

# **Comparing Long-Duration Gamma-ray Flares and High-Energy Solar Energetic Particles**

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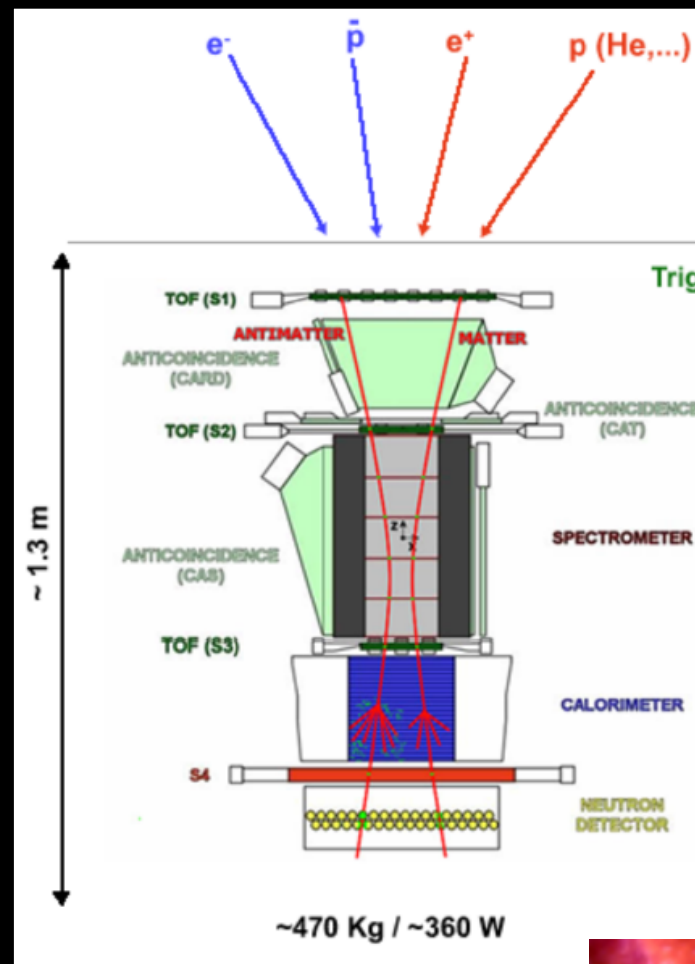
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**On behalf of the PAMELA Collaboration**

# Astrophysics Heliophysics



## Payload for Matter-Antimatter Exploration and Light Nuclei Astrophysics (PAMELA)

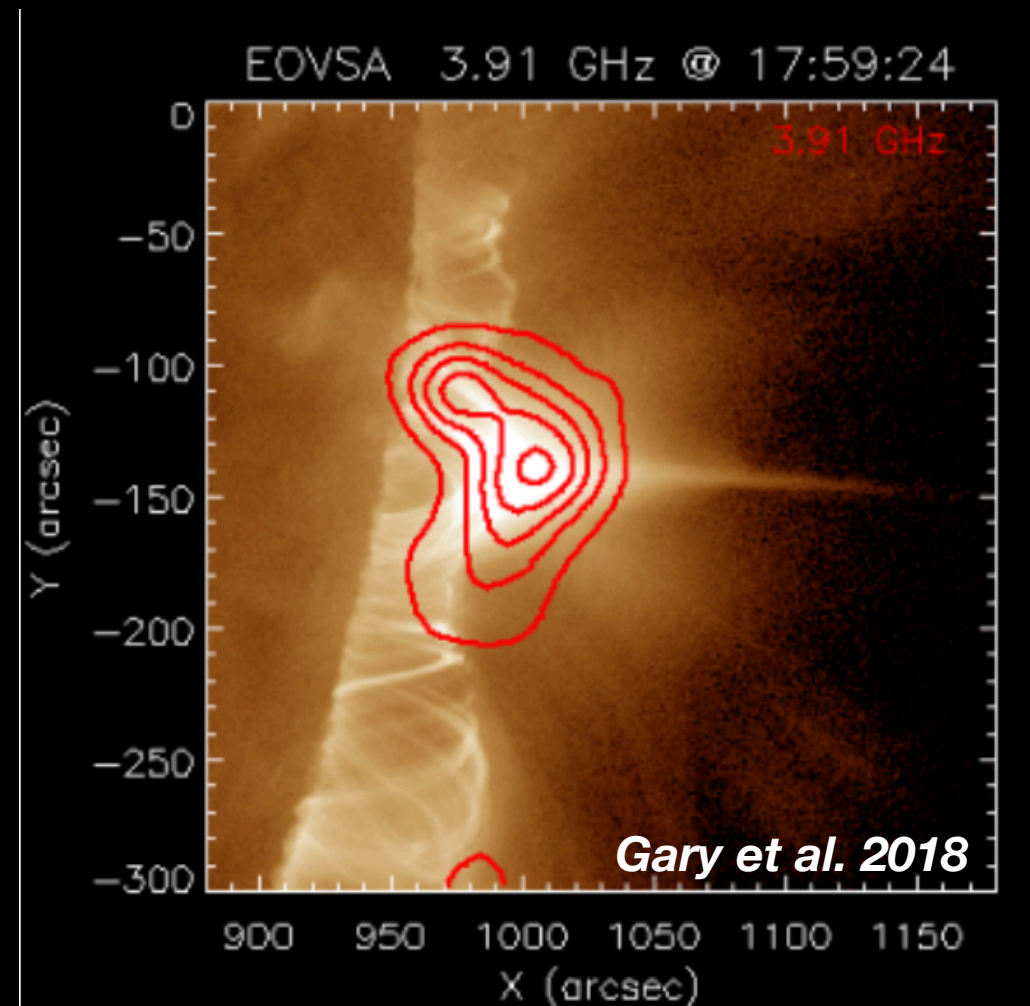
Magnetic spectrometer with silicon tracking system, a ToF, and EC to measure GCRs from tens of MeV up to hundreds GeV  
Also detect SEPs (see Bruno et al. 2018, this ICRC as well)

## Fermi/LAT

pair-conversion telescope with sensitivity to  $\gamma$ -rays between 20 MeV and 300 GeV & duty cycle for solar events of ~20%



G. A. de Nolfo



Ions producing the LDGRFs are in the SAME energy range as that observed by PAMELA!

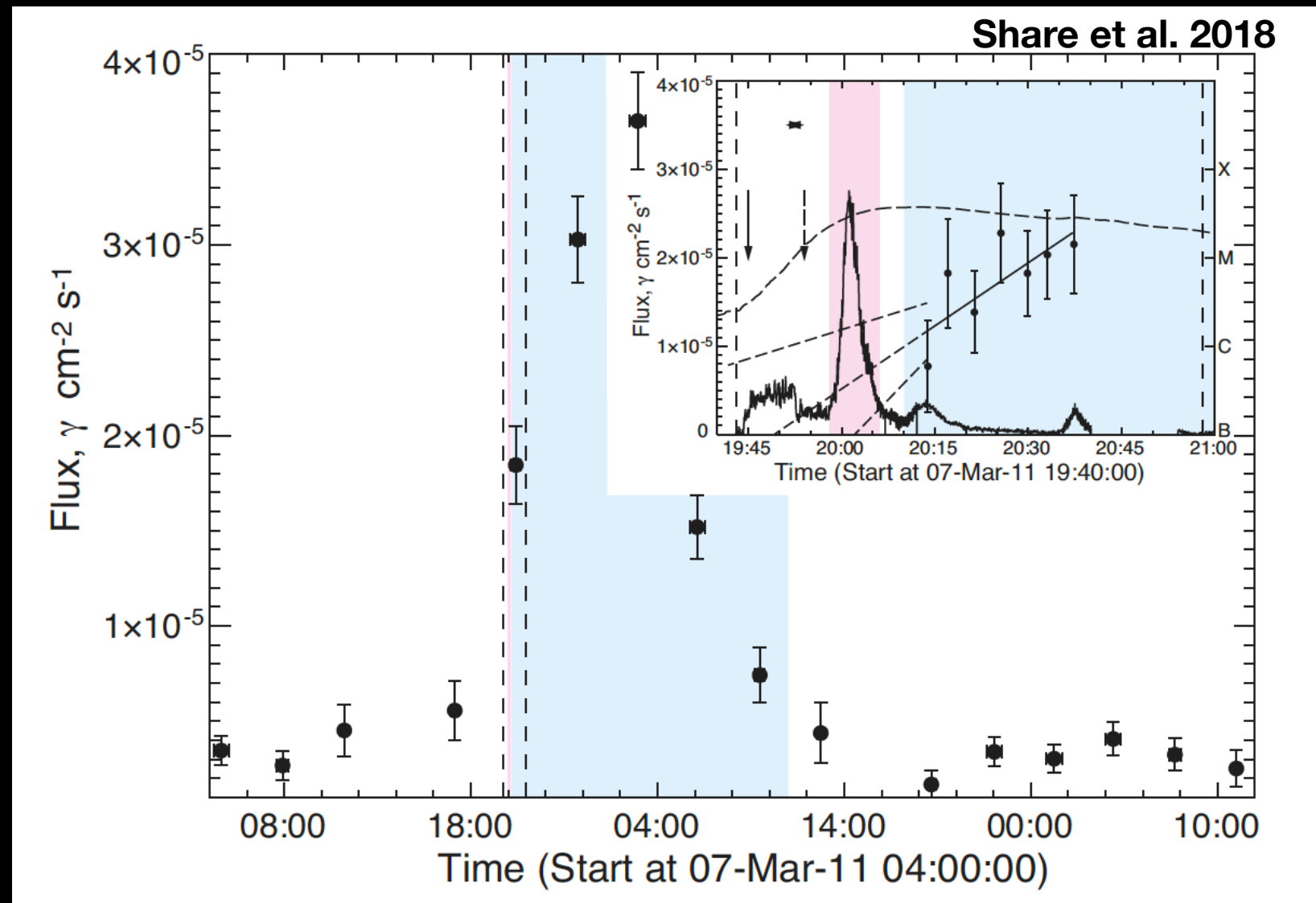
Possible to address the question of the origin of LDGRFs with PAMELA, STEREO, and Fermi/LAT for the first time!

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# Introduction : Long Duration Gamma-ray Flares

- \* Delayed and prolonged  $> 50$  MeV  $\gamma$ -ray emission well after the impulsive phase (durations up to 20 hrs!).
- \* Associated with  $> 100$  keV X-ray emission, CMEs, Type II & III radio emission (see *G. Share et al. 2018*)
- \* More (x10) fluence in delayed phase than impulsive phase



Time history of  $> 100$  MeV gamma-ray flux from Fermi/LAT. Inset compares with GBM 100-300 keV & dashed curve is soft x-rays

- \* High-energy gamma-ray emission ( $> 100$  MeV) is thought to originate primarily from the decay of pions, produced by protons (and alphas) above  $\sim 300$  MeV (above  $\sim 200$  MeV). See, for instance, *Ackermann et al. 2017*

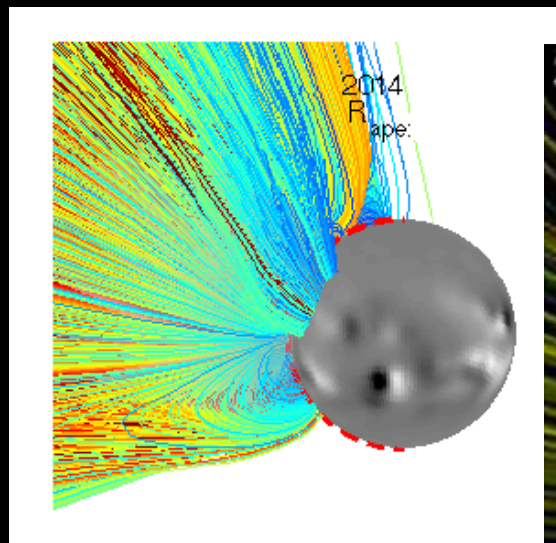
***The origin is still unknown & the challenge to theory is to explain the extreme energies & long durations !***

# Two Competing Theories

## **CME back-precipitation Scenario**

[Cliver et al. 1993; Kocharov et al. 2015]

- Attributed to CME-shock-accelerated protons that make their way back to the photosphere.
- Show that the reconstructed shock fronts become magnetically connected to visible solar surface just before onset of  $> 100$  MeV  $\gamma$ -ray emission



Plotnikov et al. 2017

Jin et al. 2018

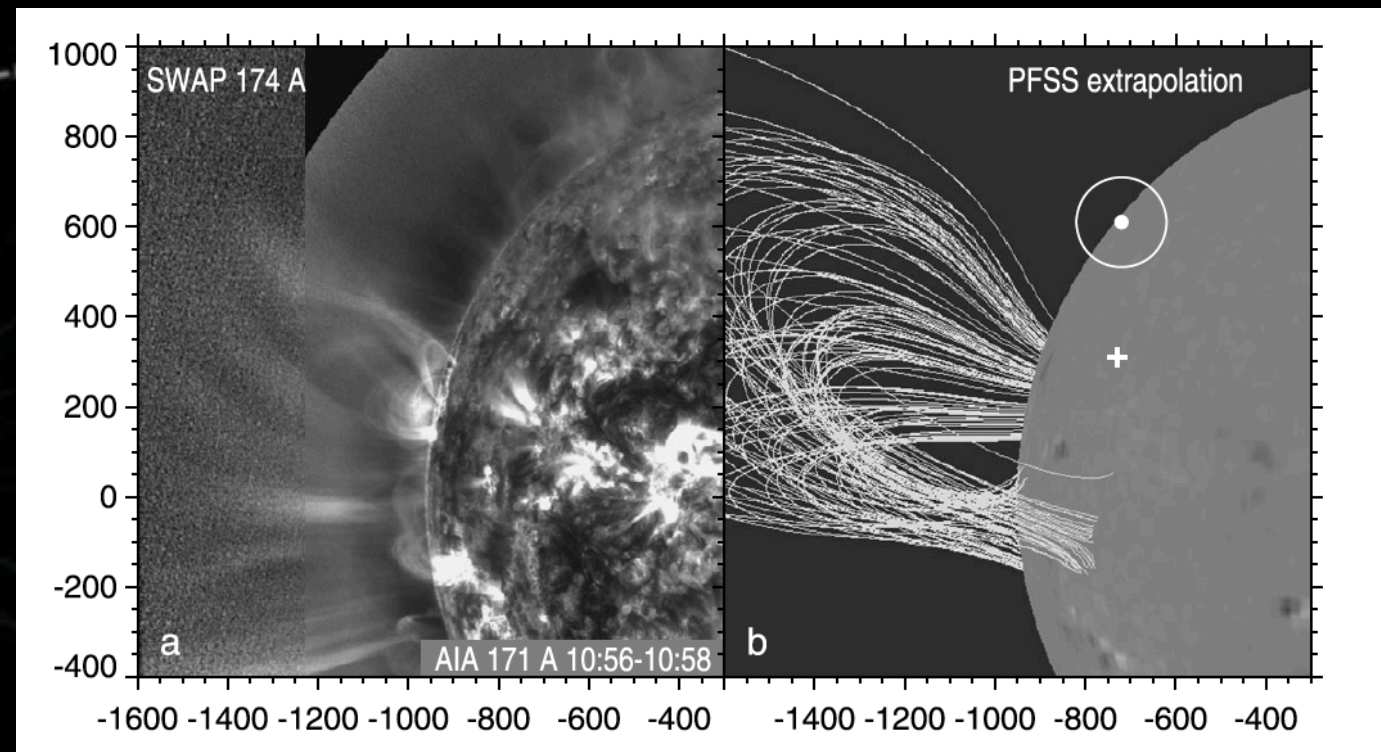
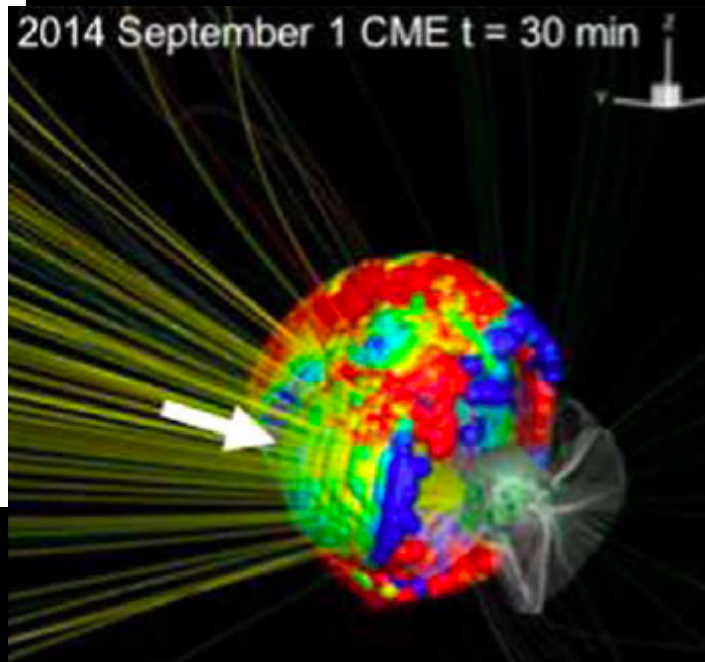
Compelling correlations between LDGRFs, CME speed, and Type II radio emission (Winter et al. 2018 ; Gopalswamy et al. 2018).

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## **Trapping & Acceleration in Large Coronal Loops**

[Chupp & Ryan 2009; Ryan & Lee 1991; Mandzhavidze & Ramaty 1992 ]

Consider injection and acceleration of particles along large coronal loops (precipitating in the photosphere) where pitch-angle scattering from magnetic turbulence may serve to further accelerate the particles.

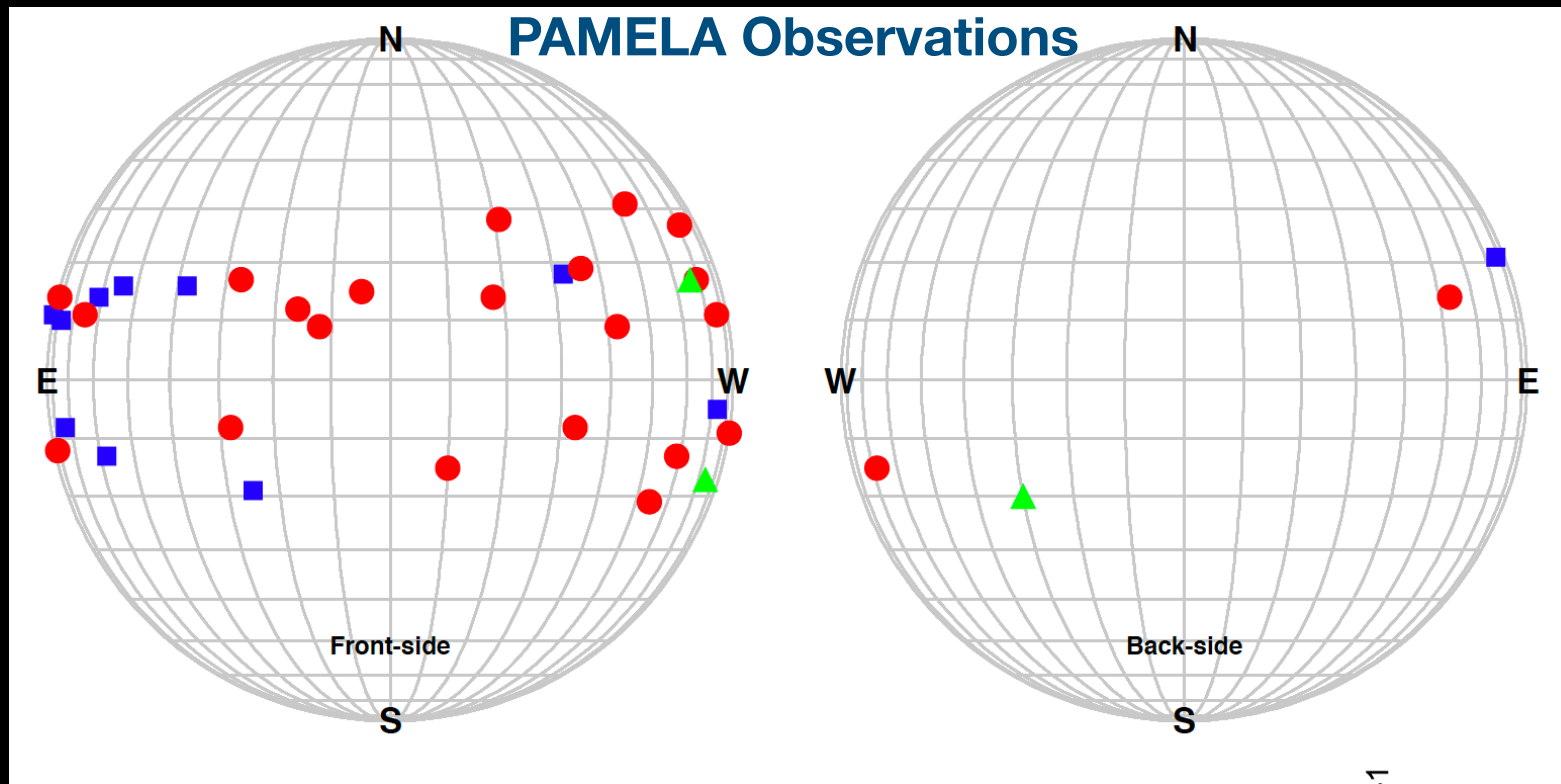


Grechnev et al. 2018, see also Ryan et al. this ICRC

⇒ **both models have supporting observations**

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⇒ Gain some insight into origin of LDGRFs by comparing with SEPs measured by PAMELA (either the populations are related or result from distinct processes)



Fermi/LAT (>100 MeV) and PAMELA (>500 MeV)

Red : Fermi & PAMELA

Blue : Fermi/LAT only (preponderance of eastern events)

Green PAMELA only (backside events and poor LAT coverage)

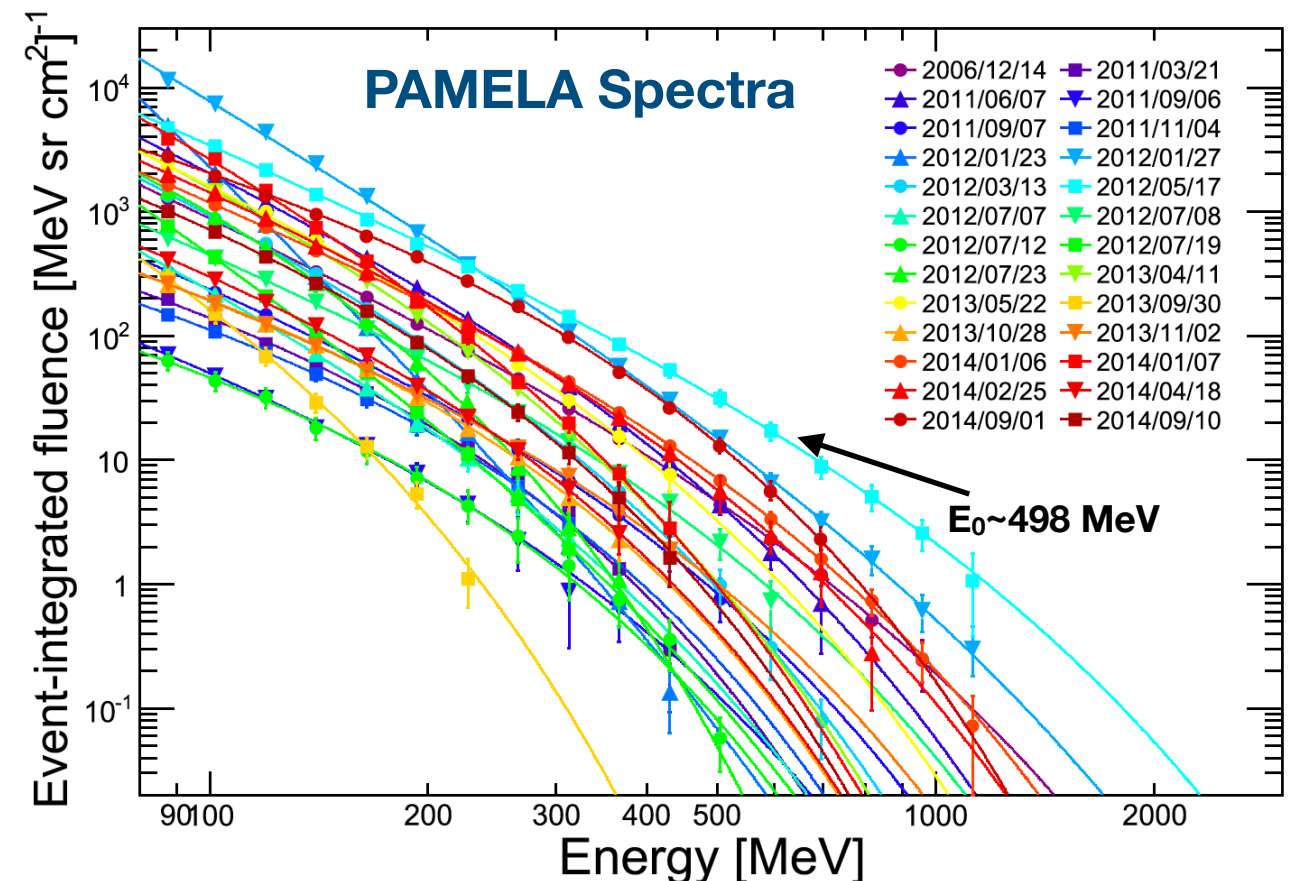
In summary, 18 out of the 25 SEP events observed by PAMELA were associated with LDGRFs by Fermi/LAT

Fluxes are consistent w/ the Ellison & Ramaty (1985) functional form consisting of a power-law with exponential cutoff

$$\Phi_{sep}(E) = A \times (E/E_s)^{-\gamma} \times e^{-E/E_0},$$

\* PAMELA measured spectra for 26 SEP events (see: Bruno, A. et al. (2018), ApJ 862:97 also, Bruno et al. this ICRC)

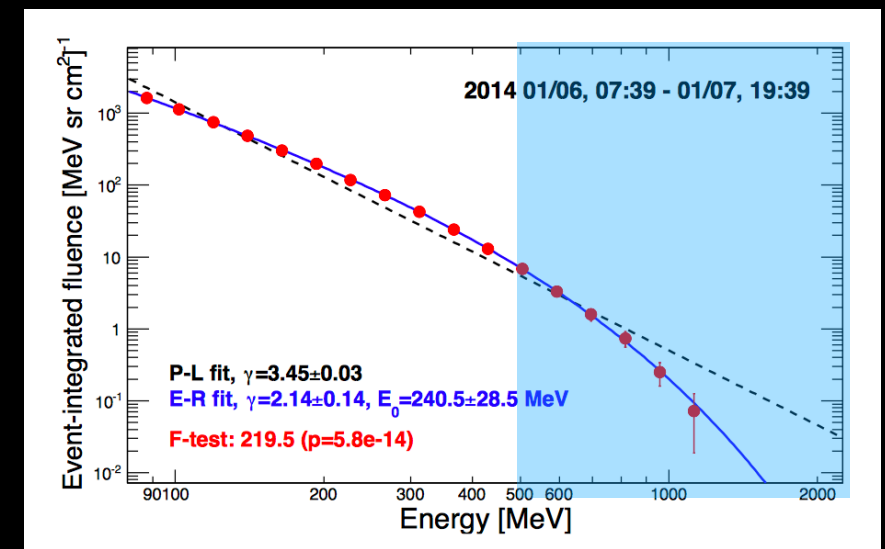
\* 14 SEP events were associated with LDGRF emission (see de Nolfo et al. 2019)





# Approach: Compare Total Proton Numbers at the Sun and in Space

1. Compute >500 MeV fluencies based of ER fits, accounting for spectral roll-overs



2. Compute number of protons assuming particle spatial distributions that is characterized by a periodic Gaussian ( $G(\delta)$ ) & integrate over a heliocentric spherical surface,  $S$ , at 1 AU.

$$N_{SEP} = \overline{N}_{cross}^{-1} \int_{4\pi} d\Omega \int_S dS (\mathbf{J} \cdot \mathbf{n}) \quad (1)$$

$$= \overline{N}_{cross}^{-1} \int_{4\pi} d\Omega \int_S dS \cos(\theta) J(\Omega, S),$$

$d\Omega = d\varphi d\vartheta \sin(\vartheta)$  is the solid angle element of the particle velocity direction at a point centered on the sphere and  $J$  is the event-integrated intensity

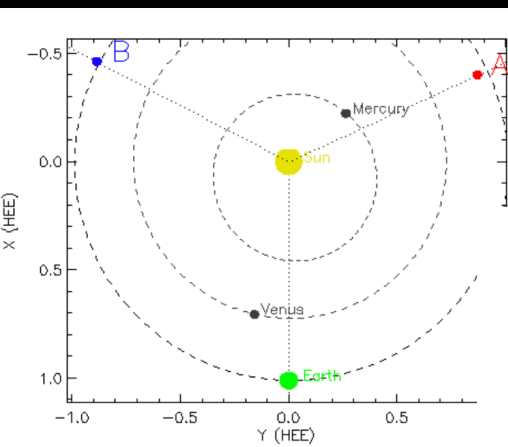
$$N_{SEP} = 2\pi \overline{N}_{cross}^{-1} J_{max} \int_S dS G(\delta) \quad (2)$$

$$= 2\pi \overline{N}_{cross}^{-1} J_E S_J C_{spa},$$

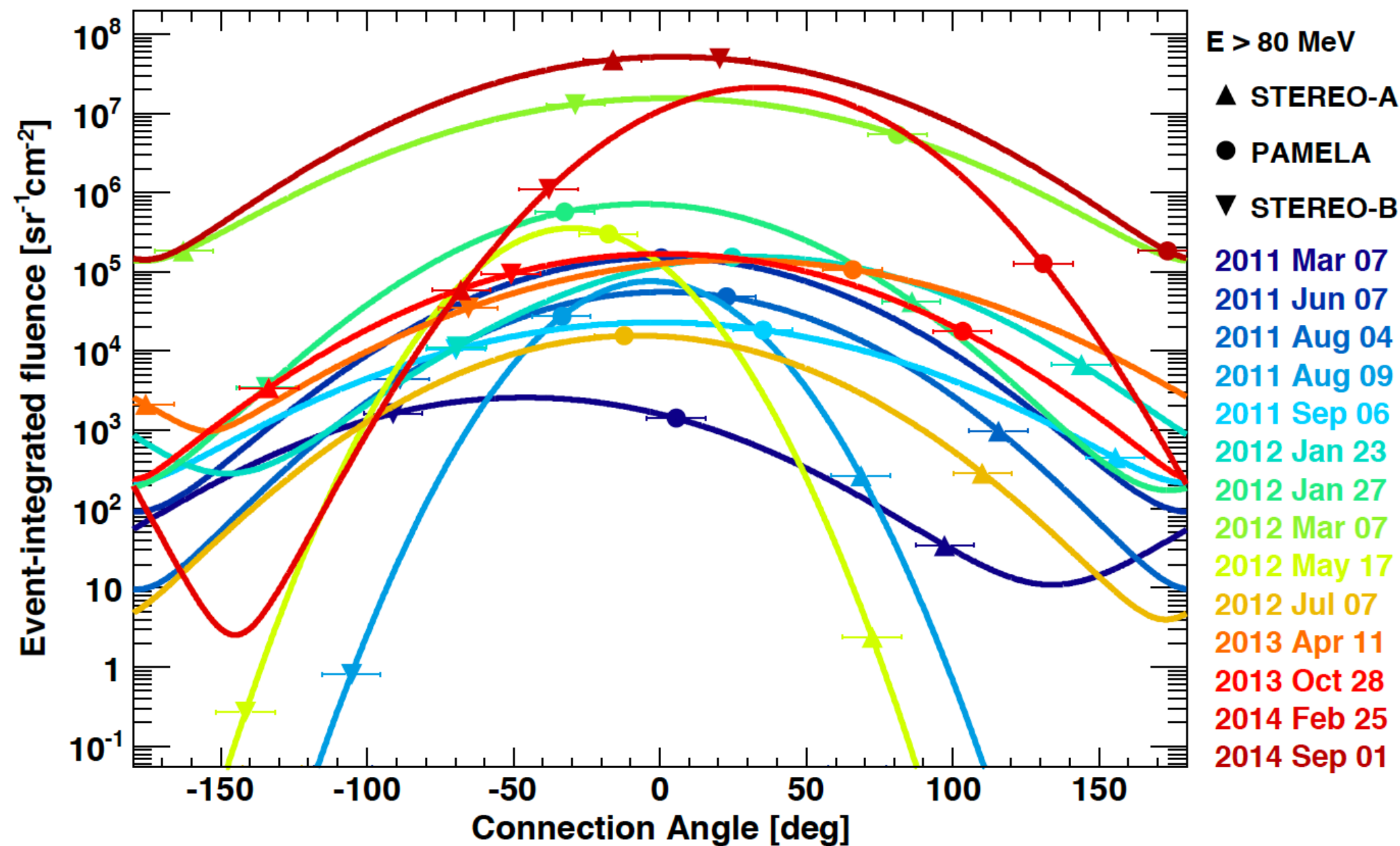
$J_E$  is the > 500 MeV event-event-integrated intensity observed by PAMELA  
 $S_J$  is the spherical area weighted by the particle spatial distribution  
 $\delta$  is the great-circle distance wrt the peak of the SEP partial distribution

Need to account for two important corrections:

1.  $C_{spa}$  accounts for PAMELA's observations not being made on interplanetary magnetic field lines that connect with the peak of the particle distribution,
2.  $N_{cross}$  takes into account multiple measurements of the same particles (beam vs. isotropic)



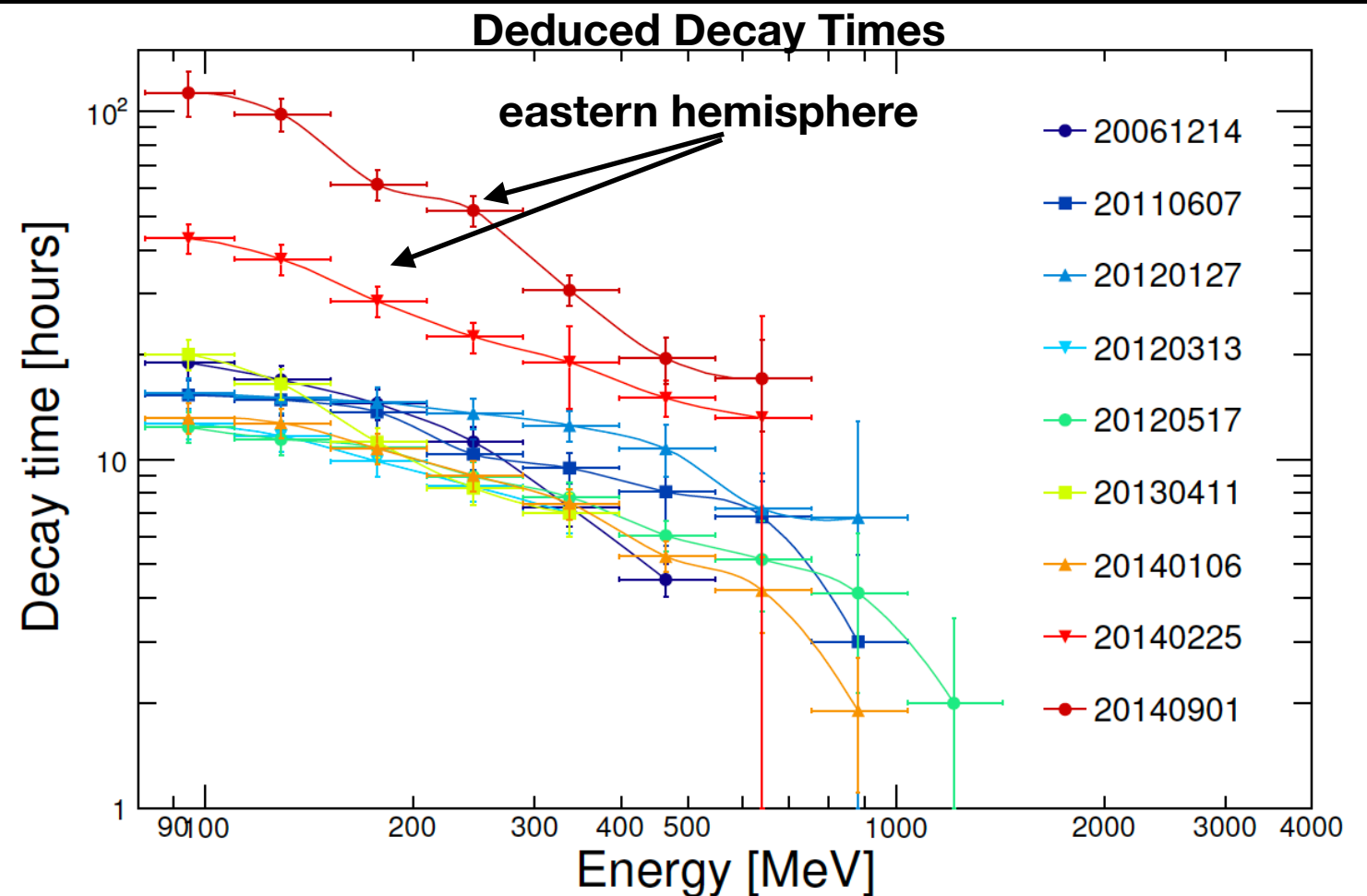
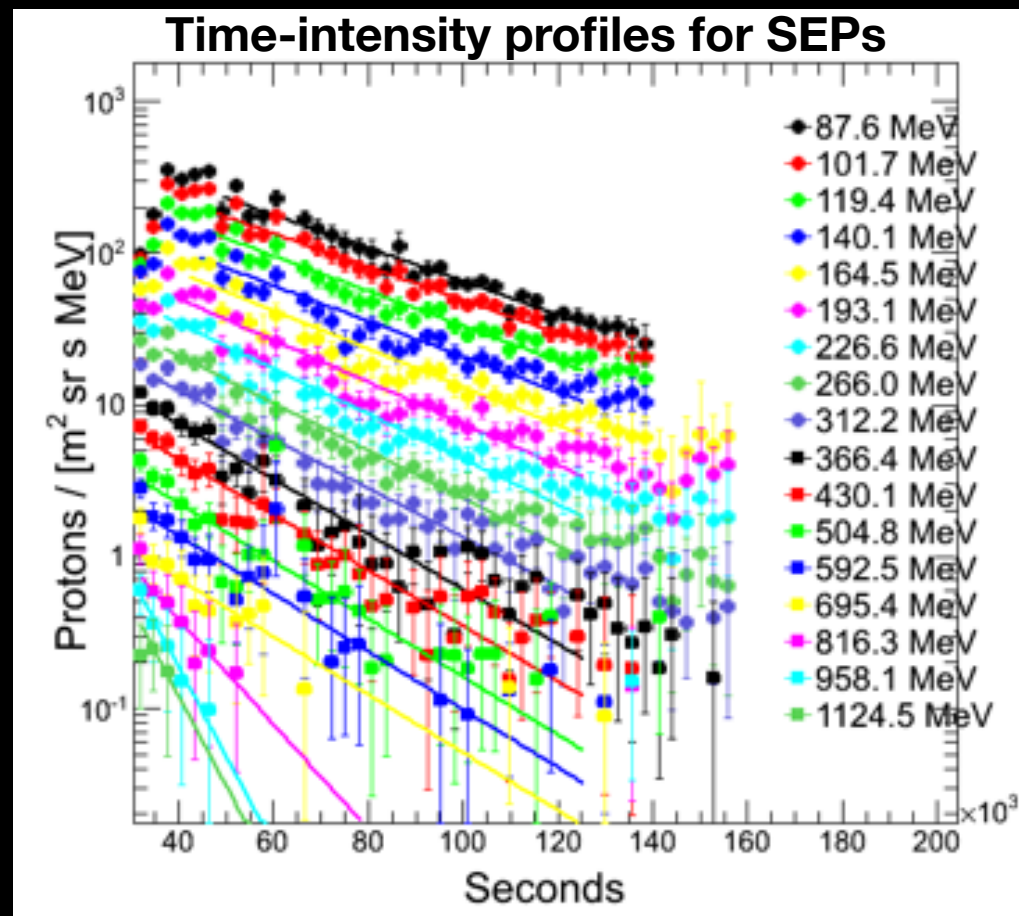
# Periodic Gaussian Fits at $> 80$ MeV Fluences



Longitudinal extent of SEP events determined from the fits of the event-integrated intensities ( $>80$  MeV) measured by PAMELA and STEREO A/B as a function of connection angle between the S/C magnetic footpoint at  $30 R_s$  & the location of the parent flare.

# PAMELA Observations also help to constrain transport

- \* SEP transport is governed by both large scale magnetic topology & scattering from small scale magnetic turbulence
- \* The amount of scattering affects the SEP intensity and anisotropy distributions



*Such trends are extremely helpful in constraining amount of scattering for SEPs*





# Modeling SEP Transport & Multiple Crossings

⇒ Depending on the amount of scattering, SEPs may cross 1 AU several times and this multiple scattering needs to be taken into account

—> Can determine  $N_{\text{cross}}$  through simulations of particle propagation under a variety of scattering conditions

Consider 2 test particle models

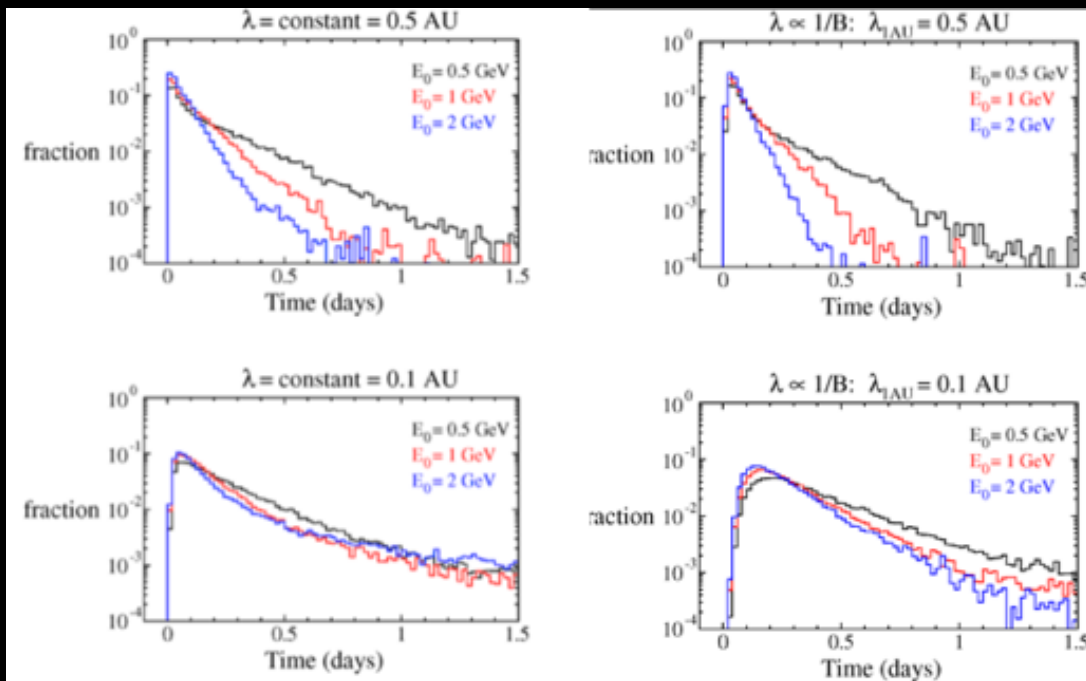
- 1) Simulation by Chollet et al. 2010
- 2) Simulation by Barttarbee et al. 2018

Both assume impulsive injection of mono-energetic isotropic particles at 0.1 AU, following the particles for 10 days, and both include magnetic focusing & scattering off of an unspecified plasma turbulence field.

# Modeling SEP Transport

For the Chollet et al. 2010 model, we assume two forms for the turbulence 1) uniform or 2) proportional to the gyro-cyclotron radius (Chollet et al. 2010)

Battarbee et al. 2018 model includes the effects of different configurations of the Heliospheric Current Sheet (HCS) and solar magnetic polarity.

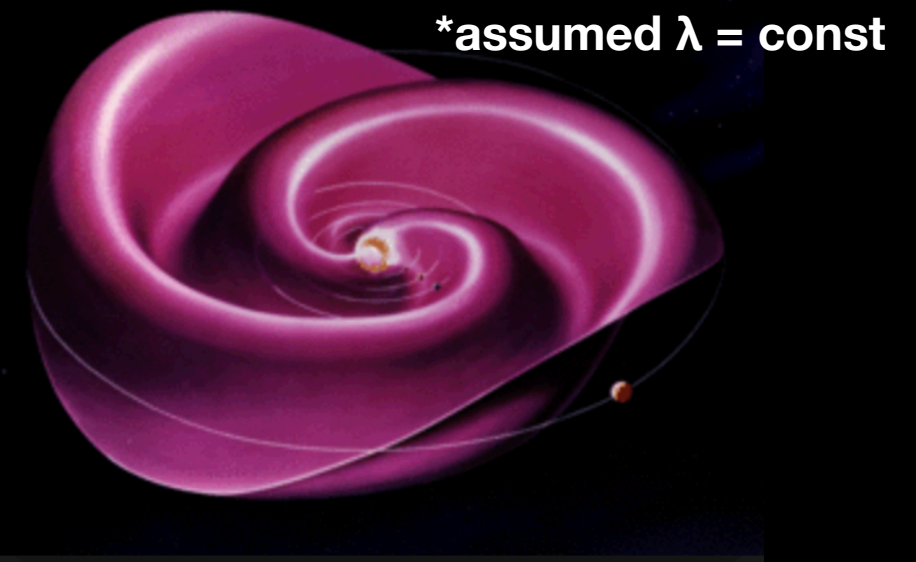


Predictions for the time-dependent development and decay of the intensity at 1 AU

The degree of scattering is adjusted to increase or decrease the anisotropy and associate decay time

\*assumed flat HCS and A+

$\Rightarrow$  short  $\lambda$   
results in  
longer decay  
times



These calculations show that  $N_{\text{cross}}$  varies for :

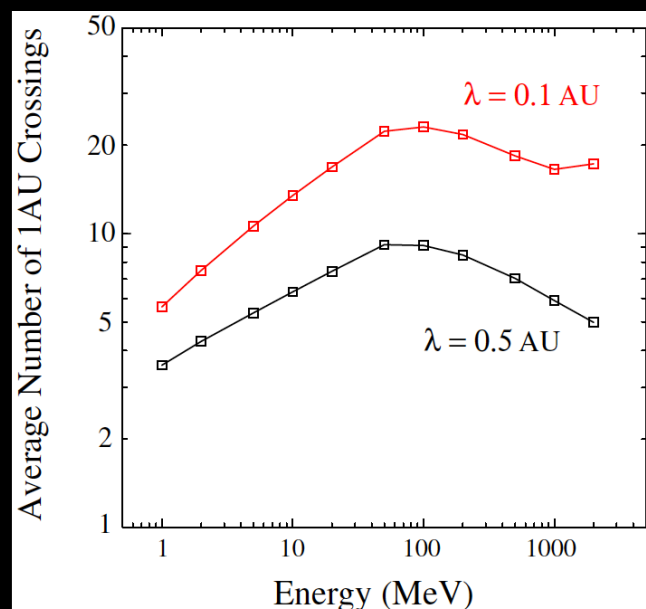
1) different configurations of the HCS (none, flat, or wavy).

2) magnetic polarity , A+ / A-

$\Rightarrow$  Large differences in  $N_{\text{cross}}$  for different polarities is due to particle drift along the HCS (e.g., A+ helps protons outward from the inner heliosphere faster)

$N_{\text{cross}}$  for flat HCS, A+ is consistent with results of Chollet et al. simulations for similar conditions.

Full simulation of 2012 May 17 is consistent with PAMELA for  $\lambda=0.3$  AU (Dalla et al., in prep)



# Computation of Total Proton Numbers

$$\begin{aligned} N_{SEP} &= 2\pi \overline{N}_{cross}^{-1} J_{max} \int_S dS G(\delta) \\ &= 2\pi \overline{N}_{cross}^{-1} J_E S_J C_{spa}, \end{aligned}$$

Compute number of protons assuming an isotropic flux & integrate over spherical surface at 1 AU.

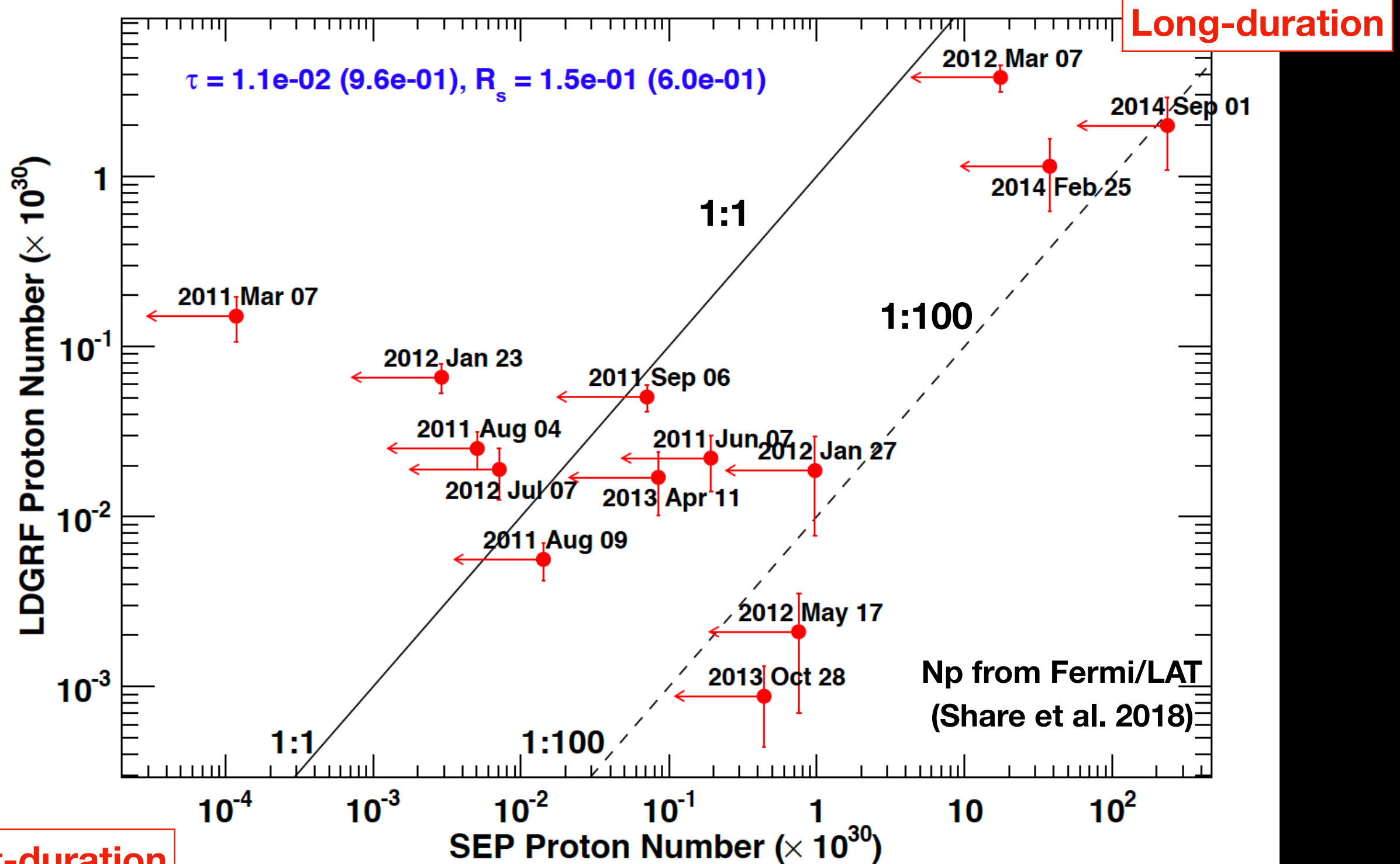
## Important Assumptions :

1. Use  $> 80$  MeV proton distributions to define longitudinal extent
2. Assume the same angular distribution for latitudinal dependence
3. Assume  $\lambda_0 \sim 0.5$  AU & wavy HCS  $\Rightarrow N_p \sim 8-11$

$\Rightarrow$  Compute **upper limits for  $N_p$**

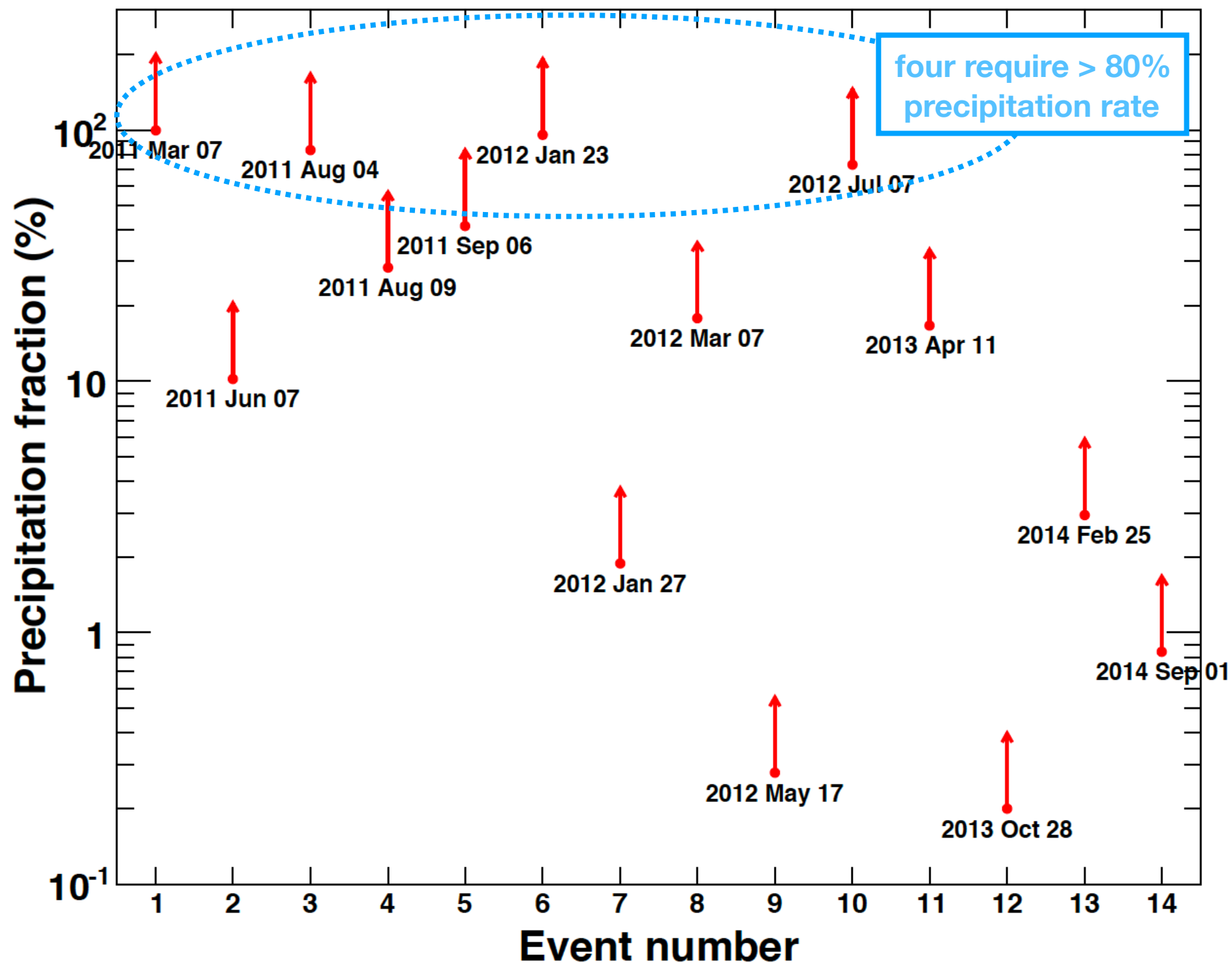


# Comparing $> 500$ MeV $N_p$ in space and at the Sun



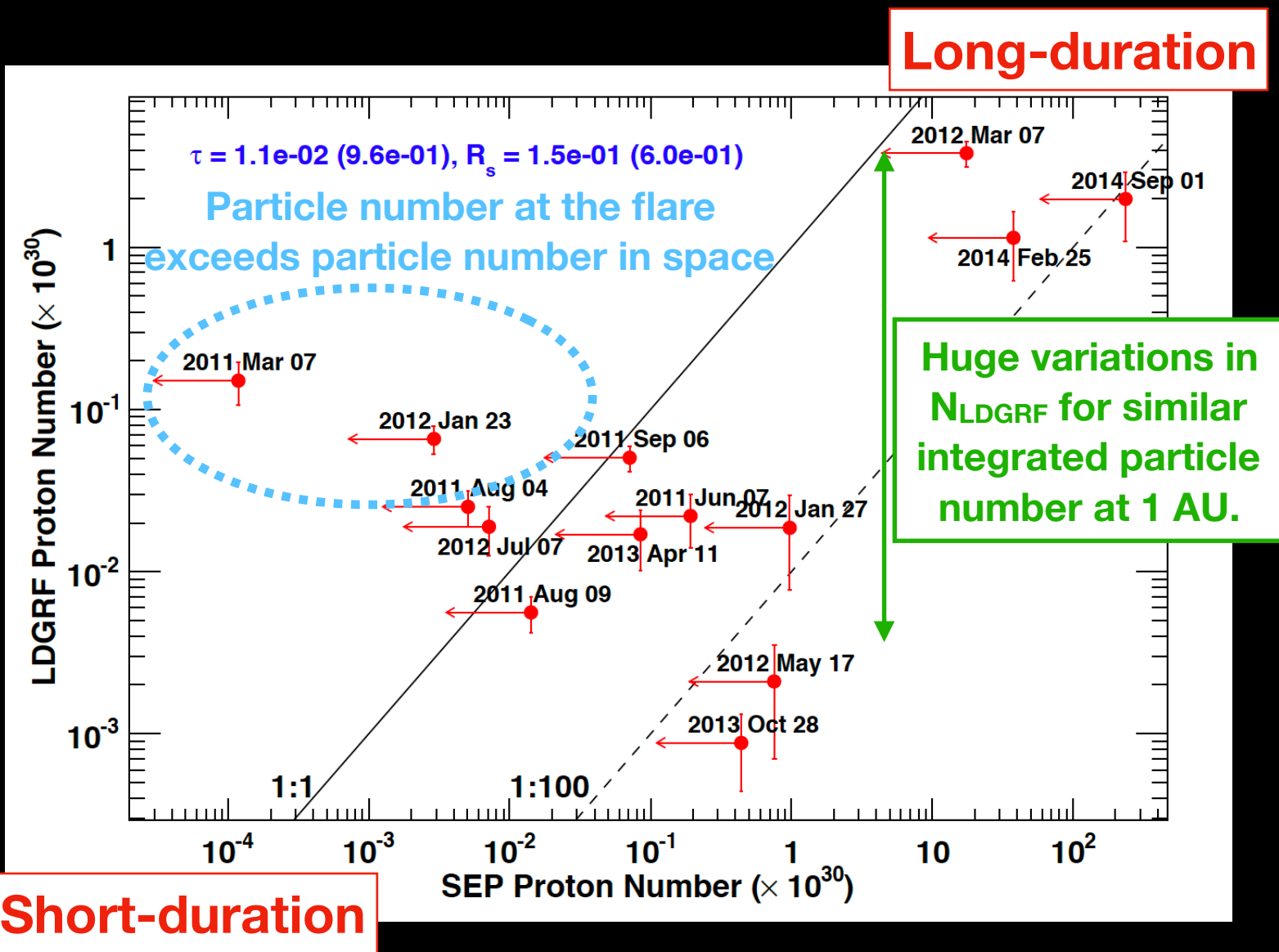
- \* No correlation (low values of the Kendall's  $\tau$  and Spearman rank correlation coeffs).
- \*  $N_{SEP}/N_{LDGRF}$  ratio spans  $> 5$  decades of magnitude from  $7.8 \times 10^{-4}$  to  $\sim 5.0 \times 10^2$

# Precipitation Rate : $N_{LDGFR}/(N_{LDGFR}+N_{SEP})$



Total number of protons (those that escape as SEPs plus those that produce LDGRFs) that would have to precipitate to account for the LDGRF emission.

# Does Back Precipitation from CME-driven Shocks Work?



— Huge variations could be the result of sporadic & unpredictable magnetic connectivity, although such widely varying connectivity isn't supported by the smoothly decaying LDGRF emission from Fermi/LAT

- Large  $N_{LDGRF}$  number with nearly 80% precipitation would imply :
- 1) an enormous loss channel for the shock
  - 2) high shock formation heights, resulting in a weakening shock, adding to the challenge of accelerating particles to high-energy.

— Additionally challenges:

2012 Oct 23 & 2012 Nov 27 exhibit LDGRF emission but have no CMEs (and likewise examples of fast, full-halo CMEs with no  $> 100$  MeV  $\gamma$ -ray emission)

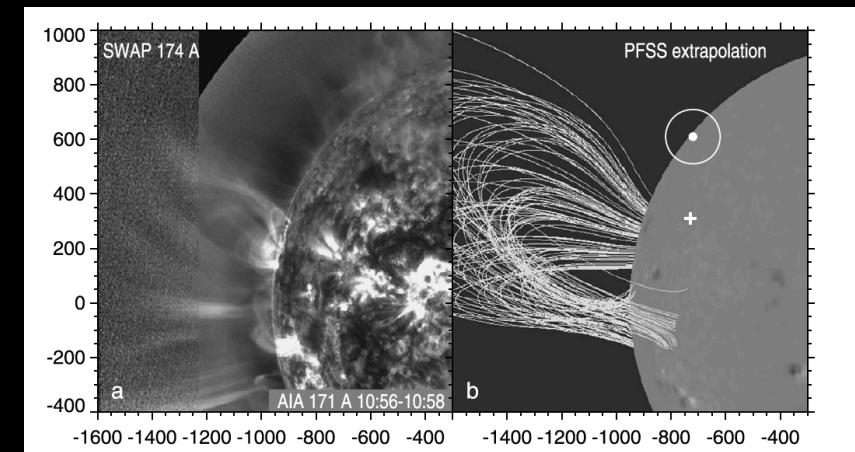


# Alternate Scenario: Trapping in Large Coronal Loops

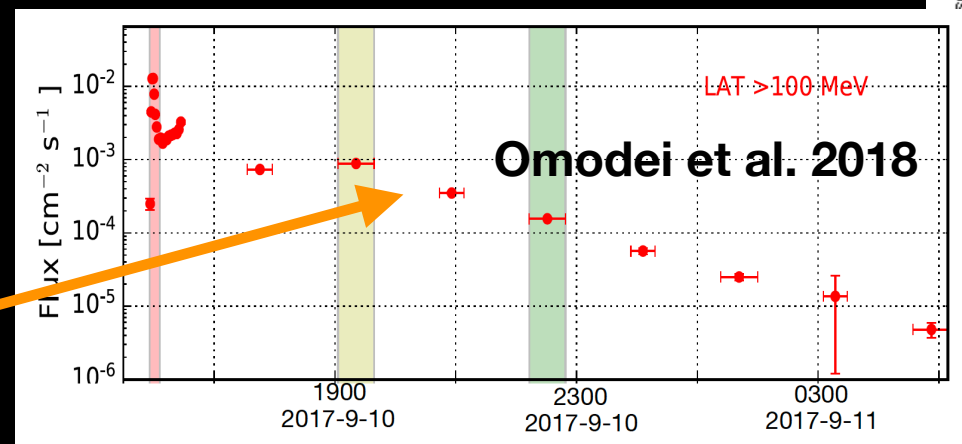
Particle acceleration occurs via second-order *Fermi* mechanism & trapping occurs locally within extended coronal loops, & ions diffusing to the denser photosphere to radiate (Ryan & Lee 1991).

1) Grechnev et al. (2019) provided evidence through radio observations (NRH) that the behind-the-limb flare of 2014 Sep 1 involved two distinct quasi-static loops of different sizes with emission consistent with prolonged confinement (and perhaps reaccelerating)

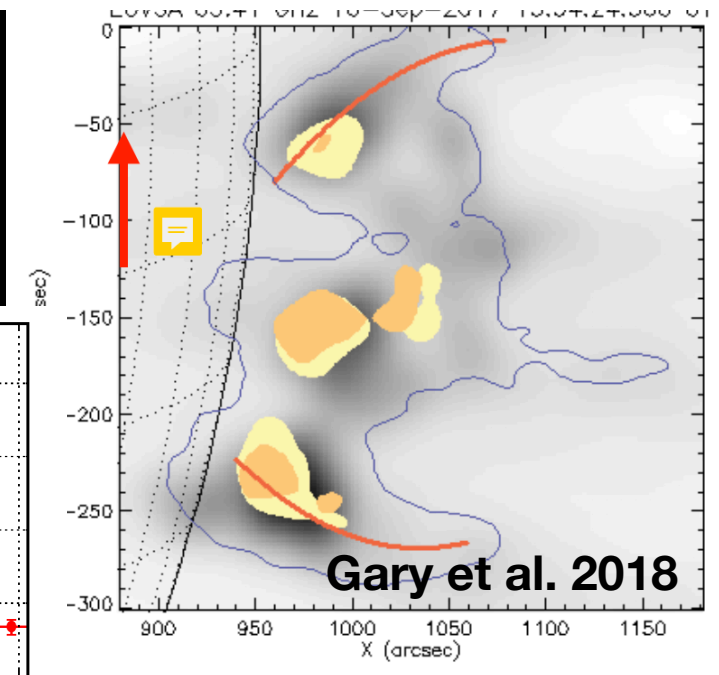
Grechnev et al. 2018



2) Gary et al. (2018) used microwave (EOVSA) observations to show footpoints of a large coronal loop for the 2017 Sep 10 flare with circular length of  $1.4 R_s$  & microwave emission persisting into the period of the  $> 100$  MeV  $\gamma$ -ray emission.



Omodei et al. 2018



Gary et al. 2018

⇒ Smooth, robust exponential decay argues for coronal trap scenario, with spatial and momentum diffusion governing the precipitation (see Ryan et al. this ICRC)

# What Have We Learned?

- Unique observations of PAMELA SEPs cover the energy range of interest for studying LDGRFs (above pion production threshold of  $\sim 300$  MeV)
  - $\Rightarrow$  Observe spectral roll-overs important as well as properly accounting for the spatial distribution and transport ( $N_{\text{cross}}$ )
- $N_{\text{SEP}}$  is not correlated with  $N_{\text{LDGRF}}$ 
  - $\Rightarrow$  Observe large variations (ratio spans 5 orders in magnitude)
- $\Rightarrow$  Precipitation rates place challenging constraints on CME shocks as the source of LDGRFs (see *de Nolfo et al. 2019*).
- $\Rightarrow$  An alternate explanation for LDGRF emission is coronal trapping/acceleration which decouples the SEPs from the interacting protons and where the effects of diffusion are consistent with smooth, exponentially decaying  $\gamma$ -ray light curves. Recent observations support the existence of large, persistent coronal loops and modeling efforts are promising (see *Ryan et al. this ICRC*).

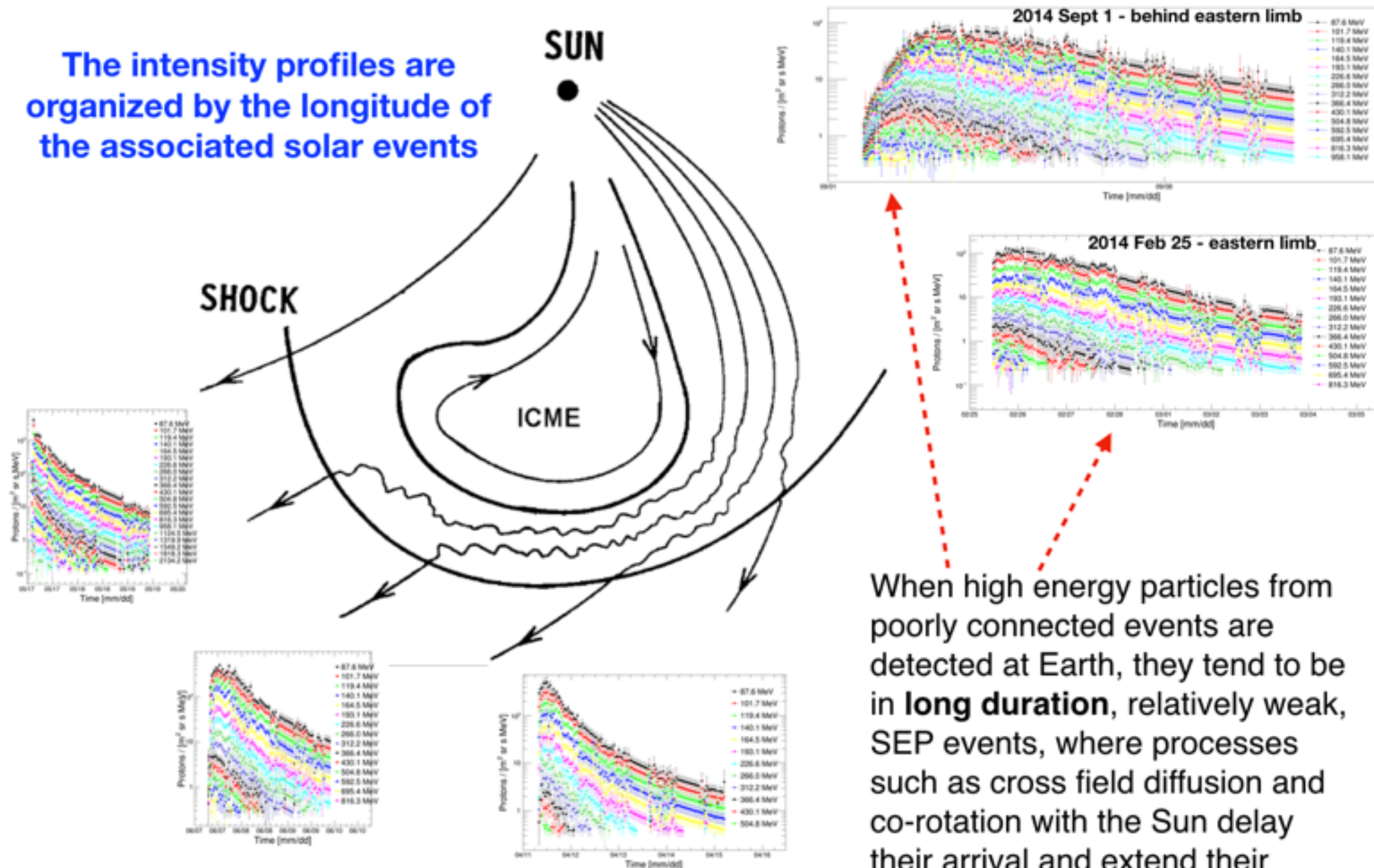
# Back ups



# SEP Transport

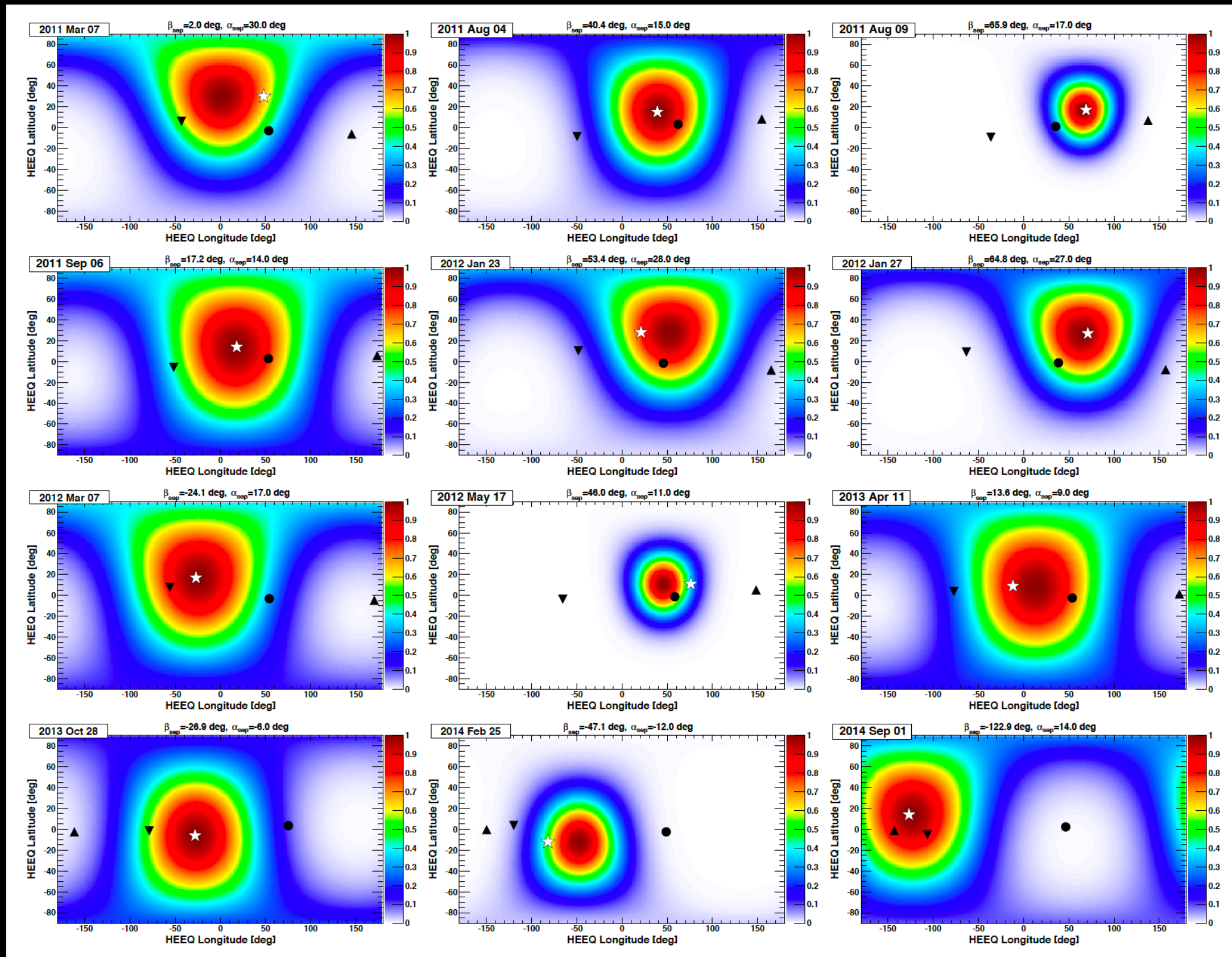
Longer time profiles for events that are not well connected

The intensity profiles are organized by the longitude of the associated solar events



When high energy particles from poorly connected events are detected at Earth, they tend to be in **long duration**, relatively weak, SEP events, where processes such as cross field diffusion and co-rotation with the Sun delay their arrival and extend their duration at Earth.

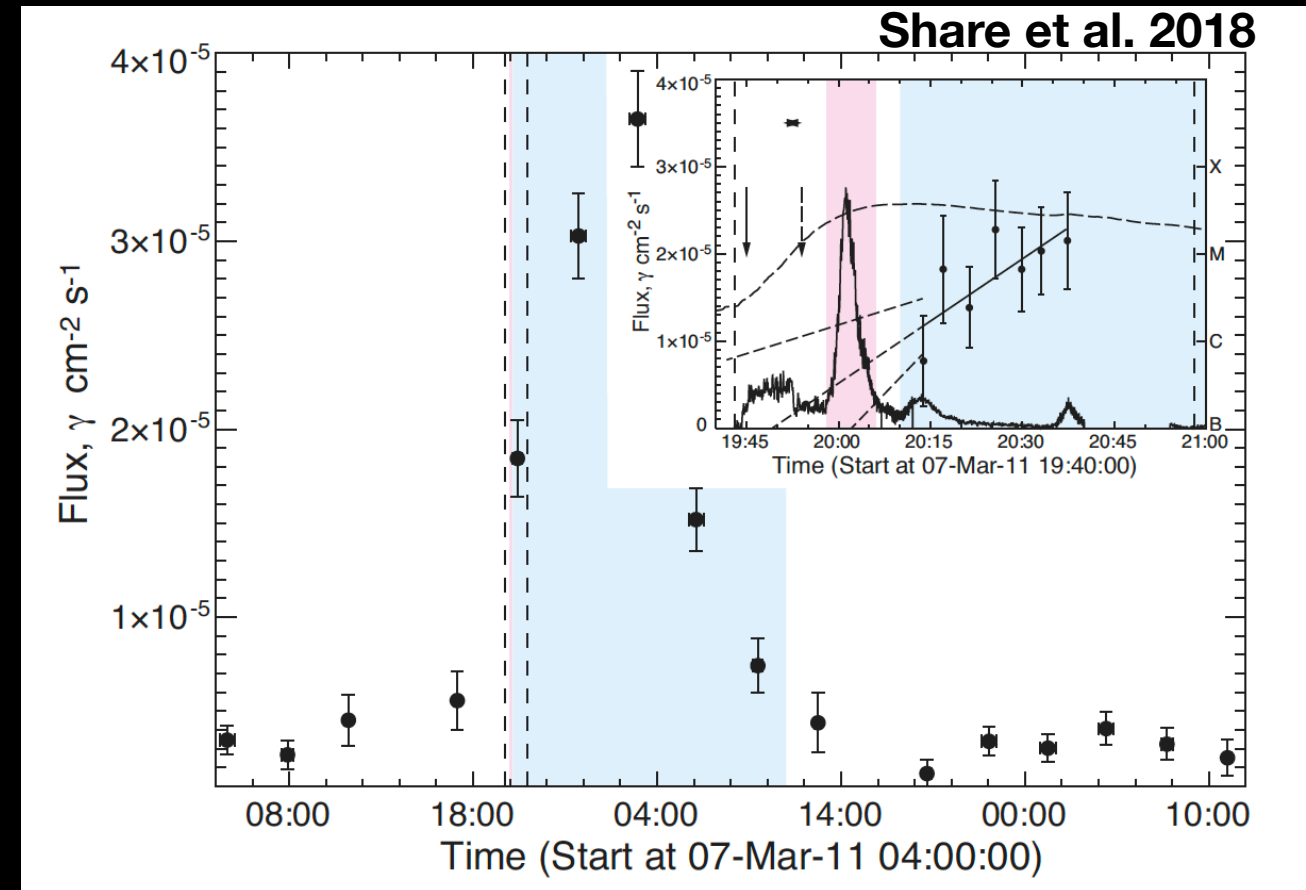
# SEP spatial distribution in HEEQ coordinates based on event-integrated fluences $> 80$ MeV



- ★ flare location
- PAMELA
- ▲ STEREO

# Introduction : Long Duration Gamma-ray Flares

- \* Delayed and prolonged  $>50$  MeV  $\gamma$ -ray emission well after the impulsive phase.
- \* High-energy gamma-ray emission ( $>100$  MeV) is thought to originate primarily from the decay of pions, produced by protons (and alpha) particles above  $\sim 300$  MeV (above  $\sim 200$  MeV). See for instance Ackermann et al. 2017)



Time history of  $> 100$  MeV gamma-ray flux from Fermi/LAT. Inset compares with GBM 100-300 keV & dashed curve is soft x-rays

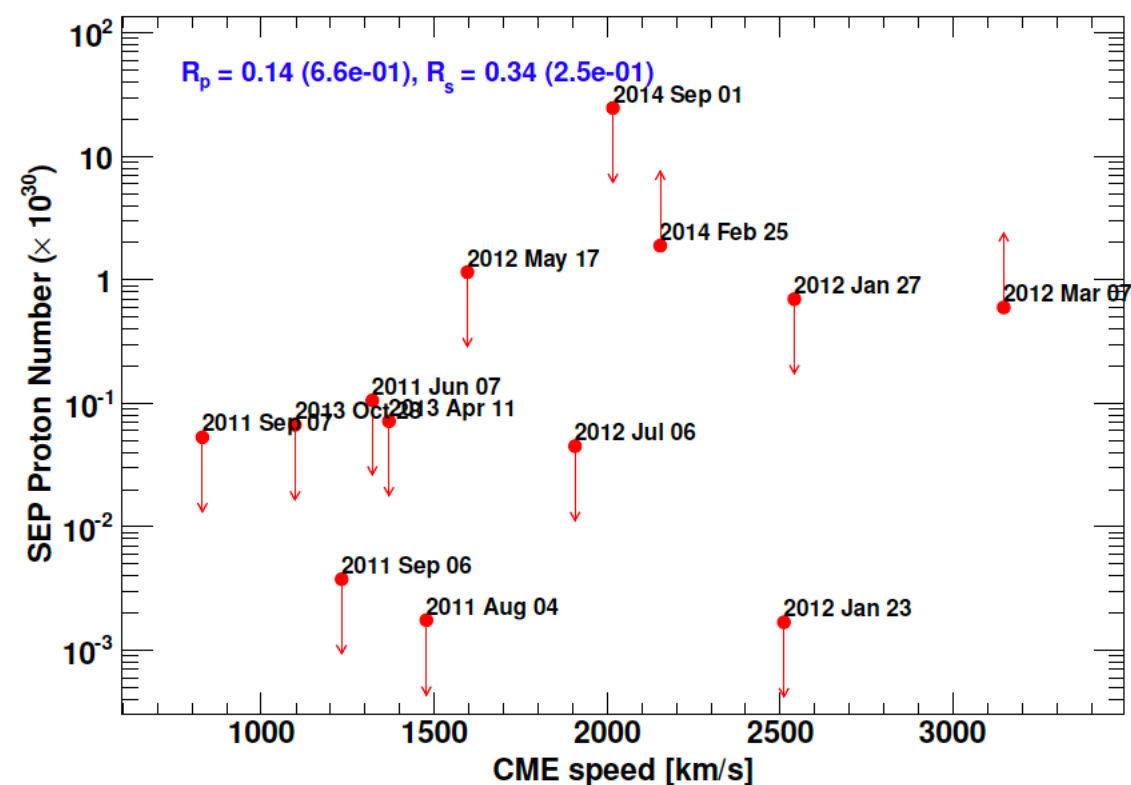
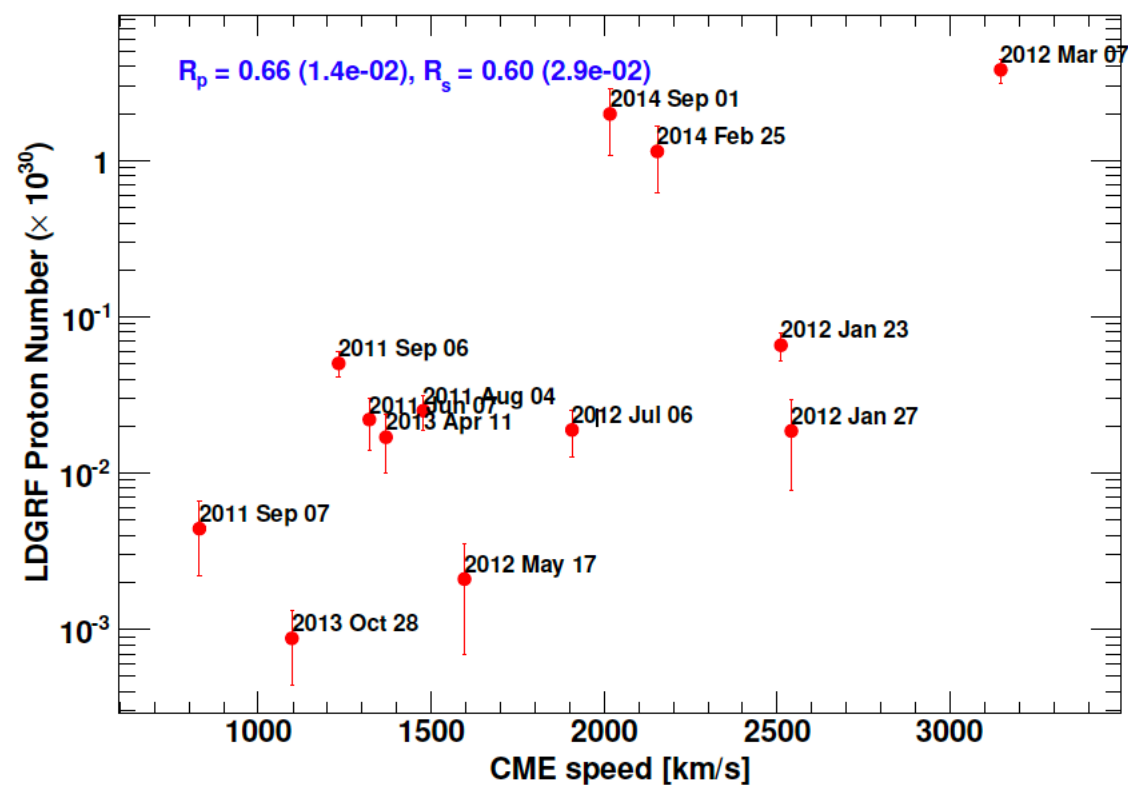
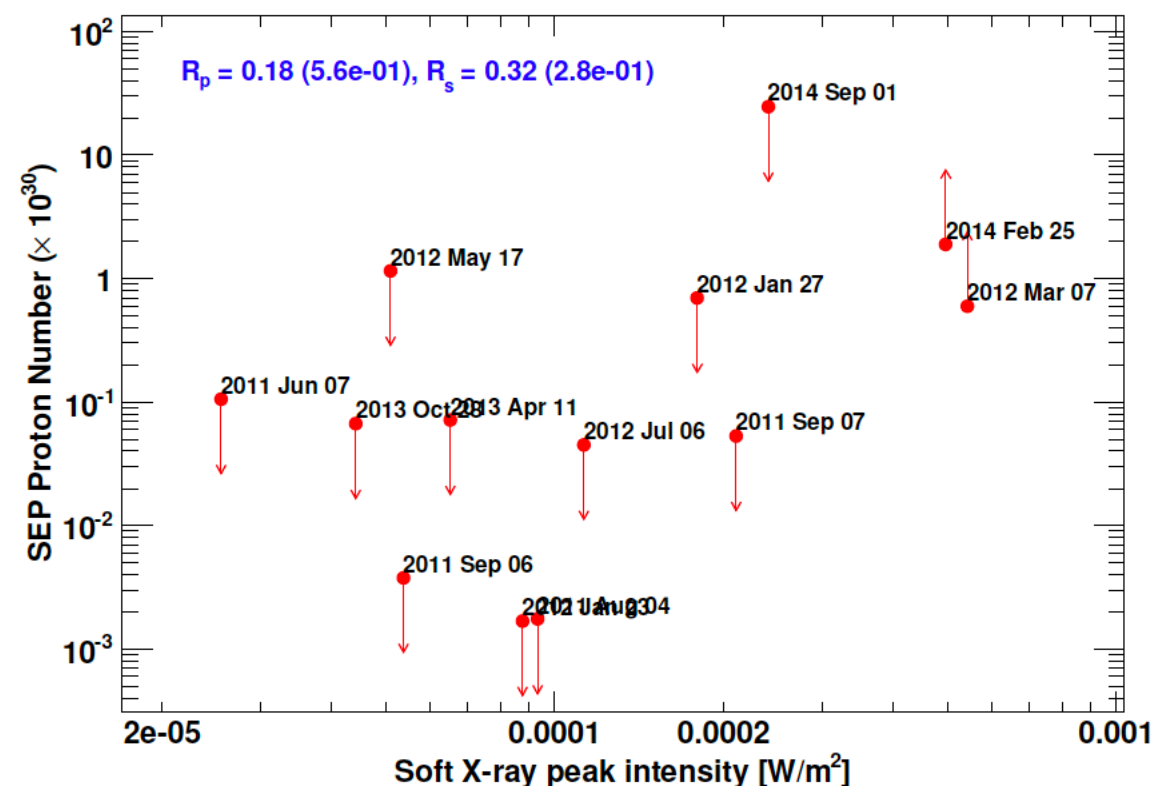
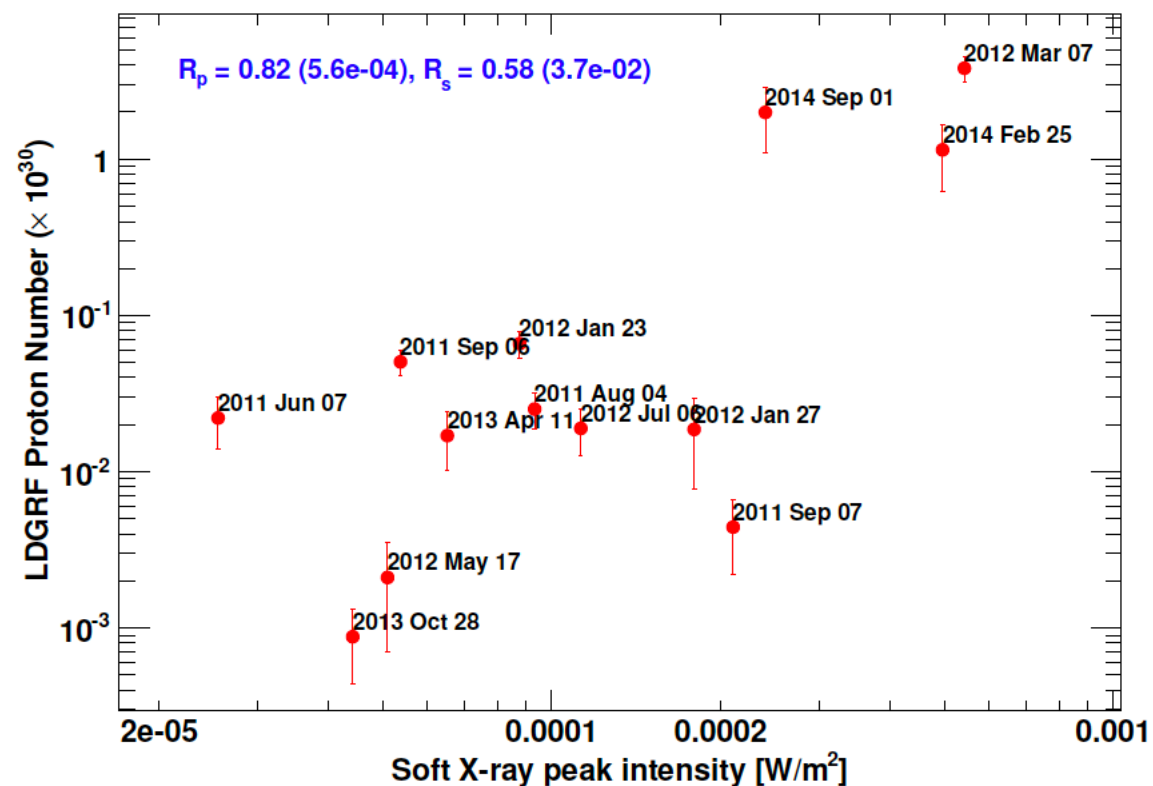
Recent comprehensive work by Share et al. 2018 on “Late Phase Gamma-ray Emission”

1. Delayed  $\gamma$ -ray emission is associated with impulsive  $> 100$  keV x-ray emission
2. There are x10 more protons accelerated in delayed phase than impulsive
3. There is a delay from CME onset of a few minutes to  $\sim 6$  hours
4. Most LPGREs were associated with CMEs, Type II & III radio emission
5.  **$> 500$  protons needed to produce LPGRE range from 0.1 to 50% the number of protons in SEPs**

=> Use direct observations of high-energy spectra from PAMELA and improve statistics

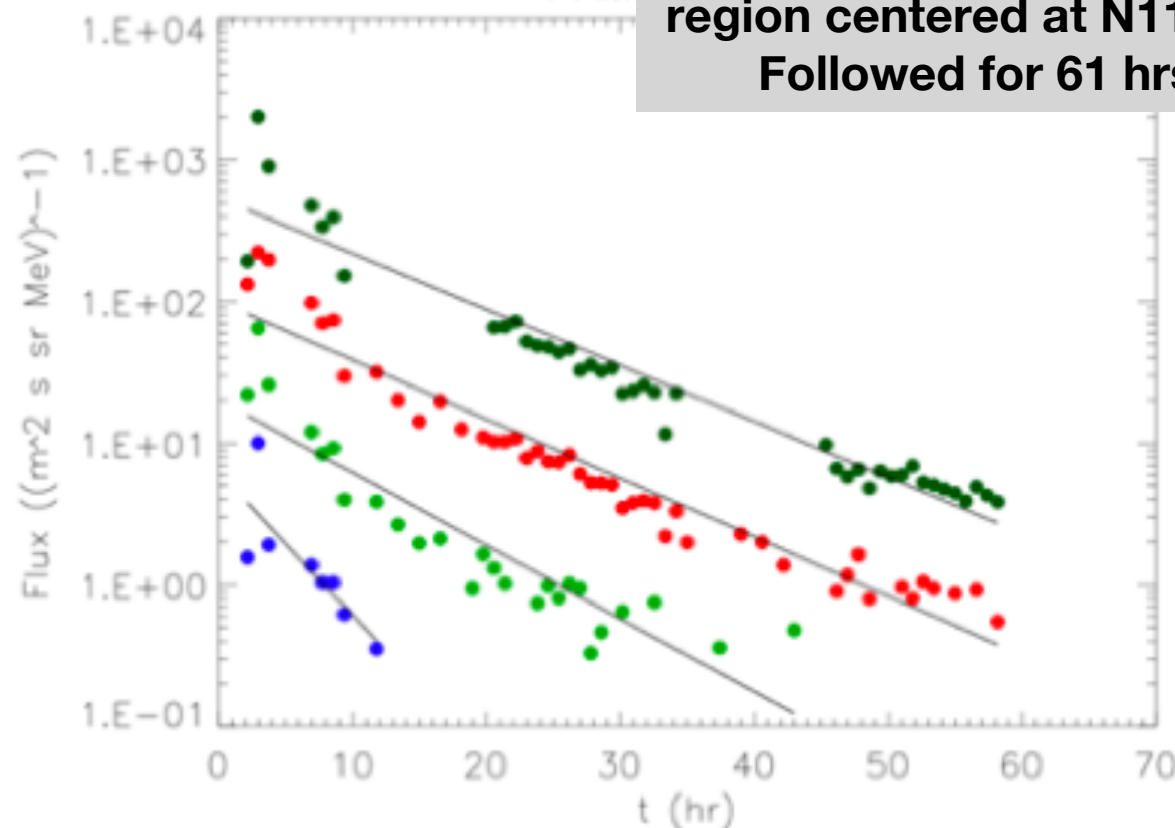


# Comparing $N_p$ to CME, X-ray Flare Parameters



**Battarbee et al. 2018 model  
simulation for GLE 2012 May 17**

**2012 May 17**



injection @ 2  $R_s$  over  $40 \times 40^\circ$   
region centered at N11W76  
Followed for 61 hrs

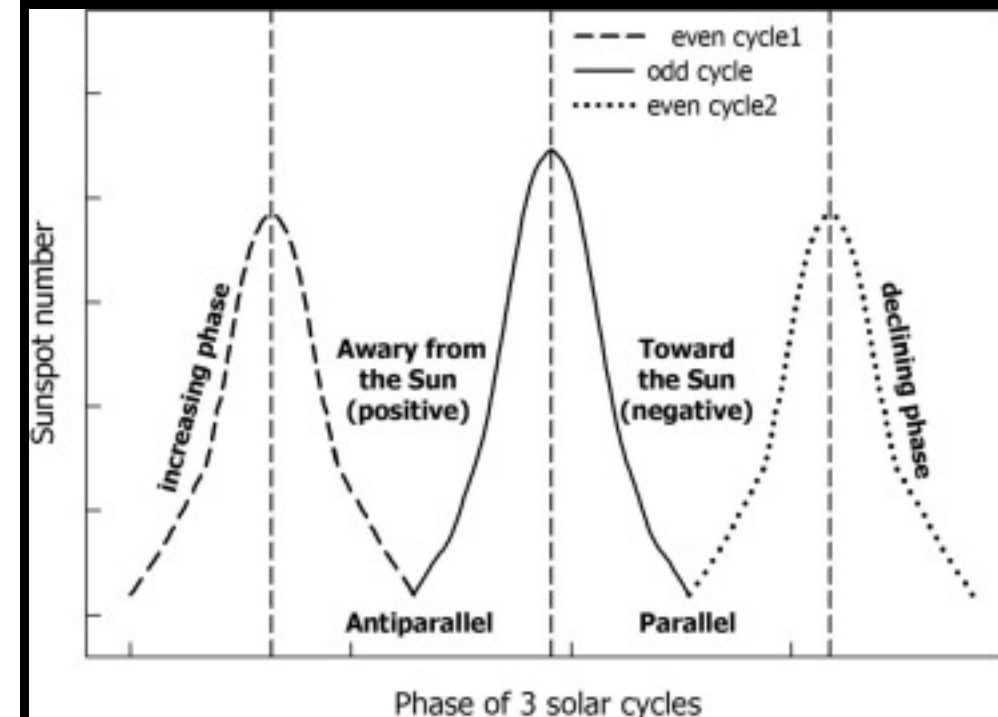
**Full simulation of the 2012 May 17 event, with initial proton distribution given by a power law, shows good agreement with PAMELA intensity time profiles for the case where  $\lambda = 0.3$  AU (Dalla et al., in preparation).**

**$\lambda = \text{const}$ , crossings over entire 1 AU sphere are added together and averaged over the monoenergetic population considered.**

**We assume  $\lambda = 0.5$  AU (and wavy HCS) conservatively in order to compute upper limits to the proton numbers in space from SEPs**

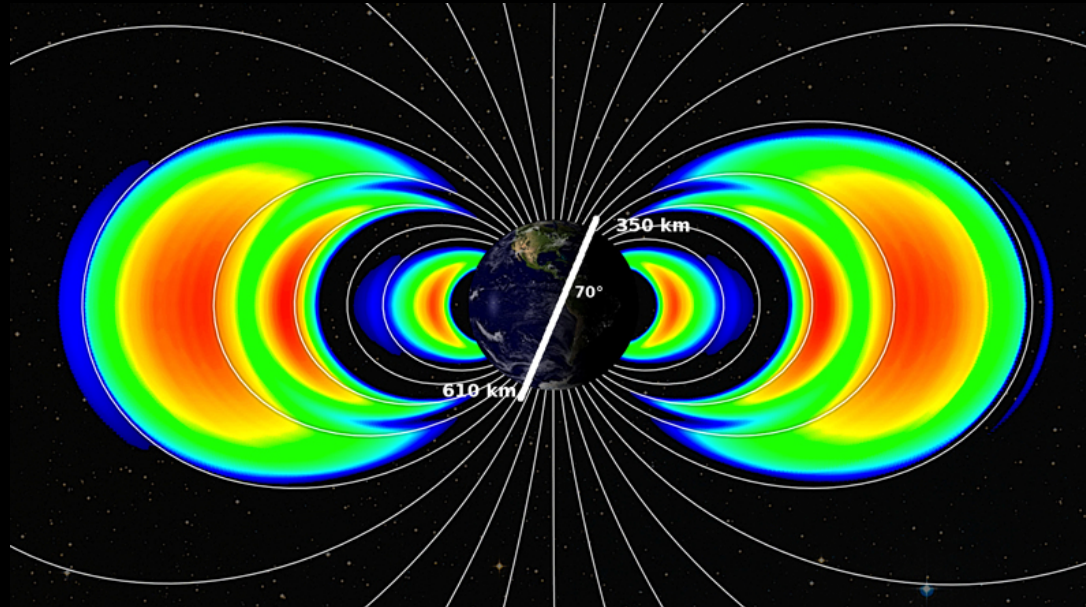
# HCS Polarity Effects

Energy (MeV)	$\lambda$ (AU)	HCS configuration	Polarity	$N_{crossings}$
1000	0.1	no HCS	+ve everywhere	31
1000	0.1	no HCS	-ve everywhere	28
1000	0.1	flat HCS	$A^+$	15
1000	0.1	flat HCS	$A^-$	28
1000	0.1	wavy HCS	$A^+$	17
1000	0.1	wavy HCS	$A^-$	29
500	0.1	flat HCS	$A^+$	19
500	0.1	flat HCS	$A^-$	29
500	0.5	flat HCS	$A^+$	8
500	0.5	flat HCS	$A^-$	11



**The number of crossings depends on the solar polarity and varies for different configurations of the Heliospheric Current Sheet (HCS)**

# Payload for Antimatter Matter Exploration and Light-nuclei Astrophysics (PAMELA)



PAMELA: ~circular orbit (580 km altitude);  
70° inclination

PAMELA is sensitive to the  
energy range (~80 MeV to several  
GeV) that corresponds to the  
interacting ions at the Sun that  
produce the LDGRFs!

