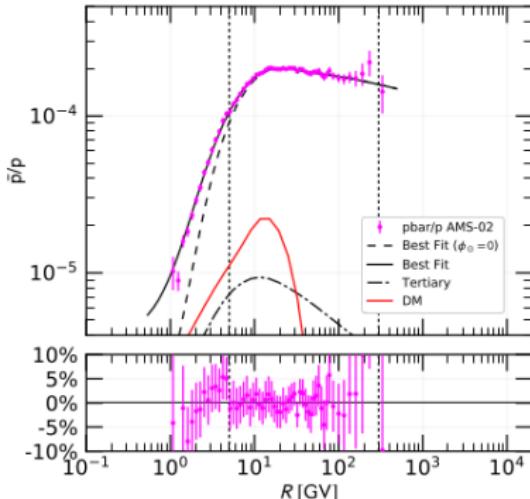


# Solar Modulation of Cosmic Rays in a Semi-Analytical Framework

*Marco Kuhlen, Philipp Mertsch*

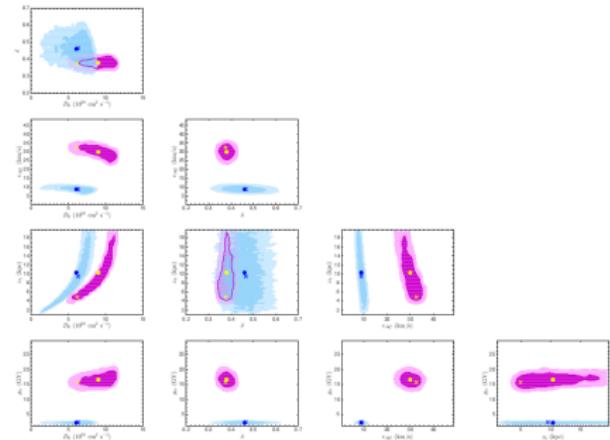
Institute for Theoretical Particle Physics and Cosmology (TTK)

# Importance of Solar Modulation



Understand cosmic ray fluxes at low energies to interpret potential dark matter signals.

Cuoco et al. arXiv:1903.01472



Disentangle modulation of galactic cosmic rays from processes in the heliosphere.

Johannesson et al. arXiv:1602.02243

# Solving the Transport Equation

computational expense

## Force-field

- ✓ fast
- ✗ inaccurate
- ✗ local

Gleeson & Axford 1968

## Numerical codes

- ✗ slow
- ✓ accurate
- ✓ global

Aslam et al.,  
arXiv:1811.10710,  
Boschini et al.,  
arXiv:1704.03733,  
Vittino et al.,  
arXiv:1707.09003,  
Kappl, arXiv:1601.02832

# Solving the Transport Equation

computational expense



Force-field	Semi-analytical	Numerical codes
✓ fast	✓ fast	✗ slow
✗ inaccurate	✓ accurate	✓ accurate
✗ local	✗ local	✓ global

Gleeson & Axford 1968

Aslam et al.,  
arXiv:1811.10710,  
Boschini et al.,  
arXiv:1704.03733,  
Vittino et al.,  
arXiv:1707.09003,

Kappl, arXiv:1601.02832

# Force Field Approximation

Gleeson & Axford, Caballero-Lopez & Moraal

Rewrite the transport equation as

$$\frac{\partial f}{\partial t} + \nabla \cdot (C\mathbf{V}f - \mathbf{K} \cdot \nabla f) + \frac{1}{3p^2} \frac{\partial}{\partial p} (p^3 \mathbf{V} \cdot \nabla f) = Q,$$

with Compton Getting factor  $C \equiv -\frac{p}{3} \frac{1}{f} \frac{\partial f}{\partial p}$ .

## Assumptions:

- ▶ Steady state,  $\partial f / \partial t = 0$
- ▶ No sources,  $Q = 0$
- ▶ No average momentum loss in lab frame,  $\langle \dot{p} \rangle = \frac{1}{3} \mathbf{V} \cdot \nabla f / f = 0$

Zero streaming condition:

$$C\mathbf{V}f - \mathbf{K} \cdot \nabla f = 0.$$

# Force Field Solution

Gleeson & Axford, Caballero-Lopez & Moraal

Assuming spherical symmetry:

$$\frac{\partial f}{\partial r} + \frac{Vp}{3\kappa} \frac{\partial f}{\partial p} = 0,$$

Method of characteristics:

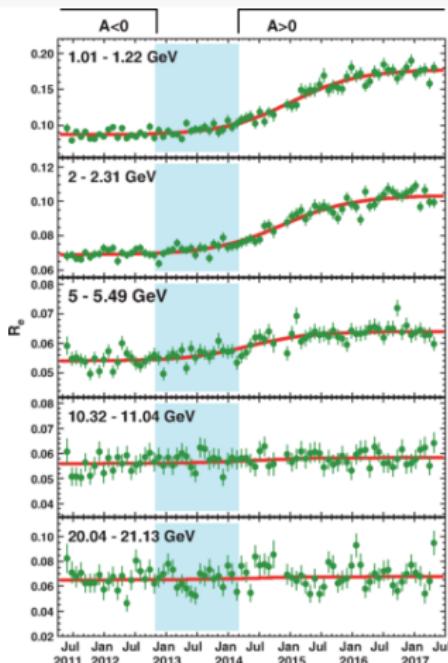
$$\int_{p_{TOA}}^{p_{LIS}} \frac{\beta \kappa_{p'}}{p'} dp' = \int_{r_{TOA}}^{r_{LIS}} \frac{V}{3\kappa_{r'}} dr' \equiv \phi(r),$$

For  $\kappa_p \propto p$  and  $\beta \approx 1 \rightarrow \phi = p_{LIS} - p_{TOA}$ .

Conservation of the phase-space density  $f$  leads to:

$$\frac{J_{TOA}}{p_{TOA}^2} = \frac{J_{LIS}}{p_{LIS}^2}.$$

## Time-Dependent Experimental Data



AMS Collaboration

Phys. Rev. Lett. 121, 051102

$$R_e = \frac{\Phi_{e^+}}{\Phi_{e^-}}$$

Explanation of current data requires charge sign dependent effects.

⇒ Importance of drifts.

To explain the AMS-02 data we make modifications to the force field model.

# Changes to Force Field Model

Where we introduce angular averages

Starting again from divergence free streaming

$$\int_S (C\mathbf{V}f - \mathbf{K} \cdot \nabla f) \cdot d\mathbf{S} = 0$$

Solve the transport equation in 2D including gradient curvature drifts.

$$\tilde{f} = \int_0^{\pi/2} d\theta \sin \theta f$$

$$\tilde{V} = \left( \frac{\partial \tilde{f}}{\partial p} \right)^{-1} \int_0^{\pi/2} d\theta \sin \theta V \partial_p f$$

$$\tilde{K}_{rr} = \left( \frac{\partial \tilde{f}}{\partial r} \right)^{-1} \int_0^{\pi/2} d\theta \sin \theta K_{rr} \partial_r f$$

$$\tilde{v}_{gc,r} = \left( \tilde{f} \right)^{-1} \int_0^{\pi/2} d\theta \sin \theta v_{gc,r} f$$

# Changes to Force Field Model

After angular averages this reduces  
to

$$\frac{\partial \tilde{f}}{\partial r} + \frac{p \tilde{V}}{3\tilde{K}_{rr}} \frac{\partial \tilde{f}}{\partial p} = -\frac{\tilde{v}_{gc,r}}{\tilde{K}_{rr}} \tilde{f}$$

With  $p_{\text{LIS}}(r, p)$  the solution to the  
initial value problem

$$\frac{dp}{dr} = \frac{p \tilde{V}}{3\tilde{K}_{rr}},$$

Can be solved using the method of  
characteristics

$$\tilde{f}(r, p) = f_{\text{LIS}}(p_{\text{LIS}})$$

with  $p_{\text{LIS}}(R, p) = p$ .

$$e^{-\int_0^r dr' \frac{\tilde{v}_{gc,r}(r', p'_{\text{LIS}})}{K_{rr}(r', p'_{\text{LIS}})}}$$

# Changes to Force Field Model

After angular averages this reduces to

$$\frac{\partial \tilde{f}}{\partial r} + \frac{p \tilde{V}}{3\tilde{K}_{rr}} \frac{\partial \tilde{f}}{\partial p} = -\frac{\tilde{v}_{gc,r}}{\tilde{K}_{rr}} \tilde{f}$$

Can be solved using the method of characteristics

$$\tilde{f}(r, p) = f_{\text{LIS}}(p_{\text{LIS}}) e^{-\int_0^r dr' \frac{\tilde{v}_{gc,r}(r', p'_{\text{LIS}})}{\tilde{K}_{rr}(r', p'_{\text{LIS}})}}$$

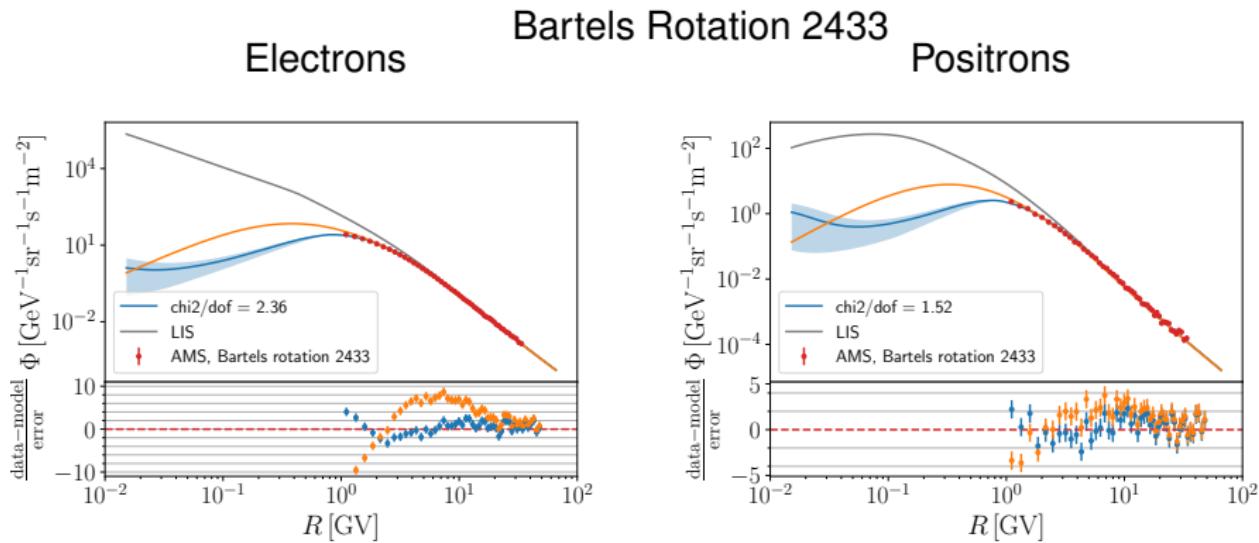
We parametrize them as

$$\tilde{V} = V_0(1 + \Delta V \theta(p - p_b))$$

$$\tilde{K}_{rr} = K_0 R^a \left( \frac{R^c + R_k^c}{1 + R_k^c} \right)^{(b-a)/c}$$

$$\tilde{v}_{gc,r} = \kappa_0 \frac{\beta p}{3B_0} \frac{10 p^2}{1 + 10 p^2}$$

## Example: Fit to AMS-02 Data

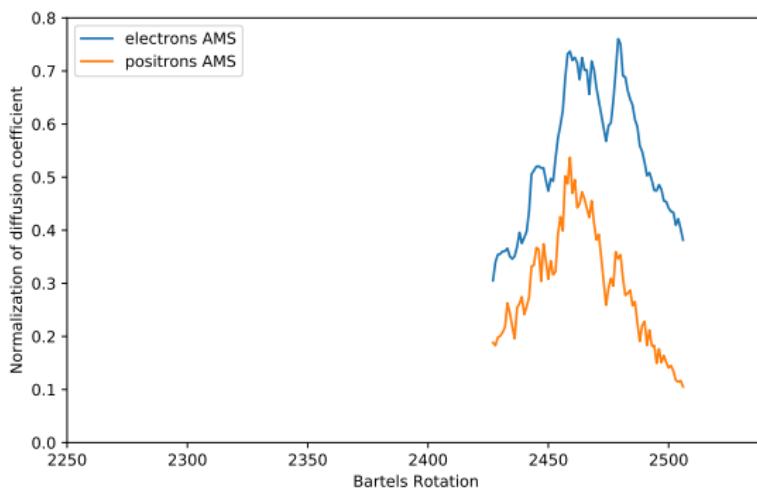


Can explain data accurately while the conventional force field model fails.

LIS from Vittino et al. arXiv:1904.05899

# Correlation with Solar Wind Parameters

We find a correlation between the tilt angle and the diffusion coefficient.



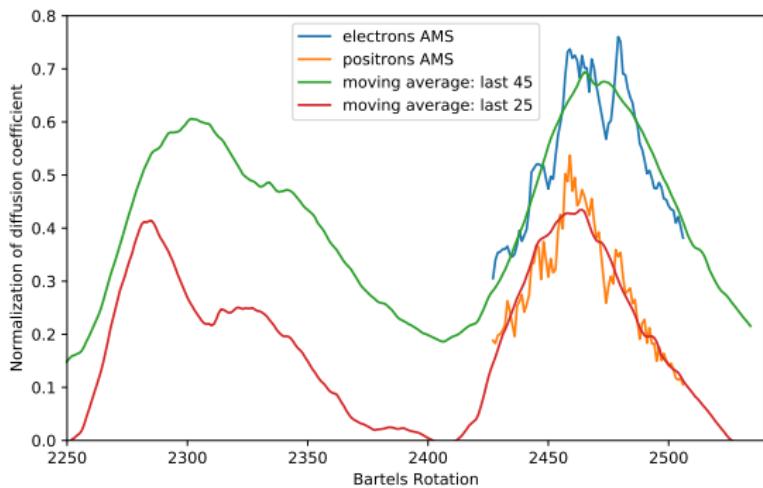
$$\tilde{K}_{rr}^- = K_0^- R^a \left( \frac{R^c + R_k^c}{1 + R_k^c} \right)^{(b-a)/c}$$

$$\tilde{K}_{rr}^+ = K_0^+ R^a \left( \frac{R^c + R_k^c}{1 + R_k^c} \right)^{(b-a)/c}$$

Tilt angle from <http://wso.stanford.edu/>

## Correlation with Solar Wind Parameters

We find a correlation between the tilt angle and the diffusion coefficient.



$$\tilde{K}_{rr}^- = K_0^- R^a \left( \frac{R^c + R_k^c}{1 + R_k^c} \right)^{(b-a)/c}$$

$$K_0^- = a^- \langle \alpha \rangle_{\Delta t} + b^-$$

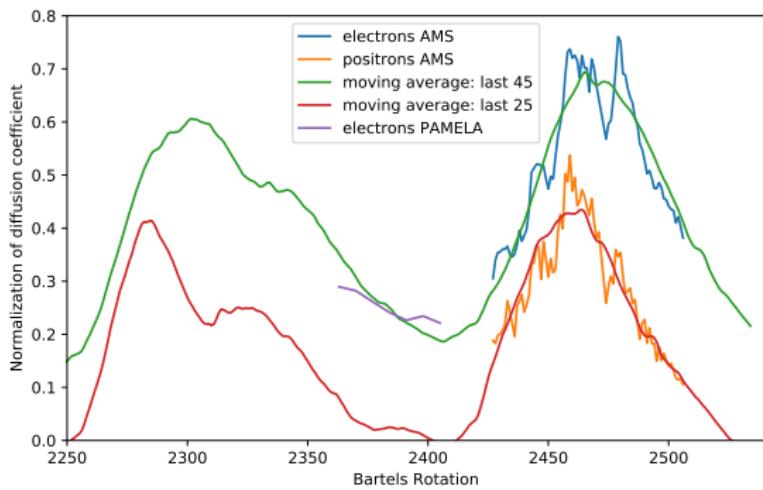
$$\tilde{K}_{rr}^+ = K_0^+ R^a \left( \frac{R^c + R_k^c}{1 + R_k^c} \right)^{(b-a)/c}$$

$$K_0^+ = a^+ \langle \alpha \rangle_{\Delta t} + b^+$$

Tilt angle from <http://wso.stanford.edu/>

## Correlation with Solar Wind Parameters

We find a correlation between the tilt angle and the diffusion coefficient.



$$\tilde{K}_{rr}^- = \textcolor{blue}{K}_0^- R^a \left( \frac{R^c + R_k^c}{1 + R_k^c} \right)^{(b-a)/c}$$

$$\textcolor{green}{K}_0^- = a^- \langle \alpha \rangle_{\Delta t} + b^-$$

$$\tilde{K}_{rr}^+ = \textcolor{orange}{K}_0^+ R^a \left( \frac{R^c + R_k^c}{1 + R_k^c} \right)^{(b-a)/c}$$

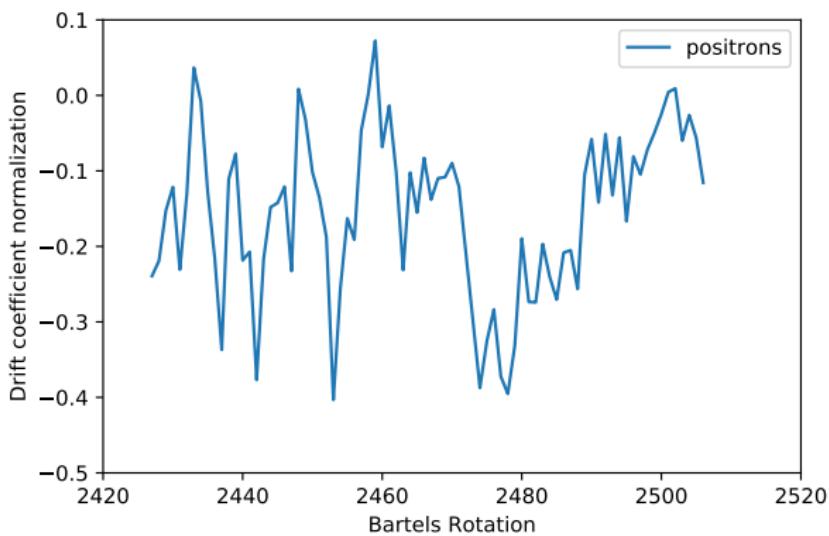
$$\textcolor{red}{K}_0^+ = a^+ \langle \alpha \rangle_{\Delta t} + b^+$$

$$\tilde{K}_{rr}^- = \textcolor{violet}{K}_0^- R^a \left( \frac{R^c + R_k^c}{1 + R_k^c} \right)^{(b-a)/c}$$

Tilt angle from <http://wso.stanford.edu/>

# Correlation with Solar Wind Parameters

We find a weaker correlation between the magnetic field strength and the normalization of the drift coefficient.

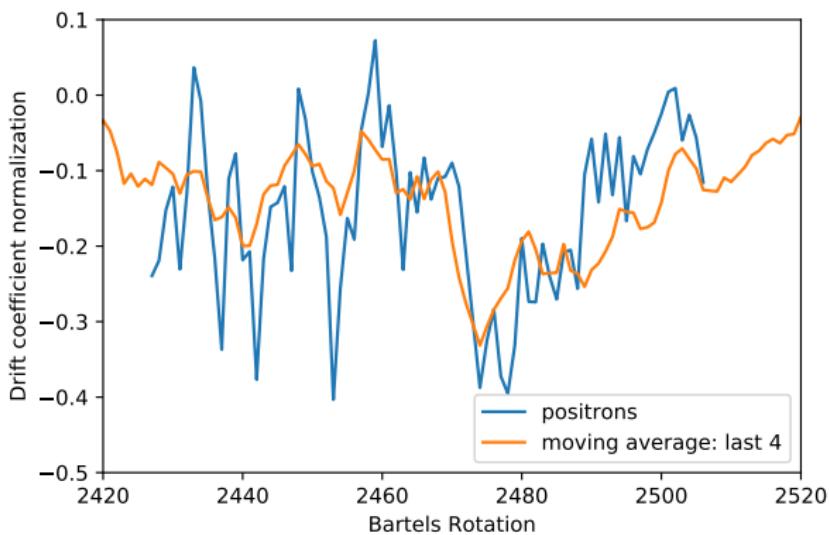


$$\tilde{v}_{gc,r}^+ = \kappa_0^+ \frac{\beta p}{3B_0} \frac{10 p^2}{1 + 10 p^2}$$

Magnetic field strength from <http://www.srl.caltech.edu/ACE/>

## Correlation with Solar Wind Parameters

We find a weaker correlation between the magnetic field strength and the normalization of the drift coefficient.



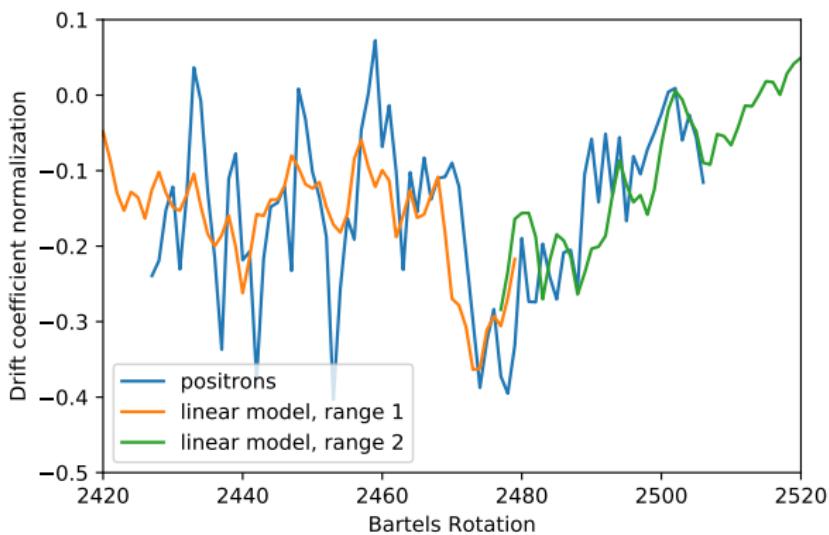
$$\tilde{v}_{gc,r}^+ = \kappa_0^+ \frac{\beta p}{3B_0} \frac{10 p^2}{1 + 10 p^2}$$

$$\kappa_0^+ = c_0^+ \langle B_0 \rangle_{\Delta t} + d_0^+$$

Magnetic field strength from <http://www.srl.caltech.edu/ACE/>

## Correlation with Solar Wind Parameters

We find a weaker correlation between the magnetic field strength and the normalization of the drift coefficient.



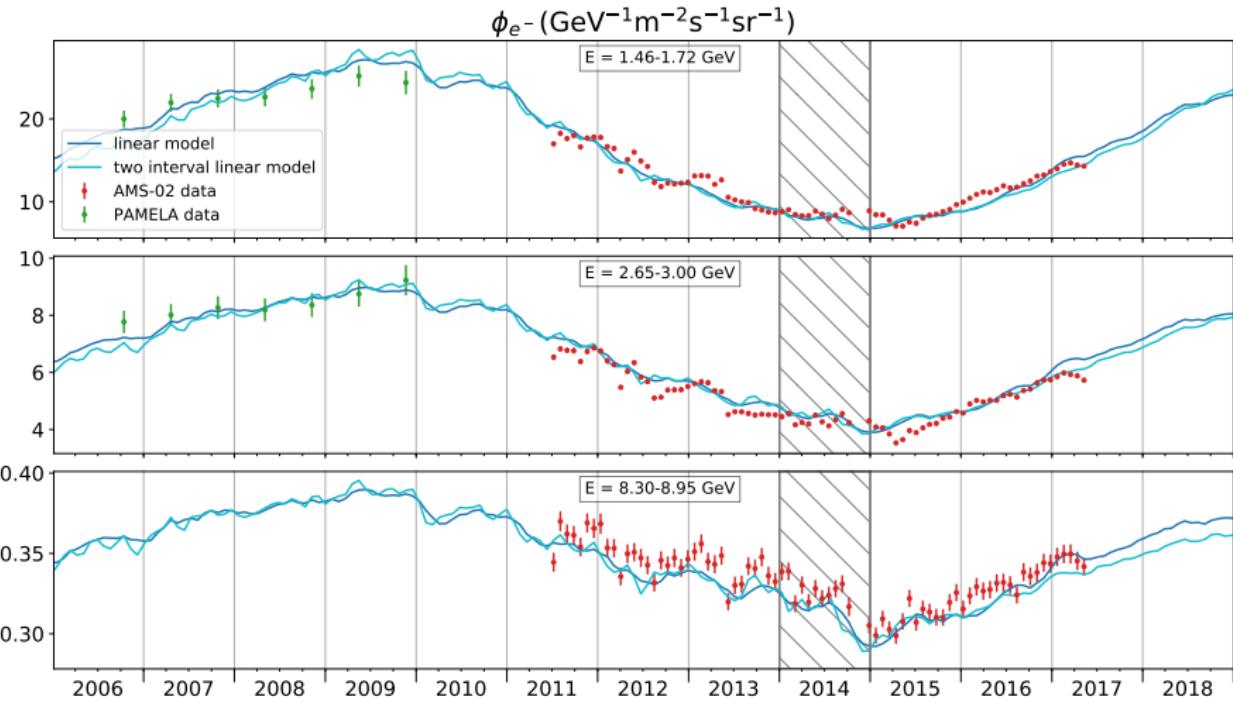
$$\tilde{v}_{gc,r}^+ = \kappa_0^+ \frac{\beta p}{3B_0} \frac{10 p^2}{1 + 10 p^2}$$

$$\kappa_0^+ = c_0^+ \langle B_0 \rangle_{\Delta t} + d_0^+$$

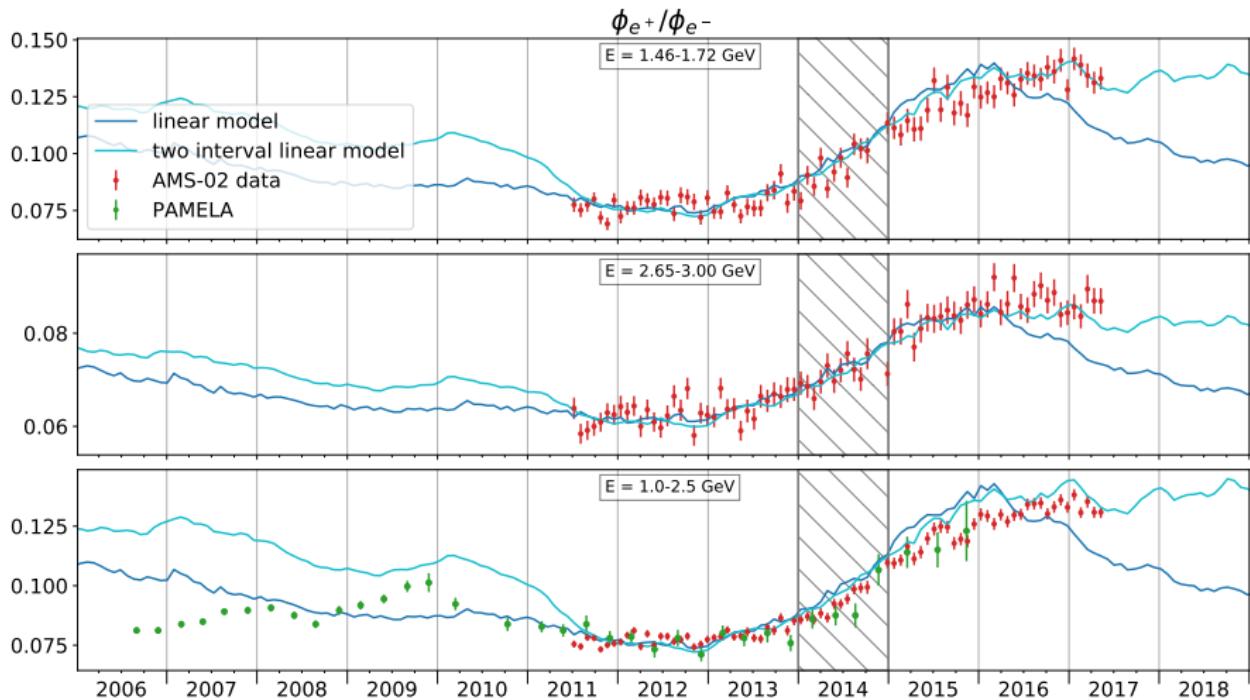
$$\kappa_1^+ = c_1^+ \langle B_0 \rangle_{\Delta t} + d_1^+$$

Magnetic field strength from <http://www.srl.caltech.edu/ACE/>

## Prediction of Electron Flux



## Prediction of Positron Ratio



# Conclusion

- ▶ We have developed a semi analytical method to solve the 2D transport equation.
- ▶ Our method runs significantly faster than fully numerical model ( $\sim 20$  ms).
- ▶ We are able to reproduce AMS-02 electron and positron fluxes.

Thank you for your attention!

Download our code at  
<https://git.rwth-aachen.de/kuhlenmarco/effmod-code>