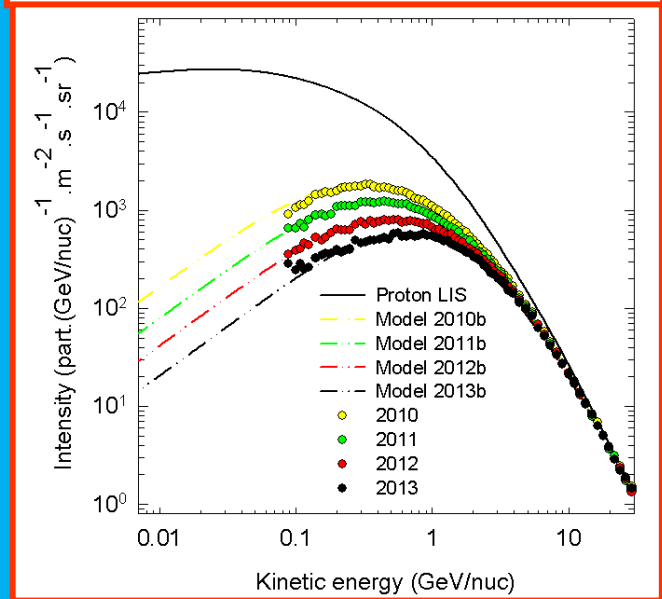
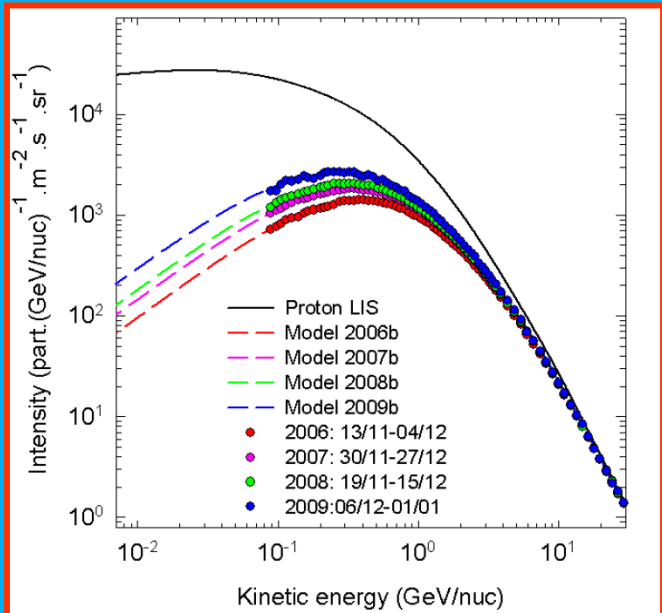


Effects of scattering parameters on charge-sign dependent cosmic ray modulation

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Introduction



- We apply the Measurement-Validated 3D numerical model (M-V Model) of Potgieter & Vos (2017) based on Parker's transport equation (TPE).

$$K_T = K_{A0} \frac{\beta P}{3B_m} \frac{(P/P_{A0})^2}{1 + (P/P_{A0})^2} \rightarrow K_T = \frac{\beta P}{3B_m} \left[\frac{(\omega\tau)^2}{1 + (\omega\tau)^2} \right]$$

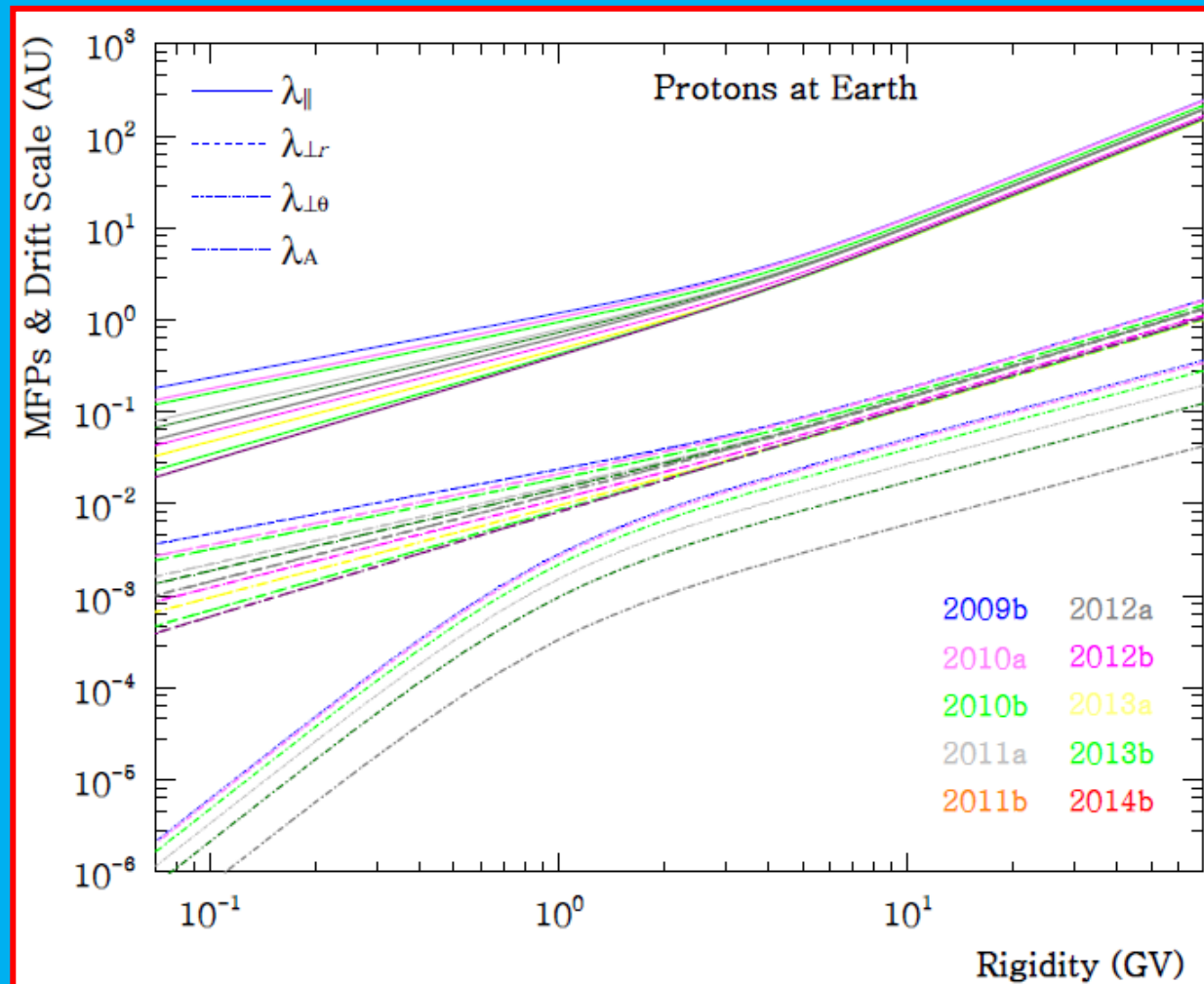
- To make meaningful comparison the modulation of both protons and anti-protons is done using the same set of modulation parameters and diffusion coefficients.
- We illustrate and discuss differences that exist between protons and anti-protons in their intensity-time profiles and ratios (p/pbar) from 2006 till the end of 2014 due to the assumed spatial dependence of $\omega\tau$.

Corresponding modulation parameters from 2009 to 2013

Table: Summary of the modulation parameters used to reproduce the proton measurements 2009b-2013b from PAMELA.

Parameters	2009b	2010b	2011b	2012b	2013b
α (degree)	9.50	21.01	45.21	62.09	65.78
B (nT)	3.91	4.52	5.11	5.58	5.19
τ_z (AU)	80.0	82.0	84.0	86.0	88.0
$\lambda_{ }$ (AU)	1.185	1.090	0.948	0.798	0.565
K_{A0}	0.90	0.80	0.40	0.0	0.0
P_{A0} (GV)	0.90	0.90	0.90	0.90	0.90
$K_{\perp r}^0$	0.02	0.02	0.02	0.02	0.02
$K_{\perp \theta}^0$	0.02	0.02	0.02	0.02	0.02
c_1	0.70	0.77	0.89	0.96	1.09
$c_{2 }$	1.52	1.52	1.52	1.52	1.52
$c_{2\perp}$	1.14	1.14	1.14	1.14	1.14
c_3	2.50	2.50	2.50	2.50	2.50
P_k (GV)	4.00	4.00	4.00	4.00	4.00
$d_{\perp \theta}$	6.00	6.00	6.00	6.00	6.00

Corresponding Mean Free Paths and Drift Scale



Transport equation for the modulation of cosmic rays in the heliosphere

Parker's (1965) Transport Equation (TPE):

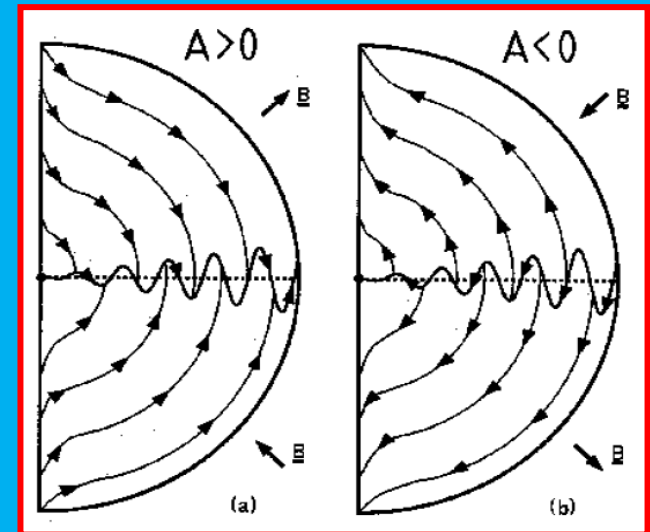
$$\frac{\partial f}{\partial t} = \underbrace{\nabla \cdot [\mathbf{K} \cdot \nabla f]}_{\text{Diffusion}} - \underbrace{\mathbf{V} \cdot \nabla f}_{\text{Convection}} - \underbrace{\langle \mathbf{v}_D \rangle \cdot \nabla f}_{\text{Particle drifts}} + \underbrace{\frac{1}{3}(\nabla \cdot \mathbf{V}) \frac{\partial f}{\partial \ln p}}_{\text{Adiabatic energy changes}} + \underbrace{Q(r, p, t)}_{\text{sources}} \quad (1)$$

$f(\mathbf{r}, p, t)$ is the cosmic ray distribution function.

\mathbf{K} is the diffusion tensor:

$$\mathbf{K} = \begin{bmatrix} K_{rr} & K_{r\theta} & K_{r\phi} \\ K_{\theta r} & K_{\theta\theta} & K_{\theta\phi} \\ K_{\phi r} & K_{\phi\theta} & K_{\phi\phi} \end{bmatrix}$$

$\mathbf{V}(r, \theta) = V(r, \theta)\mathbf{e}_r$ is the solar wind velocity vector



\mathbf{v}_D is the averaged gradient and curvature drift velocity

The drift coefficient

In general: $\langle \mathbf{v}_d \rangle = \nabla \times K_T \frac{\mathbf{B}}{B_m}$ (1)

$K_T = \frac{\beta P}{3B_m} f_s$, where $f_s = \frac{(\omega\tau)^2}{1 + (\omega\tau)^2}$ with ω the gyro-frequency and τ some time scale defined by scattering.
 Bieber&Matthaeus, 1997

Re-writing Equation (1):

$$\langle \mathbf{v}_d \rangle = \frac{\beta P}{3} \left[f_s \nabla \times \frac{\mathbf{B}}{B_m^2} + \nabla f_s \times \frac{\mathbf{B}}{B_m^2} \right] \quad (2)$$

WS drifts modification can be accomplished through assuming:

CASE 1: $\omega\tau = \text{constant}$ (similar to Potgieter et al., 1989).

CASE 2: $\omega\tau \sim \text{constant} \times P$ (see Burger et al., 2000).

CASE 3: $\omega\tau$ has spatial dependence (see Bieber & Matthaeus, 1997; Burger & Visser, 2010; Engelbrecht & Burger, 2015; Ngobeni & Potgieter, 2015).

CASE 2: The rigidity dependence

e.g. Ferreira, 2002; Potgieter & Vos, 2017 and etc

$$K_T = k_A \frac{\beta P}{3B_m} \frac{(P/P'_0)^2}{1 + (P/P'_0)^2}$$

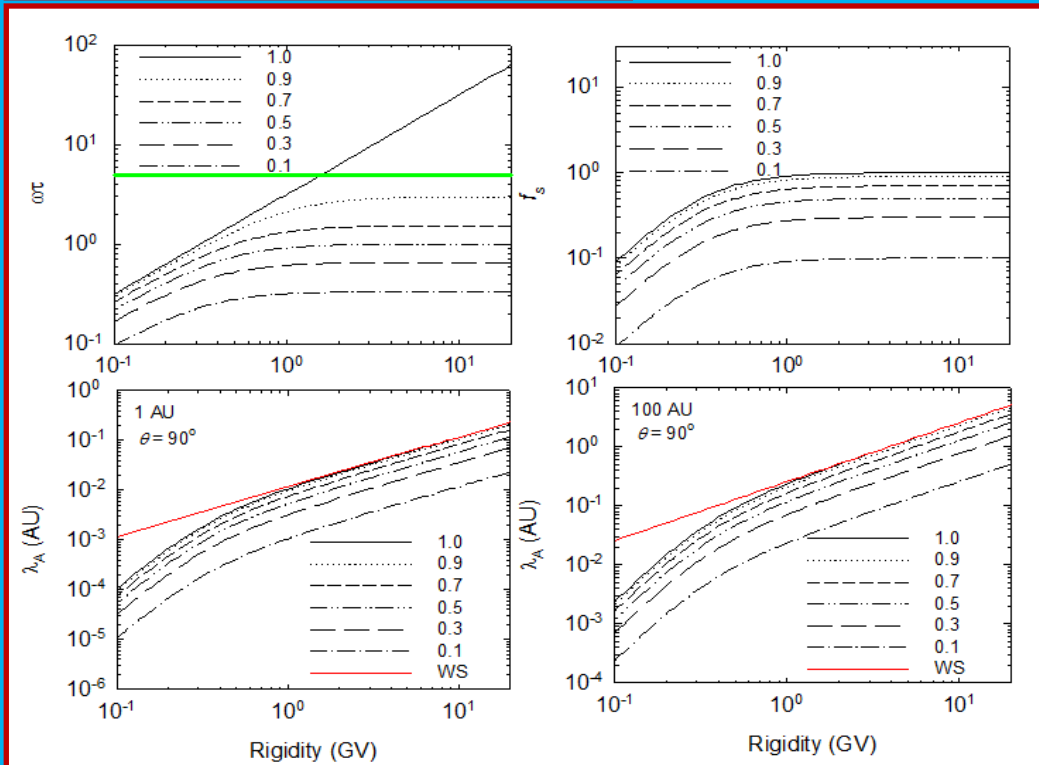


$$f_s = k_A \frac{(P/P'_0)^2}{1 + (P/P'_0)^2}$$

then

$$\omega\tau = \sqrt{k_A \frac{(P/P'_0)^2}{1 + (1 - k_A) (P/P'_0)^2}}$$

for $k_A = 1.0$, $\omega\tau$ reduces to P/P'_0 (assumption made by Burger et al. 2000)



Three important points

- ❖ When $k_{A0} < 1.0$ drifts are also reduced at $P > 1.0$ GV.
- ❖ For any value $\omega\tau > 5$, f_s remains ~ 1.0 (indicating no substantial drift reduction).
- ❖ WS drift scale ($\lambda_A = 3K_T/v$) is reduced by factor k_{A0} at $P > 1.0$ GV.

Extracting the spatial dependence of f_s

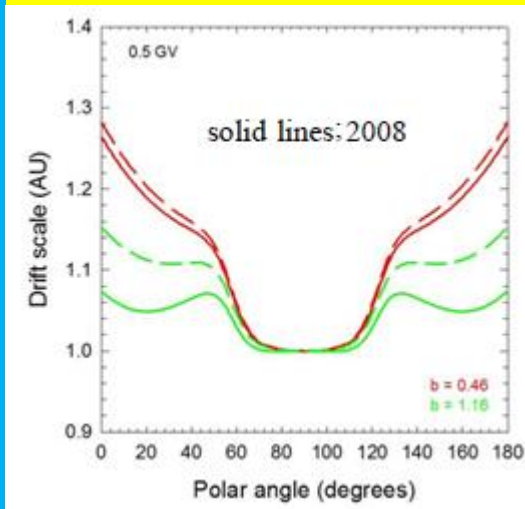
The analytical expression of the drift coefficient from Tautz and Shalchi is given as:

$$K_T = \frac{\beta P}{3B_m} \frac{1}{1 + a \left[\frac{\delta B}{B_m} \right]^{2b}} \rightarrow f_s = \frac{1}{1 + a \left[\frac{\delta B}{B_m} \right]^{2b}} = \frac{(\omega\tau)^2}{1 + (\omega\tau)^2} \rightarrow \omega\tau = \frac{1}{\sqrt{a} \left[\frac{\delta B}{B_m} \right]^b}$$

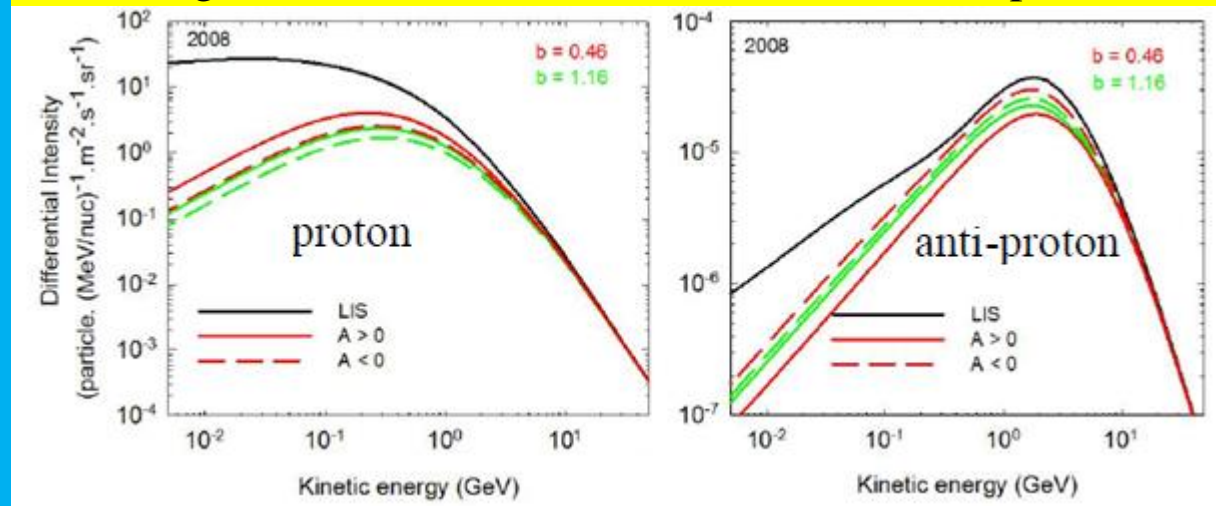
$$= \frac{\beta P}{3B_m} f_s$$

Best fit was achieved with $a = 1.09 \pm 0.52$ and $b = 0.81 \pm 0.35$

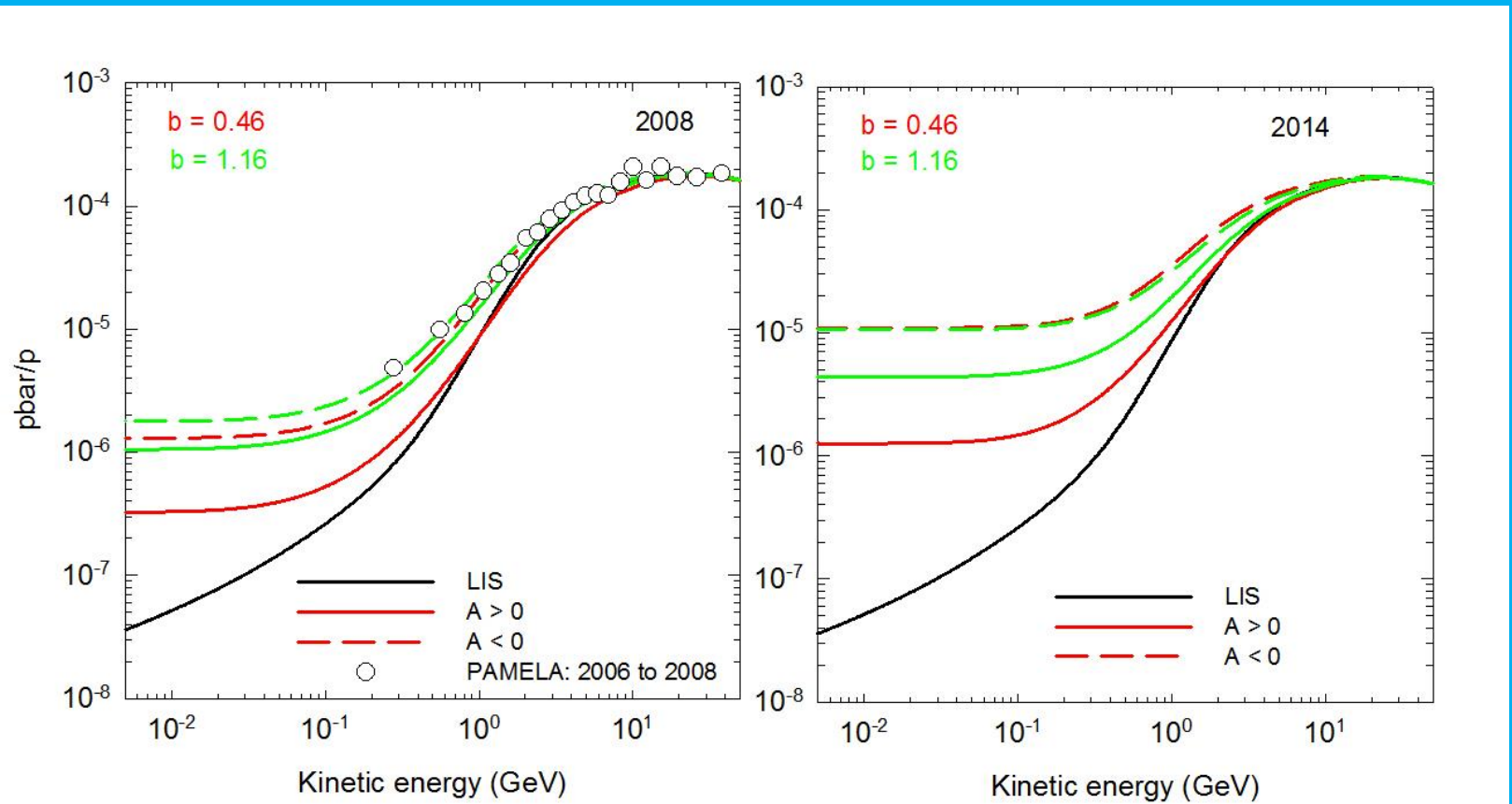
Normalized drift scale



Modeling results at the Earth: $A > 0$ and $A < 0$ comparison

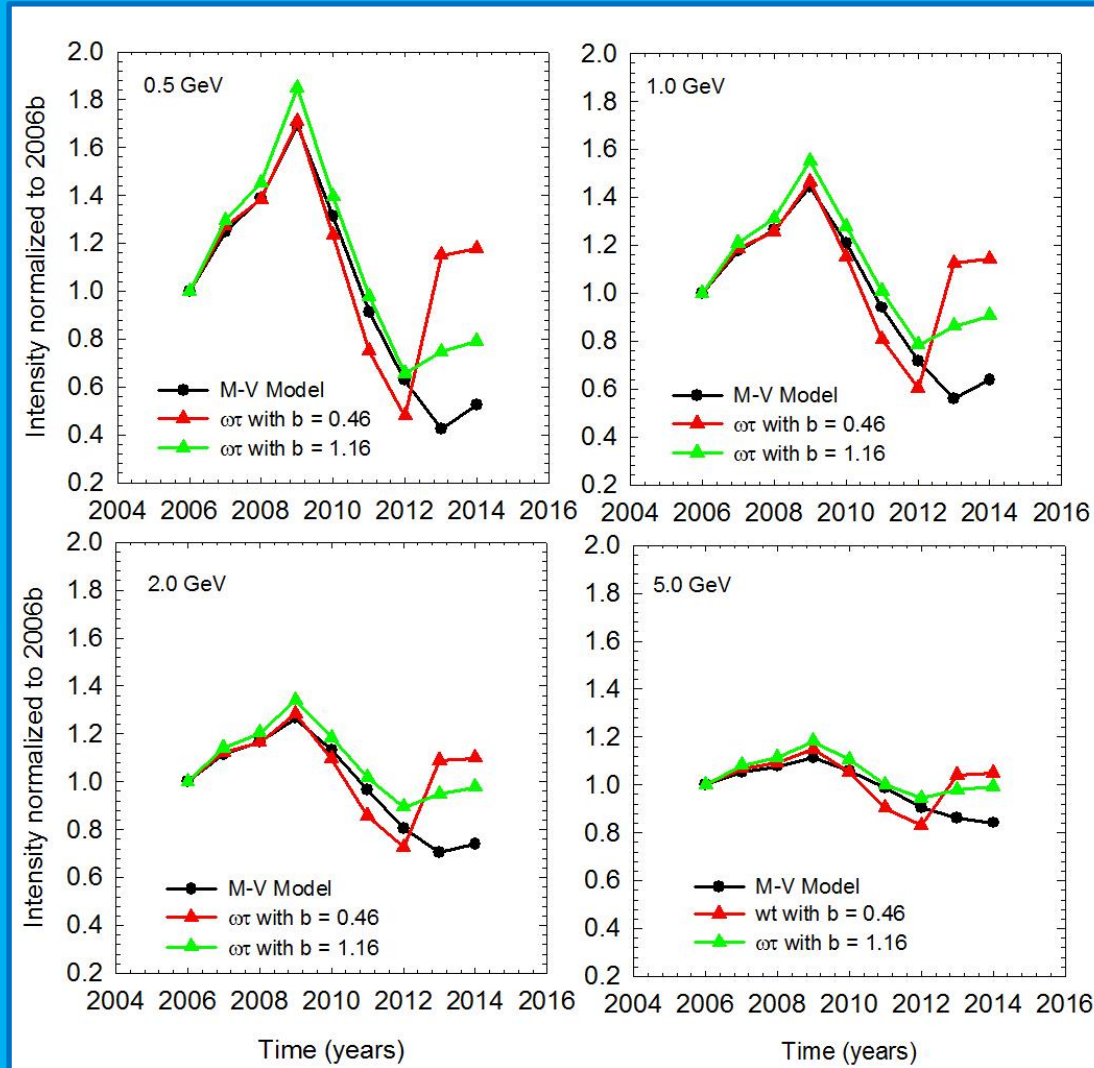


Energy dependence of anti-protons/protons ratios



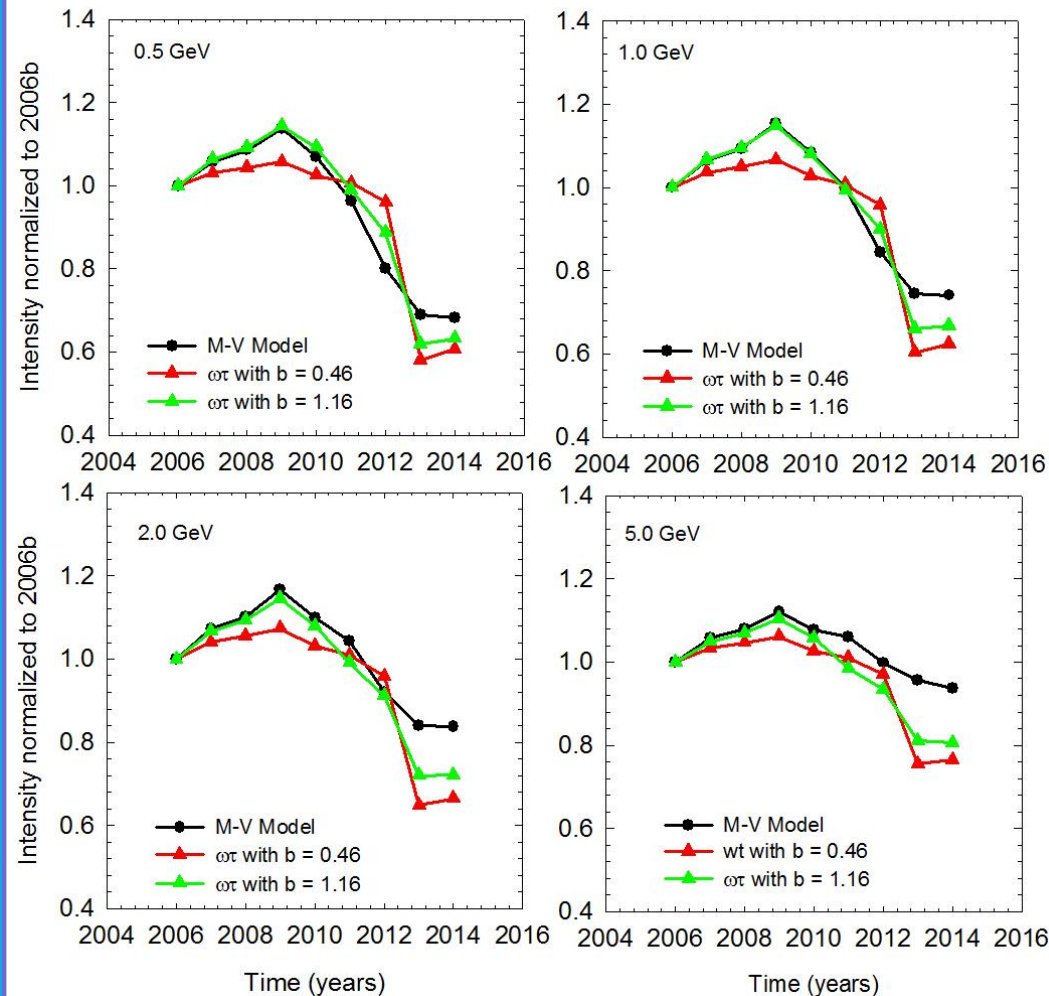
Effects of spatial dependence of $\omega\tau$ are more prominent in the $A > 0$ cycle

Comparison between M-V Model and spatial dependence of $\omega\tau$: protons



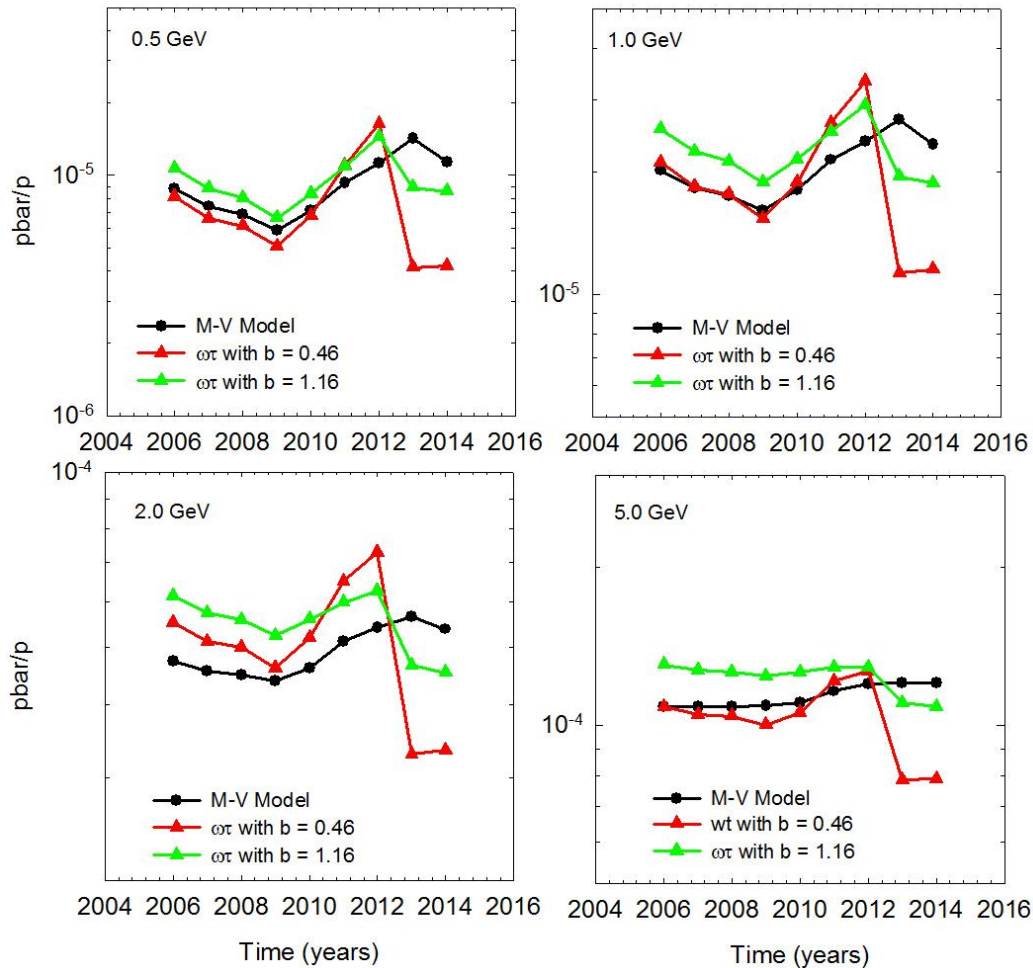
$\omega\tau$ with $b = 0.46$ describes the M-V Model well from 2006 to 2010. While $b = 1.16$ gives better estimation of MV-Model between 2010 and 2012

Comparison between M-V Model and spatial dependence of $\omega\tau$: anti-protons



$\omega\tau$ with $b = 1.16$ describes the M-V Model well from 2006 to 2012 below 5.0 GeV.

Modelling anti-proton to proton intensity ratios over time



- ✓ The proton intensity increased relatively more than anti-proton intensity until 2009 (corresponding to decreasing tilt angle).
- ✓ After 2009 (until 2012) proton intensity decreased relatively more (corresponding to increasing tilt angle).
- ✓ The large decreases in the ratio after 2012 is due to inadequate particle drift reduction and deviate largely with M-V Model.

Summary and Conclusions

- ❑ Using the self-consistent measurement validated 3D numerical model of Potgieter & Vos (2017) that includes particle drifts, the modulation of both protons and anti-protons was studied from 2006 to 2014 using.
- ✓ PAMELA proton observations together with numerical modeling confirmed that drifts played a significant role in modulation of GCRs from 2006 until around 2012.
- ✓ The intensity-time profile of protons from 2006 to 2012 can be qualitatively described with the assumed spatial dependence of $\omega\tau$, on the drift coefficient, by adjusting values of a and b .
- ✓ Both assumptions made about $\omega\tau$ are inadequate, and thus unsuitable, to describe $p\text{bar}/p$ ratios between 2012 and 2014, as required by the MV-Model.